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Institute of Computational Science

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Diffusion in the Endoplasmic Reticulum

Theoretical Foundation, Computer Simulations, Models

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Abstract

Diffusion processes in the lumen of the endoplasmic reticulum of tissue derived cells are investigated using a computational approach based on the method of particle strength exchange. 3D reconstructions of the shape of the endoplasmic reticulum are made from micrograph sections and are then used to conduct direct numerical simulations of fluorescence recovery experiments. The influences of the local geometry (i.e. the restriction of diffusion to it) on the time behavior of diffusion are discussed. Using concepts from fractal geometry, these influences are described with a single parameter that is proven to be sufficient. Based on this theory and physical principles, novel mathematical models for fluorescence recovery curves are developed and validated. They are shown to approximate the recovery curve at least one order of magnitude better than current models on the same geometry.

Zusammenfassung

Diffusionsprozesse im Lumen des Endoplasmatischen Retikulums lebender Zellen werden mittels Rechnersimulationen untersucht, wobei die "particle strength exchange"-Methode verwendet wird. Basierend auf mikroskopischen Schnittbildern werden 3D-Rekonstruktionen der Geometrie des Endoplasmatischen Retikulums gemacht, welche dann benutzt werden, um direkte numerische Simulationen von "fluorescence recovery"-Experimenten zu machen. Die Einflüsse der lokalen Geometrie (d.h. der Einschränkung der Diffusion auf diese) auf das zeitliche Verhalten der Diffusion werden diskutiert. Unter Nutzung von Konzepten aus der fraktalen Geometrie werden diese Einflüsse dann mittels eines einzigen Parameters beschrieben, welcher als dafür hinreichend bewiesen wird. Basierend auf dieser Theorie und physikalischen Grundgesetzen werden neue mathematische Modelle für "fluorescence recovery"-Kurven hergeleitet und validiert. Es wird gezeigt, dass diese die Kurven für die selbe Geometrie mindestens eine Grössenordnung besser wiedergeben als die heute verwendeten Modelle.

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to my Mother

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List of symbols

a	First edge vector of a triangle, $a = P_2 - P_1$ / a coefficient / a point
(a,c)	Lower-left corner coordinates of bleached box
\hat{a}	Least squares estimate of the slope
a_0	A coefficient
A	A set / integration constant/ bleached area
b	Second edge vector of a triangle, $b = P_3 - P_1$ / a point
(b,d)	Upper-right corner coordinates of bleached box
b_x, b_y, b_z	Size of a bin or cell in all three directions
\hat{b}	Least squares estimate of the y -axis intersection
B	The bleached box $/$ a set $/$ integration constant
B: :	Bin with indices (i, j)
c	Concentration / centroid of a triangle / a constant
с С-	Concentration at centroid of a grid cell
CV	Centroid of a voyel
C	Cell with indices (i, j, k)
$C_{i,j,k}$ C(E)	Set of all continuous functions on E
D	Diffusion constant / a bounded open region
$D_{\rm D}$	Box counting dimension
D_B	Hausdorff's dimension
D_H	Minkovski dimension
	Components of the anisotropic diffusion tensor
$D_{xx}, D_{xy}, \ldots, D_{zz}$	Spectral dimension
d_s	Dimension of the walk
e(n)	Parallel speed-up efficiency on <i>n</i> processors
E(n)	A set geometrical object / quadratic error
	Pre-fractal Euclidean approximations of order k
E_k	BMS error
E _{RMS}	δ -neighborhood of E δ -narallel body to E
\mathbf{E}_{0}	Expectation value of a random variable
f	A function
J F	A function / FRAP value
F _o	Pre-bleach FRAP value
F_	FRAP value just after bleaching
F _{ac}	Steady-state FBAP value for $t \to \infty$
1 f_	Negative part of a function
f_{\perp}	Positive part of a function
a	A function
G	Green's function / a function / number of grid points
\hat{h}	Grid spacing / inter-particle spacing
$H_{i,i,k}$	Head of chain list for cell (i, j, k)
\mathcal{H}^{s}	The <i>s</i> -dimensional Hausdorff measure
i, j, k	Indices, index counters / Integer grid coordinates in simulation
I	Identity matrix
k	A constant / particle index
K	Heat Kernel (time dependent Green's function) / convolution kernel
l	Particle index
l	Linked list / load distribution coefficient
L	Edge length of cubic domain
$L_{i,i}$	Triangle list associated to bin (i, j)
$L_{i,j,k}$	Triangle list associated to cell (i, j, k)
\mathcal{L}^{n}	The n -dimensional Lebesgue measure
m	Index / a matrix / an element / mass
m_t	Total mass in the domain
M	Number of triangles in surface / a system matrix
	number of terms in a finite series expansion

