Abstract. To probe the dynamics and size of lipid rafts in the membrane of living cells, the local diffusion of single membrane proteins was measured. A laser trap was used to confine the motion of a bead bound to a raft protein to a small area (diam $\leq 100$ nm) and to measure its local diffusion by high resolution single particle tracking. Using protein constructs with identical ectodomains and different membrane regions and vice versa, we demonstrate that this method provides the viscous damping of the membrane domain in the lipid bilayer. When glycosylphosphatidylinositol (GPI)-anchored and transmembrane proteins are raft-associated, their diffusion becomes independent of the type of membrane anchor and is significantly reduced compared with that of nonraft transmembrane proteins. Cholesterol depletion accelerates the diffusion of raft-associated proteins for transmembrane raft proteins to the level of transmembrane nonraft proteins and for GPI-anchored proteins even further. Raft-associated GPI-anchored proteins were never observed to dissociate from the raft within the measurement intervals of up to 10 min. The measurements agree with lipid rafts being cholesterol-stabilized complexes of $26 \pm 13$ nm in size diffusing as one entity for minutes.

Key words: laser trap • lipid raft • protein diffusion • single particle tracking • thermal position fluctuation analysis

Introduction

Membrane microdomains enriched in glycosphingolipids and cholesterol and containing glycosyolphosphatidylinositol (GPI)$^1$-anchored proteins have been proposed as lateral structural components of the plasma membrane (Simons and Ikonen, 1997). These microdomains have been implicated in the polarized sorting of proteins (Keller and Simons, 1998) and cellular signaling (Stauffer and Meier, 1997; for reviews see Brown and London, 1997, 1998). Imagining these microdomains as floating complexes in the membrane, Simons and Ikonen (1997) referred to them as lipid rafts. To understand the role of rafts in cells, a characterization of their dynamics is needed. Therefore, we have studied the local diffusion of single rafts in the plasma membrane of living cells.

Rafts are thought to form by self-association of sphingolipids because of their long and mostly saturated hydrocarbon chains. The interaction between glycosphingolipids can be enhanced by hydrogen bonds between their headgroups. The voids between the hydrocarbon chains caused by the rather bulky headgroups would be filled by cholesterol, which might also participate in the hydrogen bonding to the sphingolipids. A alternative view of sphingolipid–cholesterol rafts is that they form separate liquid-ordered phases in the bilayer, which are dispersed in the liquid-disordered matrix formed by unsaturated glycerophospholipids (Brown and London, 1997; Ge et al., 1999). The liquid-ordered phase in a lipid bilayer, as characterized by highly ordered carbon chains coupled with a high degree of rotational freedom, was first described by Ipsen et al. (1987).

Biochemically, the components of lipid rafts are characterized by their insolubility in the detergent Triton X-$100$ at $4^\circ$C, forming so-called detergent insoluble glycolipid-enriched complexes (DIGs) that are enriched in cholesterol, glycosphingolipids, sphingomyelin, and saturated glycerophospholipids (Brown and Rose, 1992; Parton and Simons, 1995). Because of their high lipid content DIGs

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$^1$Abbreviations used in this paper: DIG, detergent insoluble glycolipid-enriched complex; GFP, green fluorescent protein; GPI, glycosyolphosphatidylinositol; HA, influenza virus hemagglutinin; LFPGT46, artificial transmembrane YFP; PLAP, placental alkaline phosphatase; SPT, single particle tracking; TFR, transferrin receptor; TPF, two-photon fluorescence; YFP, yellow color variant of green fluorescent protein; YFP-GLGPI, artificial GPI-anchored YFP.
are found in the low density fraction after gradient centrifugation (Brown and Rose, 1992; Friedler et al., 1993). Proteins that are DIG-associated after the detergent extraction are assumed to have been raft-associated in the membrane. Depletion of cholesterol from cells renders raft-associated proteins detergent soluble (Klisdonk et al., 1995; Scheiffele et al., 1997; Keller and Simons, 1998).

The existence of lipid rafts in the plasma membrane of living cells has been demonstrated recently, but controversy persists and the exact structure of rafts remains unclear. Varma and Mayor (1998) used fluorescent resonance energy transfer between GPI-anchored proteins in living CHO cells to show that these proteins were clustered in lipid domains that were likely to be smaller than 70 nm in diameter explaining why rafts have never been observed in light microscopy. While another fluorescent resonance energy transfer study failed to detect any significant clustering of GPI-anchored proteins in MDCK cells (Kenworthy and E didin, 1998), Friedrichson and K urzhalia (1998) have shown that GPI-anchored proteins in MDCK cells exist in clusters by using short chemical cross-linkers. Harder et al. (1998) used antibodies against raft-associated proteins to cocluster the small rafts into large super rafts, which were detectable by light microscopy. To understand the role of rafts in cell signaling and protein sorting, a characterization of the stability and mobility of solitary rafts in the plasma membrane of living cells is needed. Studying the local diffusion of single raft proteins over minutes with high resolution should provide information about the dynamic properties of rafts, because the protein diffusion in a simple lipid bilayer depends on the size of the protein and the viscosity of the surrounding lipid bilayer as described by the Saffman-Delbrück relation (1975). Diffusion measurements in cell membranes using FRAP or video-based single particle tracking (SPT) have so far deviated from this simple relation. This discrepancy is thought to be caused by collisions of the diffusing protein with immobile obstacles and fences in the cell membrane (Saxton, 1982, 1989; K usumi and Sako, 1996). Locally, on a length scale smaller than the average distance between obstacles, the cell membrane can be approximated as a lipid bilayer. Therefore, we have measured the local viscous drag of single raft proteins using a novel microscopic method (Florin et al., 1998; Pralle et al., 1998). A microsphere, attached to the protein studied, is confined by a laser trap and the sphere’s thermal position fluctuations inside the trapping potential are tracked with subnanometer and microsecond resolution. The damping of these fluctuations yields directly the viscous drag \( \gamma \) on the sphere and the attached protein. The viscous drag determined with a temporal resolution of 0.3 s relates to the diffusion coefficient \( D \) of a freely diffusing protein via the Einstein relation: 

\[
D = k_B T / \gamma \quad (k_B T \text{ being the thermal energy})
\]

The local diffusion coefficients obtained with this method for protein diffusion in the plasma membrane are the first to agree with the coefficients obtained with this method for protein diffusion measurements in cell membranes using FRAP or Spectra Physics, is directed by a scanning mirror mounted on a triple-axis piezo (PiezoSystems Jen) into the microscope and focused in the object plane by an objective immersion objective lens (Plan Neofluor 100×, NA 1.3; Carl Zeiss). The condenser back focal plane (BFP) is projected onto a quadrant photodiode (QPD) (model S9981; Hamamatsu) via a dichroic mirror. To reduce the mechanical noise, the microscope’s stage is replaced by a custom made stage to which the objective lens and the condenser with the detection system are rigidly connected. The two-photon fluorescence (TPF) intensity is measured by a photomultiplier (model R2949; Hamamatsu). All signals recorded are amplified and low-pass-filtered (50 kHz) by amplifiers developed at EMBL. A transputer data acquisition board (model ADWin F5; Jäger Electronics) is used for recording and to provide the axial feedback controlling the interaction force between the sphere and the membrane. To move the trapping focus along the optical axis, the objective lens is mounted on a piezo (model Physik Instrumente). In addition, the temperature-controlled sample chamber is placed on a piezo scan stage with capacitive position sensors and digital feedback (model NPS-XY-100A; Queensgate).

To determine the average axial position of the sphere in the trap, the fluorescence intensity emitted by the fluorophores inside the spheres excited by the trapping laser is measured (Florin et al., 1996). This signal is used as input for the feedback minimizing the force exerted by the sphere on the membrane. For the high resolution particle tracking, we detect the interference of the forward scattered laser light with the reference beam with a QPD in the BFP. The lateral position is the difference of two halves of the QPD (Gittes and Schmidt, 1998), whereas the axial position is determined from the intensity of the scattered laser light reaching the QPD (Pralle et al., 1999). For a sphere with 0.2–μm diam, the lateral and axial resolution at 50 kHz bandwidth are better than 1 and 5 nm, respectively. Because the response of the QPD depends on the properties of sphere and the laser focus, it is necessary to calibrate the signal with the sphere used for an experiment at a location near the actual measurement. The local detector sensitivity \( \beta \) is determined from the thermal position fluctuations using the Stokes drag \( \gamma \) of the sphere. A analogous to the local viscosity measurements (see below), the autocorrelation time of the position movements \( \gamma \) and the spring constant of the trap are calculated, being now an uncalibrated spring constant \( k \) in units Nm/Volt instead of N/m. Because \( \gamma = k \) and \( \gamma = k \beta \), the sensitivity \( \beta \) is determined from the square amplitude of \( F(t, t) = F(t) F(t) \) with the mean square amplitude

\[
F^2 = k_B T / \gamma \quad \text{and the correlation time } \tau = /\gamma .
\]

Thus, the local viscous drag \( \gamma \) and the diffusion coefficient \( D = k_B T / \gamma \) of a sphere in a harmonic potential as long as the position fluctuations remain within the linear response range of the detector, and the calibration is performed in the solution. We calibrated the detector to be at least ten times the diameter of the sphere away from any surface, so the influence of the surface is smaller than 2% (Happe and Brenner, 1965; Pralle et al., 1998).

### Local Viscosity Measurement

The motion of a Brownian particle in a harmonic potential is characterized by an exponentially decaying position autocorrelation function \( \langle \gamma (t) \rangle = \gamma (t) \exp (-t/\tau) \) with the mean square amplitude

\[
\tau^2 = k_B T / k \quad \text{and the correlation time } \tau = /\gamma .
\]

Thus, the local viscous drag \( \gamma \) and the diffusion coefficient \( D = k_B T / \gamma \) of a sphere in a harmonic potential are calculated from the measured correlation time \( \tau \) of the motion and the stiffness \( k \) of the potential (Pralle et al., 1998).

The stiffness of a potential can be determined by measuring the position distribution of a trapped particle (Florin et al., 1998) using the Boltz-
mann probability $P(r)\,dr$ to find a thermally excited particle in a potential $V(r)$ at position $r$ in the interval $(r, r + dr)$ given $P(r) = c \cdot \exp[-V(r)/kT]$, with $c$ determined normally $P(0\,\text{r}) = 1$. Conversely, the trapping potential is given by the probability distribution as $V(r) = -kT \cdot \ln P(r) + kT \cdot \ln c$, where $c$ is an offset. This method allows to profile the trapping potential even below the thermal energy with a temporal and spatial resolution given by the strength of the potential and the bead size, while requiring only minimal knowledge about the system, i.e., the temperature. In our experiments, the lateral spring constant of the laser trap was adjusted to about $k \approx 1 \mu N/m$ for a sphere of $0.2-\mu$m diameter. The sample chamber was maintained at $36 \pm 1\°C$, which leads to lateral position fluctuations of $\pm 60\,\text{nm RMS}$ displacement.