\mathcal{M}	Family of μ -measurable sets
n	Topological dimension of a space / outer normal vector / index
	number of measurement points / number of processors
N	Total number of particles in a simulation
N_B	Number of particles inside the bleached box, area
N_i	Number of intersections
N_{Iter}	Number of iteration steps
N_{δ}	Measured length in multiples of δ
Na	Number of grid points inside the bleached box
$N(\lambda)$	Eigenvalue distribution function
N_{π} , N_{η} , N_{τ}	Number of grid points, bins or cells in each direction
Nc_{π} , Nc_{π} , Nc_{π}	Number of mesh cubes in each direction
\mathcal{N}	Gaussian random variable (normal distribution)
\mathbb{N}^n	<i>n</i> -dimensional space of integer numbers
IN ⁺	Set of all non negative integer numbers
overlap	Particle overlap parameter
overtap	Particle overlap parameter
p	A maritime main lie function
$p(\cdot)$	
p_t	Transition density at time t
P	A point / a particle / probability of an event
P_1, P_2, P_3	Vertex points of a triangle
q	Particle index / source term
Q	A point
Q_{ϵ}	Anisotropic integral operator for diffusion
$\overline{Q}_{\epsilon}^{n}$	Discrete particle approximation to Q_{ϵ}
r	Ratio of anisotropy (ratio of diffusion resistances in x and y)
	Radial spherical coordinate / distance / order of a particle method
r_c	Cut-off radius
R	Radius of a tube / mobile fraction
$R_n^x(r), R_n^y(r)$	Effective diffusion resistances in x and y
\mathbb{R}^n	Euclidean space of dimension n
\mathbb{R}^+	Set of all positive real numbers
s	Dimension (of a Hausdorff measure) / random walk step vector
S+	Time stretch factor
S(n)	Parallel speed-up on n processors
t (10)	Time
$t_{T}(n)$	Execution time on n processors
$t_E(n)$	Time sport on processor <i>i</i>
t_i	Half time of fluorescence recovery
$t_{1/2}$	Final time of simulation
r_{final}	Number of time stops / time constant
	Telerence
10L	Interance
И	Uniformity distributed random variable
u	Unknown variable of a differential equation / velocity vector
v v	A vector / volume of a particle
V	Volume
V_k	Vertex with index k
$V_{i,j,k}$	Voxel with indices (i, j, k)
w	Vertex index / quadrature weight
x	Location vector / position / a point / argument to a function
\overline{x}	Mean displacement
ξ	Location vector / position
x_p	Position of particle p
X	A set
$X_k(t)$	Random walk with step length 2^{-k}