To achieve diffusion coefficient measurements with errors <10%, the observation interval has to be $\approx 1,000$-fold longer than the correlation time $\tau$ (estimated using methods developed by Bartlett, 1946). Hence, the temporal resolution of the viscosity measurement is limited by the sphere’s motion and not by the bandwidth of the detection. The sphere’s position was recorded every $18.75\,\mu$s for some time-long observations the recording interval was increased to $50\,\mu$s. The three-dimensional position distribution of the bead was used to visualize the cellular surface in the area of the measurement. Only experiments in which the surface was normal to the optical axis of the microscope were selected for further analysis.

The position fluctuations along each detector axis were evaluated independently. To determine $\tau$, overlapping intervals of $0.3\,\mu$s were used. The autocorrelation function in each interval was calculated and fitted by an exponential decay from 0.15 to 25 ms. The potential $\kappa$ and the viscous drag $\gamma$ were computed for each interval and plotted against time.

Over the time course of an experiment, the viscous drag of the free bead, of the same bead near the membrane, and bound are determined. The position fluctuations along each detector axis were evaluated independently. To determine $\tau$, overlapping intervals of $0.3\,\mu$s were used. The autocorrelation function in each interval was calculated and fitted by an exponential decay from 0.15 to 25 ms. The potential $\kappa$ and the viscous drag $\gamma$ were computed for each interval and plotted against time.

The transiently transfected cells had normal morphology and were viable for several days. LYFP GT46 and YFPGLGPI were expressed in Ptk2 cells by infection at 1.3°C with recombinant adenoviruses encoding for the constructs, and experiments were performed 12–36 h after transfection. For cholesterol depletion, the cells were grown on coverslips and washed in serum-free culture medium and incubated with $10\,\text{mM}$ methyl-β-cyclodextrin for 30 min at 37°C. This procedure was reported to extract $\approx 60\%$ of total cellular cholesterol from BHK cells (Keller and Simons, 1998). After a wash in normal culture medium, the coverslips were transferred to the microscope. The measurements were performed within 1 h of the cholesterol depletion.

### Protein Constructs and Generation of Recombinant Adenoviruses

The transient expression of PLAP in BHK and the BHK clone stably expressing PLAP has been described previously (Harder et al., 1998). YFPGLGPI, containing YFP as the ectodomain connected to a GPI anchor, was produced and characterized by Keller (Keller, P., D. Tooomey, J. White, and K. Simons, manuscript in preparation). The expression of HA in BHK cells has been characterized previously (Schellie et al., 1997). LYP FTG46 is an artificial secretory protein containing the signal sequence of rabbit lactase-phlorizin hydrolase (Mantei et al., 1988), the fluorescent protein, a consensus N-glycosylation site, the transmembrane domain of the human LDL receptor (Yamamoto et al., 1984), and the cytoplasmic tail of CD46 (Müller et al., 1997; Teuwen et al., 1999) containing the sequence stretch FTSL that allows basolateral transport. A DNA fragment encoding the signal sequence was generated by PCR with an Nhel site and a consensus Kozak sequence at the 5’ end, and a Pst site at the 3’end. Primers were used as follows: 5’-TA GC TCT GAC CAA CAG TG TCC TTT TAG CTG AGA G-3’ (forward primer) and 5’-ATC TCT GAG A ATG CTG TCT CTA TGC-3’ (reverse primer). This PCR fragment was ligated into Nhel-Pst-digested YEP-NL yielding the secreted protein LYP FTG46. Next, we inserted the CD46 cytoplasmic tail into LYP FTG46. This was accomplished with two codon-optimized complementary oligonucleotides creating a BsrGI site at the 5’ end and a NotI site at the 3’ end. The forward oligonucleotide was 5’-CTC AGA TCT GAC CAA CAG TG TCC TTT TAG CTG AGA G-3’ (forward primer, N-glycosylation site NGS in bold, LDL receptor sequence underlined) and 5’-CTT GCT TAC GAC CCG CCA GTA AAT TTG GGA AGA GAA TTA TTT CTT CTC TCG TGC GGT TCG GTG GTA GGG ATG CAC-3’ (TCC replacing the TAA stop codon of YFP is underlined). The reverse oligonucleotide was 5’-GGG CCG CTT TCG AAG GAA AGA GAA TTA TTT CTT CTC TCG TGC GGT TCG GTG GTA GGG ATG CAC-3’ (forward primer, N-glycosylation site NGS in bold, LDL receptor sequence underlined) and 5’-CAT GCT TGC AAG GAC TCT TAC CTT CTT GA AAC CAC TCT TGC GGA C-3’ (forward primer, N-glycosylation site NGS in bold, LDL receptor sequence underlined). This PCR fragment was directionally ligated into BsrGI-Not-digested LYP FTG46, yielding LYP FTG46. Finally, a DNA fragment encoding a consensus N-glycosylation site and the human LDL receptor transmembrane domain was generated by PCR. Primers used were as follows: 5’-CTG TAA CAA GCT TAC GAT GTC CCA GCA AAT TTG GGA AGA GAA TTA TTT CTT CTC TCG TGC GGT TCG GTG GTA GGG ATG CAC-3’. This PCR fragment was directionally ligationed into BsrGI-digested LYP FTG46, yielding LYP FTG46. Recombinant adenoviruses were generated with the system described by He et al. (1998). In brief, LYP FTG46 was released as a Nhel-Xbal fragment that was directionally ligated into XbaI-digested pShuttle-CMV. A titer homologous recombination with pAdefase-1, recombinant adenoviruses were generated in 293 cells as described (He et al., 1998). It was sorted to the basolateral membrane when expressed in MDCK cells and did not associate with rafts (Keller, P., unpublished results).

The cDNA for the EGF used for the blocking of the antibodies was subcloned by PCR from pEGFP-N1 (CLONTECH Laboratories) into a derivative of pET9 (Stratagene, modified by Gunter Stier, EMBL). 6xHis-tagged EGF was purified from Echerichia coli B121(DE3) using nickel affinity chromatography.

### Bead Coating and Antibodies

Range fluorescent (530 nm ex/560 nm em.) carbocyl-modified latex beads from Molecular Probes with a nominal diameter of 0.2 μm (radius = 108 ± 4 nm) were used. The beads were coated with the antibodies by adsorption, based on modified procedures of Sako and Kusum (1995). Before coating, the spheres were washed three times in 0.2 M boric acid buffer (adjusted to pH 9 using NaOH), and then they were incubated with 1 mg/ml antibody in a 50-mM MES buffer, pH 6, for 30 min at room temperature. After coupling, the spheres were incubated for 30 min with 10 mg/ml FSG and washed twice in 10 mg/ml FSG in PBS; another wash was added immediately before the experiment. To optimize the spheres for single membrane protein binding, they were incubated with free ligand, i.e., EGF (White, EMBL) or PLAP (Sigma Chemical Co.). In the

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case of EGFP, 15 nl of the EGFP solution (1 mg/ml) was added to 10 μl of a 1% aqueous suspension of antibody-coated spheres. Adding 50 nl of GFP inhibited binding to cells completely. Because of the lack of soluble forms of HA, the antibody concentration on the surface of the sphere was diluted by coadsorption of an unspecific antibody at a ratio of 1:4.

mAbs against human PLAP were obtained from DAKO. Rabbit polyclonal antibodies were raised against human PLAP (Sigma Chemical Co.). mAbs HA were prepared as described (Matlin et al., 1981; Scheiffele et al., 1997). Polyclonal antibodies against XFP were raised in mice (LeBot, EMBL).

Results

The viscous drag of an individual membrane protein in the plasma membrane of living cells, maintained in a closed microscope chamber heated to 36 ± 1°C, was measured by observing the thermal position fluctuations of an attached microsphere (radius = 108 nm; Fig. 1). The damping of the motion is expected to be dominated by the viscous drag on the membrane domain of the protein because of the 1,000-fold higher viscosity of the lipid bilayer compared with the aqueous medium. We used a laser trap to capture the sphere and place it onto the cell membrane while controlling the interaction force. The motion of the sphere was first recorded in the bulk solution, then near the membrane, and finally after binding to the membrane protein. The comparison of these three measurements allowed us to separate the influence of the sphere diffusing unbound near the membrane from the binding of the bead to the membrane protein. The continuous observation with high temporal resolution enabled us to directly observe the binding event. To demonstrate that the measurements were not influenced by the sphere size, some experiments were performed with smaller spheres (radius = 55 nm), yielding similar results (data not shown).

High Resolution Particle Tracking in a Trapping Potential

We implemented an optical trap in an inverted microscope using a near infrared laser that simultaneously excites the fluorophores inside the trapped microsphere via a two-photon process (Fig. 2). The dependence of the TPF intensity on the sphere’s position relative to the focus of the trapping laser provides an axial displacement sensor used to control the force exerted by the sphere on the membrane below 0.1 pN. To record the thermal position fluctuations, we developed a fast three-dimensional position sensor based on the interference of the transmitted laser light with the light scattered by the trapped particle (Gittes and Schmidt, 1998; Pralle et al., 1999). The interference pattern was detected with a quadrant photodiode in the backfocal plane of the condenser lens (Fig. 2).

Local Viscous Drag of Single Membrane Proteins

The particle tracking in a trapping potential allowed us to calculate the viscous drag and the trapping potential at a temporal resolution of 0.3 s and for an area of 100 nm in diameter. The viscous drag γ of the particle was computed as the product of the autocorrelation time τ of the position fluctuations and the spring constant k of the trapping potential (see Materials and Methods).

Fig. 3 depicts a typical measurement of the lateral viscous drag γ, spring constant k of the trapping potential, and autocorrelation time τ of the movement plotted

Figure 1. Scaled model of the experimental situation: a sphere (r = 108 nm) bound via an adsorbed antibody to a GPI-anchored protein that is part of a raft domain. The lipid bilayer is symbolized by the double row of gray dots with black sections symbolizing raft domains. The extent of the thermal position fluctuations observed in the experiments (± 60 nm) is marked. It is much smaller than the smallest estimates of the spacing of immobile cytoskeleton-anchored obstacles to free diffusion of 300-500 nm (Sako and Kusumi, 1995). The cytoskeletal elements drawn here are 250 nm apart.

Figure 2. Optical paths in our instrument, built around an inverted microscope with DIC equipment whose wavelength of 700 nm is chosen to reduce photon damage on the cells. The IR-laser trapping beam is focused on the sample by an oil immersion objective lens mounted on a piezo. The forward-scattered laser light is collected by the condenser lens and projected by a dichroic mirror onto the quadrant photodiode (QPD) for the particle tracking. The two-photon fluorescence (TPF) is detected confocally by a photomultiplier.
against time. The measured viscous drag $\gamma$ is the sum of the Stokes drag of the sphere $\gamma_s = 6\pi \eta_m r$ and the viscous drag of the protein in the lipid bilayer $\gamma_p$. The Stokes drag of the sphere was measured in the bulk solution, $\sim 2\ \mu m$ above the plasma membrane (Fig. 3 A, I). Near a surface, like the membrane, diffusion is reduced because of the spatial confinement (Happel and Brenner, 1965). Therefore, the viscous drag of the sphere was measured again after positioning it near the membrane (Fig. 3 A, II), and after binding to the membrane protein (III). The $\gamma_s$ (B) remains unchanged during the experiment, whereas $\tau_s$ (C) changes as $\gamma$ (A).