Coordinates in a cartesian system Coordinates in a cartesian system

Brownian motion

X(t)x, y, z ξ, η, ζ

y	Position / Argument to a function
Y	Intersection point / vector of measured values
α	Order of continuity / model parameter / line parameter
α_k	Time between subsequent steps in a random walk of order k
β	Model parameter / line parameter
γ	Order of continuity
δ	Size of length scale / uniform grid spacing
δ_{ij}	The Kronecker delta symbol
$\delta(\cdot)$	Dirac's delta distribution
δt	Time step size
$\delta x, \delta y, \delta z$	Cartesian grid spacings in each direction
Δ	Difference / determinant
Δ_k	Triangle k
Δz	Distance between subsequent slices
ϵ	Kernel core size for particle method
ε	Residual / error
η	Number of voxels in a voxel representation / kernel function
η_{ϵ}	Regularized kernel function for a particle method
$\eta_{\alpha}, \eta_{\beta}$	Learning rates for gradient descent
φ.	Eigenfunction / azimuthal angle / angular coordinate on tube
ϕ_0	Volume-filling coefficient
Φ	Regression matrix
	Spherically symmetric matrix cut-off function
$\frac{\tau}{\eta/2}\epsilon$	Regularized matrix cut-off function
ү К	Dimensionless constant
λ	Figenvalue / line parameter / ratio of a similarity transformation
	A measure / Micro
μ ν	Restriction of a measure / a distribution
π	Batio of circumference to diameter of a circle $(= 3.14159265359)$
σ	Standard deviation / voxel edge length / source term / anisotropic kernel
σ.	Regularized anisotropic kernel
τ	Time / time constant
29	Polar angle for spherical coordinates
Â	Vector of unknown parameters
0	Scalar cut-off function
Â	Least squares estimation of unknown parameters
Θ_{LS}	Destination of unknown parameters
ζ	Particle moliner function / axial cool difface of a tube
Se	Indicator function of set 4
χ_A	Darticle strength (in general)
ω	Particle strength (in general)
22	Triangulated gunface, i.e. its set of triangles
2	Dantial during time / hour dama of damain
	Partial derivative / boundary of domain
$\frac{V}{\nabla^2}$	Nabla Operator
$\frac{V}{\nabla^2}$	Carting and internal composition to the London composition
∇_{ϵ} ∇^2	Continuous integral approximation to the Laplace operator
$V_{\epsilon,h}$	Discrete sum approximation to the Laplace operator
∞	
• 	Absolute value / diameter of a set
$\frac{\ \cdot\ }{\langle \cdot \rangle} \infty$	Infinity form of a function
(\cdot)	Upper dimension
$\underbrace{(\cdot)}$	Lower dimension
(\cdot)	Finite sum approximation of an infinite series
$(\cdot)^{ op}$	Matrix transpose
$(\cdot)^h$	Particle approximation with inter-particle spacing h
$(\cdot)^*$	True value (as opposed to model value)
$(\cdot)'$	Derivative with respect to the function's only variable
$(\cdot)_0$	Initial value, property at time zero

$(\cdot)_{\infty}$	Steady-state value for $t \to \infty$
$(\cdot)_{x}, (\cdot)_{1}$	Cartesian x component or value in x direction
$(\cdot)_{\alpha}$ $(\cdot)_{\alpha}$	Cartesian v component or value in v direction
()y, ()2	Cartesian z component or value in z direction
$(\cdot)_{z}, (\cdot)_{3}$	Cartesian z component or value in z direction
$(\cdot)_l$	Lower bound of a variable
$(\cdot)_h$	Upper bound of a variable
(·)min	Minimum value of a variable
())	Maximum value of a variable
$(\cdot)_{max}$	
$(\cdot)_{opt}$	Optimum value of a variable
$(\cdot)_{mirror}$	Boundary mirror image
$(\cdot)_n$	At time step n
$(\cdot)_k$	At iteration step k
$O(\cdot)$	The Landau symbol
$\mathcal{O}(\cdot)$	On emotion accent at the ender of
$\overrightarrow{U(\cdot)}$	Operation count at the order of
AB	Vector pointing from A to B
×	Cartesian product / vector cross product
\forall	For all
()I	Foculty of a number
(•):	Faculty of a number
÷	Must be equal (requirement)
\times	Asymptotically equal
E	In a set, member of a set
d	Not in a set
≻ \	Sot minus
\ Ø	Empty act
ψ	
$spt(\cdot)$	Support of a measure
$int(\cdot)$	Interior of a closed set
$\operatorname{Vol}^n(\cdot)$	n-dimensional volume of an object or set
[a,b]	Interval between a and b including its boundaries
(a,b)	Interval between a and b excluding its boundaries
P(Sub-vector with components n to m of vector P
$\binom{n:m}{n}$	Multiset union of <i>n</i> sets
$\sum_{n=1}^{N}$	
$\sum_{n=1}^{j=1}$	Sum of <i>n</i> terms
$\prod_{i=1}^{n}$	Product of <i>n</i> terms
inf	Infimum of a set (multidimensional minimum)
\sup	Supremum of a set (multidimensional maximum)
min	Minimum of a set
max	Maximum of a set
lim	Limit
log	Natural loganithm (i.e. loganithm to the base o)
10g	(1.e. logarithm to the base e)
$\operatorname{celling}(x)$	Smallest integer larger or equal x
floor(x)	Largest integer less or equal x
int(x)	Integer part truncation of x
$Tr(\cdot)$	Trace of a matrix
$\#\{\cdot\}$	Counting measure of a set (number of elements in the set)
\Rightarrow	Logical implication
\Leftrightarrow	Only if (iff)
\mapsto	Maps to (mapping or transformation)
	To towards
\rightarrow	Non datama annum ant
<.>	Mandatory argument
L.J	Optional argument

Further symbols may be defined in the text as they are encountered. Name conflicts may occur. However, in such cases the meaning of a symbol is always clearly defined in the text. Vectors are not printed in bold or otherwise marked as their occurrence depends on the dimension of the underlying space.