To ensure binding of the bead to a single membrane protein, we have reduced the number of proteins present in the cells by selecting for weakly expressing ones using the YFP fluorescence. The number of active antibody sites per bead was minimized in the case of anti-PLAP and anti-GFP beads by adding free ligand, a high concentration of which would block binding completely and, for HA, by replacing the major part of the antibodies on the bead with an unrelated antibody. The concentration of the free ligand was chosen to reduce the active binding sites per bead to below 10. The probability for single binding to surface proteins is very high due to the geometric restrictions. To verify this assumption, the fraction of beads bound and the time passed in contact with the membrane before binding were analyzed using Poisson statistics. Under the blocking conditions used, only a fraction of these beads was bound after being in contact with the membrane for 50 s (25% for PLAP, 30% for YFPGLGPI and HA, and 20% for LYFPGT46), whereas 95% of the unblocked beads bound within a few seconds. A Poisson statistic for single binding fitting this behavior provides an estimate of <25 possible binding sites to ensure single binding during the observation interval, and predicts that under the conditions used, >80% of the beads would bind to a single surface protein.

**Viscous Drag of Raft-associated Proteins**

Fig. 4 displays the distribution of viscous drag measure-

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**Figure 3.** The lateral viscous drag $\gamma$ (A), the lateral spring constant $k_s$ (B) and autocorrelation time $\tau_s$ (C) of a 0.2-\(\mu\)m sphere binding to a PLAP molecule plotted against time. The time trace shows three regions corresponding to the reference measurement away from the surface (I), the approach of the bead to the membrane, the diffusion near the plasma membrane (II), and after binding to the membrane protein (III). The $\gamma_s$ (B) remains unchanged during the experiment, whereas $\tau_s$ (C) changes as $\gamma$ (A).
Effects of Cholesterol Extraction

For the raft-associated proteins PLAP, YFPGLGPI, and HA, except for PLAP in untreated cells and YFPGLGPI in cholesterol-extracted cells, the distributions of single molecule viscous drag measurements have two peaks. For the transmembranous HA expressed in BHK cells, we measured $\gamma = (2.9 \pm 0.7) \times 10^{-9}$ Pa·s·m (67%) and $\gamma = (1.4 \pm 0.4) \times 10^{-9}$ Pa·s·m (33%) (n = 12, P < 0.01, Fig. 4 A). The GPI-anchored protein PLAP had $\gamma = (3.2 \pm 0.8) \times 10^{-9}$ Pa·s·m (n = 20, 100% raft-associated, Fig. 4 B), which is the summary of measurements performed on a stably expressing BHK cell line and on transiently transfected BHK cells. 76% of YFPGLGPI expressed in PtK2 cells had $\gamma = (3.0 \pm 0.9) \times 10^{-9}$ Pa·s·m, whereas 24% of YFPGLGPI had $\gamma = (0.7 \pm 0.3) \times 10^{-9}$ Pa·s·m (n = 29, P < 0.001, Fig. 4 C). Thus, for all three proteins, the majority of molecules had a viscous drag $\sim 3 \times 10^{-9}$ Pa·s·m. This value was independent of the type of membrane anchor of the protein and remained unchanged for $\sim 1$ min, which was the maximal time period continuously accessible in these experiments. For YFPGLGPI and HA, a fraction of the proteins was found to have a lower viscous drag, which was dependent on the type of membrane anchor (P < 0.1).

Effect of Cholesterol Extraction

To investigate the effect of cholesterol depletion on rafts, cholesterol was extracted from the cells by incubation with methyl-β-cyclodextrin. Extraction of cholesterol from cells results in the dissociation of proteins from rafts and the disappearance of the proteins from the floating fraction in density gradient centrifugation (Cerneus et al., 1993; Scheiffele et al., 1997; Ledesma et al., 1998). PtK2 and BHK cells were treated according to the same protocol that had been shown to extract $\sim 60\%$ of the total cellular cholesterol in BHK cells (Keller and Simons, 1998).

The distribution of viscous drag measurements of the transmembranous raft protein HA in BHK cells was shifted slightly after cholesterol depletion: 50% of HA molecules had $\gamma = (1.2 \pm 0.6) \times 10^{-9}$ Pa·s·m, whereas 50% of HA remained unchanged at $\gamma = (3.1 \pm 0.6) \times 10^{-9}$ Pa·s·m (n = 10, Fig. 4 D). The effect on the GPI-anchored proteins was more pronounced: 40% of PLAP in BHK cells had a greatly reduced viscous drag (E, n = 15, P < 0.01), 60% yielded only slightly reduced values (the multiple peaks are not statistically significant, P > 0.1). The effect of cholesterol depletion was most pronounced for YFPGLGPI. In PtK2 all YFPGLGPI had extremely reduced viscous drags (F, n = 9, P < 0.001).

Viscous Drag of Nonraft Transmembrane Proteins

To determine the viscosity of the plasma membrane outside of rafts, the local viscous drag of the nonraft transmembrane protein LYFPGT46 was measured. This construct comprised YFP as ectodomain, the transmembrane membrane anchor of the protein and remained unchanged for $\sim 1$ min, which was the maximal time period continuously accessible in these experiments. For YFPGLGPI and HA, a fraction of the proteins was found to have a lower viscous drag, which was dependent on the type of membrane anchor (P < 0.1).
domain of the LDL receptor and the cytoplasmic tail of CD46 containing a basolateral targeting signal, which, however, does not function as an endocytosis signal (Maisner et al., 1997). In PtK₂ cells $\gamma = (1.1 \pm 0.5) \times 10^{-9}$ Pa·s·m was determined for LYFPGT46 ($n = 13$, Fig. 5 A). Using a construct of the human transferrin receptor, we measured comparable viscous drags in BHK cells (data not shown).