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Acronyms

2D	two-dimensional
3D	three-dimensional
ASCII	American Standard Code for Information Interchange
С	The C programming language
CD	Compact Disc
CD-ROM	Compact Disc Read-Only Memory
CFL	Courant-Friedrichs-Levy (condition)
DNA	Desoxy-Ribonucleic Acid
DVI	Device Independent Interface file
EPS	Encapsulated Post-Script
\mathbf{ER}	Endoplasmic Reticulum (an organelle of a biological cell)
ETH	Eidgenössische Technische Hochschule (Swiss Federal Institute of Technology)
FD	Finite Difference
\mathbf{FFT}	Fast Fourier Transformation
FRAP	Fluorescence Recovery After Photobleaching
GB	Giga Byte
GFLOPS	Billion Floating Point Operations Per Second
GFP	Green Fluorescent Protein
ICoS	Institute of Computational Science
IRIX	The Silicon Graphics Operating System
kDa	Kilo Dalton (a protein molecular mass unit)
MB	Mega Byte
MBit	Mega bit
MFLOPS	Million Floating Point Operations Per Second
MIPS	Million Instructions Per Second
MOPS	Million Operations Per Second
MPI	Message Passing Interface
MPEG	Motion Picture Experts Group
ODE	Ordinary Differential Equation
OpenDX	A 3D data visualization tool
PDE	Partial differential equation
PDF	Portable Document Format
PS	Post-Script
PSE	Particle Strength Exchange
PSE3D	Three-dimensional particle strength exchange method
RMS	Root Mean Square
SGI	Silicon Graphics, Inc.
ssGFP-KDEL	ER resident signal sequence GFP
VERO	Cell from green monkey intestine epithelium

Names of simulation runs

bip2	erp574_1
clx	erp574_2
erp57	erp5722
erp572	erp5723
erp573_1	erp5724
erp572_2	8s
erp573_3	8.2

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Foreword

This report describes the work and presents the results of a diploma project that has been hosted by the Institute of Computational Science (ICoS) at the Swiss Federal Institute of Technology (ETH) in Zürich, Switzerland. A close collaboration with the Institute of Biochemistry at ETH enabled the interdisciplinary and applicationrelated character of the project. The work is a direct continuation of the semester project [Sbalzarini (2001)], to which frequent references will be made throughout the text. The reader will be assumed to be familiar with the basic concepts, ideas and results of [Sbalzarini (2001)]. Namely preliminary knowledge about the endoplasmic reticulum (ER), its function and shape, the concept of fluorescence recovery after photobleaching (FRAP), the green fluorescent protein (GFP) and the notions of random walk, Brownian motion and diffusion is expected. Readers not familiar with any of these subjects are advised to read chapters 2 (for the ER in general as well as its function and shape), 3 (for the concept of FRAP analysis and GFP) and 4 (for diffusion and random walk) of [Sbalzarini (2001)] as an introduction to the present work. It is electronically available from the e-collection of ETH library.¹

The main purpose of this project is to further pursue the investigation of diffusion processes inside the endoplasmic reticulum of tissue derived cells. Particularly, the following goals are aimed at:

- Generation of 3D reconstructions of real ER geometries and representation as triangulated surfaces in the computer
- Development of a 3D geometry preprocessor for use with existing simulation codes
- Mathematical treatment of the theory of diffusion on fractal sets
- Identification of the relevant statistical geometry parameter and proof of its sufficiency
- Investigation of the usefulness of fractal models to describe protein diffusion in the ER
- Measurement of the fractal dimension of the ER surface in space and comparison to the findings in [Sbalzarini (2001)]
- Change of the simulation technique employed from random walk to the superior method of particle strength exchange
- Direct numerical simulations of FRAP experiments in full 3D reconstructions of real ER structures
- Development of a novel FRAP data model to fit simulation and experiments
- Comparison of simulations and models to experiments

¹http://e-collection.ethbib.ethz.ch/browse/alph/s.html

Due to the interdisciplinary nature of the project, this report is targeted at both molecular biologists and applied mathematicians. In an effort to make it (at least partially) comprehensive to both groups, certain repetitions and somewhat lengthy-seeming explanations could not be avoided in the text. To keep it as clean and general as possible, vectors are not printed in bold or otherwise marked as their occurrence depends on the dimension of the underlying space. It is always clear from the context which variables are vectors (e.g. position in spaces of more than one dimension) and which are scalar (e.g. time).

Chapter 1 contains an overview of the current state and a survey of relevant literature. Moreover, it gives a concise introduction to the strategy that will be followed for the rest of the work. Chapter 2 starts with a description of the process of 3D reconstruction of ER shapes from a stack of parallel microscopical sections. Chapter 3 describes the algorithms that have been implemented in order to check the resulting triangulation for topological and syntactical validity as defined in [Sbalzarini (2001)].

Chapter 4 represents the first piece of the theoretical part of this work. It will introduce and summarize some basic theoretical concepts along with their notation and terminology for later use. Although all of this theory is commonly available from literature, it is concisely summarized here in a single place using a notation that is consistent with the rest of the work. However, readers are free to skip the chapter if familiar with the notions presented therein or to consult their own books of choice.

In chapter 5 the concepts previously introduced will be applied to measuring the fractal dimension of the ER surface in 3D space. As every work should be built on a sound theoretical foundation, chapter 6 will be concerned with the theory of diffusion on fractal sets. This chapter relies on the information presented in chapter 4 in order to identify the geometrical parameter that is relevant to the time behavior of diffusion and prove its sufficiency for two large classes of fractals.

This concludes the theoretical part of the report. Chapter 7 will describe the simulation methods that will be used to simulate protein diffusion in reconstructed ER shapes. Particularly, the method of particle strength exchange (PSE) and its anisotropic extension are presented. However, no anisotropic runs will be conducted in this project since biologists do not yet know anything about the anisotropy of the medium filling the ER lumen. However, all simulation algorithms are set up to handle anisotropic diffusion in order to be prepared for future work. Due to its many favorable properties, the PSE method will replace the random walk code developed in [Sbalzarini (2001)] for the rest of this project. Before applying it to real ER problems, it is validated on a test case against the analytic solution, a finite difference code and the random walk simulation in chapter 8.

Chapter 9 contains results from the actual FRAP simulations in 3D reconstructed ER samples from micrographs. Moreover, the influence of position and size of the bleached spot is investigated. In order to validate the simulations, their results are compared to experiments in chapter 10 and it is demonstrated how to obtain numerical values for the diffusion coefficient out of such simulations. However, it would be much too complex to run a full direct numerical simulation for every experiment one wishes to evaluate quantitatively. Therefore, chapter 11 starts looking into new mathematical models to describe FRAP curves with a small number of parameters. In chapter 12 the new models derived are then compared to existing ones for diffusion in a cubic box as well as diffusion in reconstructed ER samples and experimental data.

In order to allow easy re-use of the simulation programs, result data sets and resources of this work, all files are contained on two CD-ROMs that accompany this report. Besides the complete program source codes, the CDs also contain the text files of this document (IAT_EX source, EPS and PS figures, xfig drawings and all pixel images) as well as complete PostScript and PDF versions of it. Moreover, color images, movies of certain runs, all simulation and experimental data as well as the PowerPoint file of the final presentation can be found on them. Appendix A contains the complete table of contents.

Appendix B lists the numerical values of the model parameters found by fitting different FRAP data models to different ER samples. The Fortran source codes of all the programs entirely developed in this work are given in appendix C. It also contains basic information about the program structures and their usage. Since an existing code has been adapted and extended for the PSE simulations, its source code will not be given in the appendix but it is contained on CD #1 nevertheless.

Each algorithm presented in this report contains marginal numbers referring to the corresponding statement labels in the Fortran routine that implements it. This allows easy understanding of the Fortran programs and immediately links the algorithms to their respective implementation for convenient reference and easy reproduction. To keep the notation consistent, all algorithms and formulae in this report use Fortran index numbering starting at 1 rather than 0.