**Effect of Cholesterol Extraction**

To study whether cholesterol extraction has an effect on the diffusive behavior of nonraft proteins, we measured the viscous drag of LYFPGT46 after depletion of cholesterol from the cells. Their viscous drag did not change, remaining at $\gamma = (1.1 \pm 0.4) \times 10^{-9}$ Pa·s·m for 100% of the molecules ($n = 8$, Fig. 5 B).

**Stability of Raft Association**

The viscous drag was measured continuously over a period of $\sim$1 min. Depending on the time that the antibody-coated bead bound to the membrane protein, this provides an observation interval for the protein diffusion of 10-50 s. During this interval, we have never observed that a raft-associated protein has dissociated from the raft ($n = 50$).

To address the question whether raft proteins would stay associated with these domains on longer time scales and would diffuse over the cell surface together with the rafts, spheres bound to GPI-anchored raft proteins were released from the laser trap after the measurement (Fig. 6 A, II). The spheres were allowed to diffuse freely for 10 min, typically for a distance of $\sim$10 µm over the cell surface. The laser trap was repositioned to trap the sphere again, and new local viscous drag measurements were performed (III). (B) The lateral viscous drag $\gamma$ of a 0.2-µm sphere bound to YFPGLGPI plotted against time. During the experiment ($t = 43-48$ s), the laser trap was used to move the sphere 1 µm laterally on the cell surface.

**Discussion**

To understand the role of rafts in cell signaling and protein sorting, knowledge about raft dynamics will be needed. Therefore, we have characterized the stability and size of rafts in cell membranes on a time scale from seconds to minutes by measuring the local viscous drag of single
membrane proteins. The size dependence of the viscous drag of a protein diffusing in a lipid bilayer was described by Saffman and Delbrück (1975). In dimyristoylphosphatidylcholine (DMPC) bilayers, their theory has been validated by Peters and Cherry (1982) who used FRAP to measure the diffusion coefficient of bacteriorhodopsin to be between $0.15 \times 10^{-8}$ cm$^2$/s and $3.4 \times 10^{-8}$ cm$^2$/s, depending on the protein/lipid ratio. In contrast, analysis of protein movement in the plasma membrane has produced various results. While assays using FRAP follow an ensemble of protein molecules, the application of SPT to plasma membrane components permits observation of movements of individual protein molecules. From many different measurements, it has become clear that the long range protein diffusion observable by video-based SPT is hindered in cell membranes by obstacles and fences, and the diffusion coefficients obtained are much slower than predicted by the Saffman-Delbrück model (Saxton, 1982, 1989; Sako and Kusumi, 1994; Kusumi and Sako, 1996). Tomishige et al. (1998) found an 8-fold increase in the diffusion of the Band3 protein in erythrocyte ghosts for the microdiffusion, when they increased the temporal resolution from 33 to 0.22 ms. In the present study, we have not only improved the temporal and spatial resolution more than 10-fold, to 18.75 μs and 5Å, but moreover, have used the laser trap to confine the motion of bead and attached protein to a small region (diam ≈ 100 nm). The latter approach provided extended observation periods of obstacle-free diffusion of single membrane proteins as verified by profiling the confining potential from the thermal position fluctuations. The obtained local diffusion coefficients of proteins in the plasma membrane in intact cells agree for the first time well with the Saffman-Delbrück model, e.g., for LYFGT46 the viscous drag measurement translates into a diffusion coefficient $D = 3.9 \times 10^{-8}$ cm$^2$/s at 36°C. The precision afforded by the excellent spatial and temporal resolution of this new method allowed us to go one step further and compare the behavior of proteins associated with lipid rafts with the one of nonraft proteins.

The distributions of the raft-associated proteins show a minority of molecules with low viscous drag, but most of the molecules have a large viscous drag. This higher value is the same for all three constructs independent of their type of membrane anchor or ectodomain. Therefore, we think that this value is correlated to the diffusion of the raft itself. However, the smaller viscous drag value measured for a few molecules depends on the type of membrane anchor of the protein. These molecules are assumed to be currently in a nonraft environment. The degree of raft association determined by our measurements (PLAP 100%, YFPGLGPI 76%) agrees with the raft association observed by the detergent extraction method: PLAP 90% (Brown and Rose, 1992), YFPGLGPI 70%, K (Keller, P., D. Toomre, J. White, and K. Simons, manuscript in preparation). After cholesterol depletion, the viscous drag values of raft proteins are reduced, with a significantly larger reduction for GPI-anchored proteins than for transmembrane proteins. The majority of molecules for each protein has viscous drag values comparable to the minority population before extraction, agreeing with the interpretation that these are molecules without raft environment. The effect of cholesterol depletion is more pronounced in PtK$_2$ cells than in BHK cells. Possibly, PtK$_2$ cells are more susceptible to cholesterol extraction because of their large plasma membrane area per cell. Cholesterol depletion had no effect on the diffusion of nonraft proteins. The distribution of viscous drags measured for LYFGT46 shows only one peak that does not shift upon cholesterol depletion.