Acknowledgments

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Zürich, February 2002

Chapter 1 Introduction

The investigation of diffusion in complex-shaped organelles has become an important area of research in recent biology and fluorescence recovery analysis is the tool of choice (see [Lippincott-Schwartz, Snapp & Kenworthy (2001)] for a comprehensive review of the methods available). While early work applied 2D and even 1D models of homogeneous and isotropic diffusion to the problem of fluorescence recovery after photobleaching (FRAP), doubts about their validity and a potential influence of the underlying geometrical structure of the ER arose soon (e.g. stated in [Olveczky & Verkman (1998)] or [Ellenberg et al. (1997)]) and biologists started looking into mathematical models and computer simulations to capture the effects (see e.g. [Olveczky & Verkman (1998)], [Dayel, Hom & Verkman (1999)] or [Siggia, Lippincott-Schwartz & Bekiranov (2000)]). However, the problem has been taken up in applied physics long before. [Axelrod et al. (1976)] started developing models for fluorescence recovery curves as early as 1976 and their publication might have fallen into oblivion by now. In lack of simulation algorithms and computing power however, they restricted themselves to two dimensional problems and simple geometries whereas the present work will deal with 3D problems in real ER geometries. Nevertheless, they discovered some important basics and were able to match experimental data with high accuracy. By application of single particle tracking methods, [Kusumi, Sako & Yamamoto (1993)] and recently [Dietrich et al. (2002)] proposed an alternative to fluorescence recovery experiments if small statistical samples are sufficient. A comparative summary of the methods available can for example be found in [Cheezum, Walker & Guilford (2001)].

In order to be able to investigate the influence of geometrical factors on the time behavior of diffusion (i.e. the eigenvalues of the Laplacian) of molecules inside the lumen of the endoplasmic reticulum¹, several building blocks from different scientific disciplines are needed:

- Cell biology
- Biochemistry and biotechnology
- Physics of transport phenomena
- Numerics and applied mathematics
- Computer simulation techniques
- Theoretical mathematics

¹Readers not familiar with the endoplasmic reticulum or diffusion are advised to read chapters 2 through 4 of [Sbalzarini (2001)] or the corresponding chapters of [Alberts et al. (1997)]

- Fractal and Euclidean geometry
- Image processing and reconstruction methods
- Fluorescence microscopy
- Optimization algorithms

Each of them is a well-known subject on its own. For an introduction to cell biology and biochemistry, see for example [Alberts et al. (1997)]. The numerics and computer simulation methods that will be applied are for example described in [Hockney & Eastwood (1988)] or [Cottet & Koumoutsakos (2000)] as well as in [Degond & Mas-Gallic (1989a)] and [Degond & Mas-Gallic (1989b)]. Alternative methods can for example be found in [McCorquodale, Colella & Johansen (2001)] and [Beaudoin, Huberson & Rivoalen (2001)]. The field of fractal geometry and the theory of partial differential equations (PDE) involving it is subject to extensive and ongoing research in theoretical mathematics. A good introduction to fractal geometry in general can be found in [Mandelbrot (1982)]. More in-detail and mathematically profound information are given in [Falconer (1985)], [Falconer (1990)] and [Falconer (1997)], which also contains some theory about PDEs on fractal domains. A concise summary of theorems and definitions relevant for fractal geometry is given in [Monks (2001)]. In addition, several research papers contain related information. They will be cited in the text where relevant. Information about image processing and 3D reconstruction from parallel sections are for example contained in [Baldock & Graham (2000)].

Despite all those efforts in the different fields, little work has been done in trying to combine them. The three most noticeable projects are the very recent one on cellular fluid mechanics by [Kamm (2002)], the application of fractal theory to image compression ([Mitra, Murthy, Kundu & Bhattacharya (2001)]) or the first – however somewhat aimless – application of fractal geometry to biological cell image analysis by [Smith, Marks, Lang, Sheriff & Neal (1989)]. The present work aims at combining all of the above areas in order to find a better understanding of molecular diffusion in the endoplasmic reticulum which hopefully will lead to novel FRAP data models that can be used to determine diffusion constants from experimental data.