Fig. 7 summarizes the means of the majority of molecules for each raft and nonraft protein studied. Important to note is that two GPI-anchored proteins with different ectodomains (PLAP and YFP) yield similar results, whereas GPI-anchored YFP diffuses differently from transmembrane-anchored YFP, showing that the local viscous drag is determined by the membrane domain of the protein. Why do our data indicate that rafts diffuse as a small unit? First, some of the raft proteins (PLAP and YFPGLGPI) are attached to the membrane by a lipid anchor that just penetrates the outer leaflet of the bilayer, whereas HA is a transmembrane protein traversing both leaflets. Nevertheless, all three proteins diffuse as if connected to a stable membrane domain of similar size with a viscous drag significantly larger than that of the nonraft transmembrane protein LYFGT46 ($P < 0.01$). Second, after depleting cholesterol from the cell membrane, the viscous drag of raft proteins is greatly reduced: the transmembrane protein HA shows a viscous drag similar to the nonraft transmembrane protein LYFGT46, which is equivalent to a threefold reduction. However, the viscous drag of the GPI-anchored proteins is reduced about five-fold and is lower than that exhibited by the nonraft transmembrane proteins ($P < 0.1$). After cholesterol depletion, GPI-anchored raft-proteins diffuse faster than transmembrane raft proteins, even though in untreated cells they behave the same. The cholesterol depletion has negligible effect on the nonraft transmembrane proteins, indicating that the viscosity of the cell membrane outside of the raft domains remains unchanged. We interpret this difference in behavior as dissociation of the raft assembly or as segregation of the proteins from the lipid raft. A tetra cholesterol depletion, the raft proteins diffuse like nonraft proteins.
with the diffusion coefficient becoming dependent on the type of the membrane anchor. An alternative explanation for this reduction of the viscous drag would be that the diffusion observed was actually within a large raft, and that the viscosity of the raft domain was changed by cholesterol extraction. However, protein diffusion in a lipid bilayer depends approximately linearly on the membrane viscosity. Therefore, the latter interpretation is not compatible with the observation that cholesterol depletion reduces the viscous drags for GPI-anchored proteins fivefold, whereas the one of transmembrane raft proteins only threefold. The larger viscous drag of raft compared with nonraft proteins indicates that they are anchored to a membrane patch with a larger diameter than the transmembrane domain of the LDL receptor. Third, the absence of any change after translocating spheres bound to a raft protein laterally over the cell shows that the lipid rafts are moved together with the protein. The other explanation for these data would be that rafts are larger than 1 μm, which would contradict the observations published previously using fluorescence methods (Harder et al., 1998; V arma and M ayor, 1998). The local viscous drag on all raft-associated proteins remained unchanged during the entire observation interval of 50 s. Thus, raft proteins stay raft-associated for at least 1 min. Experiments observing the raft protein diffusion intermittently over extended time intervals indicate that raft proteins can diffuse over the cell surface while being raft-associated for up to 10 min. Taken together, these points let us conclude that the raft association of proteins is quite stable on the time scale of at least 1 min, and that rafts diffuse as one entity, which means that the diffusion coefficient measured is dominated by the motion of the entire raft.

Consequences for the Raft Model

Our data agree with the model suggested by Simons and Ikonen (1997) of cholesterol-stabilized rafts associating with proteins in the native state of the cell membrane. Other models have suggested that proteins associate with rafts only after clustering, which might induce the separation of lipids and cause raft formation. Alternatively, the affinity of proteins associating with rafts would be so low that only clustering would stabilize the association. These latter models do not fit our observations, that proteins associate with rafts without clustering and that raft proteins are moved together with the protein. The other explanation for these data would be that rafts are larger than 1 μm, which would contradict the observations published previously using fluorescence methods (Harder et al., 1998; V arma and M ayor, 1998). The local viscous drag on all raft-associated proteins remained unchanged during the entire observation interval of 50 s. Thus, raft proteins stay raft-associated for at least 1 min. Experiments observing the raft protein diffusion intermittently over extended time intervals indicate that raft proteins can diffuse over the cell surface while being raft-associated for up to 10 min. Taken together, these points let us conclude that the raft association of proteins is quite stable on the time scale of at least 1 min, and that rafts diffuse as one entity, which means that the diffusion coefficient measured is dominated by the motion of the entire raft.

**Estimation of the Raft Size**

If the raft diffuses as a stable structure as our results suggest, then the size of the raft can be estimated from the local viscous drag measurements. The Saffman-Delbrück model provides the simplest description of the viscous drag of a protein with radius \( r \) diffusing in a lipid bilayer with the membrane viscosity \( \eta_m \) and thickness \( h \) (Saffman and Delbrück, 1975; see Materials and Methods). It was shown to be a valid model for protein diffusion in lipid bilayers (Peters and Cherry, 1982; V az et al. 1984), and bilayers containing other mobile proteins can be described by an increased effective viscosity (Cherry and Godfrey, 1981; Saxton, 1987).

Therefore, the Saffman-Delbrück model was applied to estimate the size of rafts. To obtain the membrane viscosity, we used the result of the nonraft protein LY FPG T 46, which has a single transmembrane domain. Recently, E s-kandari et al. (1998) have measured the transmembrane area of several membrane proteins and estimated that a single transmembrane helix occupies 1.4 nm². Membrane thickness \( h \) and viscosity \( \eta_m \) are coupled as \( h \cdot \eta_m \) in the Saffman-Delbrück relation. Our data for the nonraft transmembrane protein are best fit by \( h \cdot \eta_m = (6.3 \pm 1.2) \times 10^{-10} \) Pa·s·m. Therefore, the viscosity of a membrane with \( h = 5 \) nm would be \( \eta_m = 0.13 \pm 0.03 \) Pa·s. This value agrees well with the viscosity of a D M PC bilayer (\( \eta_m = 0.11 \) Pa·s; Peters and Cherry, 1982). To address possible influences from an extracellular coating, we performed some measurements with a construct of the human transferrin receptor with a modified cytoplasmic tail (data not shown). The derived membrane viscosity agrees well with the value for \( h \cdot \eta_m \) obtained here. A iso, the fact that HA in cholesterol-depleted cells diffused just like LY FPG T 46 dismisses any effect from an extracellular coating. In addition, an extracellular coating would not influence the diffusion in a protein-radius-dependent manner. Under these conditions, the estimated radius of the rafts is \( r = 26 \pm 13 \) nm. This estimate is also obtained with the formula for protein diffusion of H ughes et al. (1981, 1982), which includes coupling to the surrounding medium, because the viscosity of water and the cytoplasm are low compared with the membrane viscosity (Swam-nathan et al., 1997). The broad distribution of viscous drag values observed for the rafts studied indicates that even one raft type in one cell might have a distribution of sizes, and that the size of a single raft might be dynamic.
Nevertheless, transferring these values to other cell types and non-GPI–anchored proteins might not always be possible since raft size and stability are most likely dependent on the lipid and protein constituents. Indeed, recent data suggest that rafts in the apical membrane of M D CK behave differently from rafts in fibroblasts (Verkade et al., 2000).

A corolling to our study, rafts in the plasma membrane of fibroblast-like cells diffuse as a rather stable platform with an average area of 2,100 nm². The size estimate allows an assessment of the maximal contents of one raft. If these were composed purely of lipid molecules, having a radius comparable to phosphoethanolamine (r = 0.44 nm), one raft would contain almost 3,500 lipid molecules. How many proteins a raft contains depends on how densely packed the proteins would be. If they were as densely packed as rhodopsin molecules are in frog rods (B lasie and Worthington, 1969), or as the spikes in the envelope of Semliki Forest virus (Cheng et al., 1995), a raft would contain 55–65 proteins, respectively. Clearly, the packing density in a lipid raft in a mammalian plasma membrane would be lower. Indeed, recent results using fluorescent folate to measure the fluorescent resonance energy transfer between folates bound to folate receptors indicate that few of the receptors are clustered and many more are monomers. Further, these clusters consist of multiple GPI–anchored proteins if they exist in the same cell and are disrupted if cholesterol levels are depleted (V arma, R., and S. Mayor, personal communication).

Implications for Raft Function

The consequences of a small raft size and stable association is that proteins within rafts would be restricted in their interactions with other proteins. To use rafts as platforms in membrane trafficking (e.g., in apical transport from the Golgi complex) would imply that these rafts have to be clustered together by an apical sorting machinery to form a container comprising several rafts. A iso, the formation of caveolae at the plasma membrane would involve the invagination of several rafts to form one caveola. Often caveolae are found like grapes in large clusters, thus ten caveolae are found like grapes in large clusters, thus forming a container comprising several rafts. Also, the formation of caveolae at the plasma membrane would involve the invagination of several rafts to form one caveola. Often caveolae are found like grapes in large clusters, thus forming a container comprising several rafts. Also, the formation of caveolae at the plasma membrane would involve the invagination of several rafts to form one caveola. Often caveolae are found like grapes in large clusters, thus forming a container comprising several rafts. Also, the formation of caveolae at the plasma membrane would involve the invagination of several rafts to form one caveola. Often caveolae are found like grapes in large clusters, thus forming a container comprising several rafts. Also, the formation of caveolae at the plasma membrane would involve the invagination of several rafts to form one caveola. Often caveolae are found like grapes in large clusters, thus forming a container comprising several rafts. Also, the formation of caveolae at the plasma membrane would involve the invagination of several rafts to form one caveola. Often caveolae are found like grapes in large clusters, thus forming a container comprising several rafts. Also, the formation of caveolae at the plasma membrane would involve the invagination of several rafts to form one caveola.
Peters, R., and R.J. Cherry. 1982. Lateral and rotational diffusion of bacterio-
rhodopsin in lipid bilayers: experimental test of the Saffman-Delbrück equa-
Pralle, A., E.-L. Florin, E.H.K. Stelzer, and J.H.K. Hörber. 1998. Local viscos-
hree-dimensional high-resolution particle tracking for optical tweezers by
Saffman, P.G., and M. Delbrück. 1975. Brownian motion in biological mem-
membrane for receptor movements as revealed by a nanometer-level mo-
Sako, Y., and A. Kusumi. 1995. Barriers for lateral diffusion of transferrin re-
ceptors in the plasma membrane as characterized by receptor dragging by la-
Saxton, M.J. 1982. Lateral diffusion in an archipelago: effects of impermeable
Saxton, M.J. 1987. Lateral diffusion in an archipelago. The effect of mobile ob-
Saxton, M.J. 1989. Lateral diffusion in an archipelago. Distance dependence of
hemagglutinin with sphingolipid-cholesterol membrane domains via its
387:569-572.
Swaminathan, R., C.P. Hoang, and A.S. Verkman. 1997. Photobleaching recov-
ery and anisotropy decay of green fluorescent protein GFP-565T in solution
and cells: cytoplasmic viscosity probed by green fluorescent protein transla-
Teuchert, M., A. Maisner, and G. Herrler. 1999. Importance of the carboxyl-
terminal FTSL motif of membrane cofactor protein for basolateral sorting
and endocytosis. Positive and negative modulation by signals inside and out-
Tomishige, M., Y. Sako, and A. Kusumi. 1998. Regulation mechanism of the
lateral diffusion of band 3 in erythrocyte membranes by the membrane skel-
Varma, R., and S. Mayor. 1998. GPI-anchored proteins are organized in submi-
Yamamoto, T., C.G. Davis, M.S. Brown, W.J. Schneider, M.L. Casey, J.L.
Goldstein, and D.W. Russell. 1984. The human LDL receptor: a cysteine-
rich protein with multiple Alu sequences in its mRNA. Cell. 39:27-38.