Figure 1.1 shows the global picture of how these different techniques and areas of science will be combined. The corresponding chapter numbers of where they are treated in this report are printed in *italics*. Boxes with black print symbolize steps of the present work whereas rounded boxes with blue print signify external inputs to the process. The final goal is highlighted using red letters. The left and bottom parts only need to be done once, the right top-down line describes the repeated application in daily laboratory work.

The starting point for all mathematical modeling are stacks of micrographic slices of stained ERs. These images are then used to perform a 3D geometry reconstruction in order to obtain a triangulated surface description (see [Sbalzarini (2001)]) of the real-world ER shape in the computer (see chapter 2). This triangulated representation then needs to be checked for different criteria of validity in order to make sure that it actually represents a valid ER structure (i.e. its interior needs to be a connected closed subset of \mathbb{R}^3). The topological criteria and algorithms used are described in chapter 3. The final and valid triangulated geometry can then be used as a computational domain for subsequent numerical simulations of diffusion. The simulation techniques employed and their validations are contained in chapters 7 and 8. Choosing appropriate initial conditions then allows realistic simulations of FRAP experiments. Since for the simulation runs, the corresponding diffusion



Figure 1.1: General research strategy outline

constant is known (it has to be passed to the simulation algorithm as a parameter), these FRAP curves serve as reference curves with no unknown properties. Fitting these curves to experimental data *in the same ER geometry* would in principle allow to directly determine the unknown real-world diffusion constant (dashed arrow). Since it would be too complicated to always take a stack of micrographs and run a full-fledged computer simulation for each cell one wishes to make a fluorescence recovery experiment, data models are looked for.

The starting point to develop such models is the theoretical foundation given by mathematical set theory, the physics of transport phenomena and fractal geometry, the most important concepts of which are summarized in chapter 4. Certain knowledge about the solution of the diffusion equation on fractal domains (cf. chapter 6) and the fractal properties of the specific ER geometries at hand (chapter 5) then allow to infer novel, physically motivated FRAP data models (see chapter 11). These models can be fitted to the simulated reference FRAP curves in order to get numerical values for their parameters. As both the model parameters and the underlying diffusion constant are known, it is in principle possible to find the physically given functional relationship f between them (section 11.4).

For productive laboratory use, one starts from an experimentally measured FRAP curve with unknown diffusion coefficient. Fitting the newly found fluorescence recovery model to such experimental data yields values for the model parameters. Using the functional relationship f finally translates them into the sought-after value of the diffusion constant.

CHAPTER 1. INTRODUCTION

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Chapter 2

3D reconstruction of micrographs

Since it is the goal of this project to take the ideas in [Sbalzarini (2001)] to the third dimension, some 3D representation of the ER is needed. This should be in the form of a triangulated surface in \mathbb{R}^3 meeting the topological requirements stated in [Sbalzarini (2001)]. Fortunately, there exists a good commercial software package for 3D reconstruction from parallel sections called Imaris¹. It not only handles the reconstruction but also does the triangulation of the resulting surface. However, the problems regarding aliasing due to insufficient resolution (see [Shannon (1948)]) and geometrical ambiguity² remain. Imaris "solves" these problems by making some best guesses such that the overall result is closed, connected and as smooth as possible. In addition, Imaris does a Gaussian filtering of the geometry to remove noise, artifacts and geometrical details below a certain level of resolution. The minimum size of the triangles can be controlled as well as the resolution and smoothness of the resulting reconstruction. For the reconstructions made in this work, the data are first filtered with a Gaussian filter of kernel size 2.208 in all three directions. The reconstruction is then done using voxels³ of size $3.15 \times 3.15 \times 2.8$. Moreover, objects are required to be closed and contain more than 2000 triangles.

For illustration purposes, one example of the numerous reconstructions that have been conducted is shown hereafter. All reconstructions were made using Imaris 3.1.3 by BitPlane, Inc. running on a Silicon Graphics Octane workstation under IRIX 6.5. BitPlane kindly granted a demo license of Imaris for the duration of this project.

Figure 2.1 shows the 16 plane parallel slices of stained ER of an example VERO cell as they come from a Leica confocal microscope (see [Sbalzarini (2001)] chapter 9.2 for the detailed protocol). They are then fed to Imaris which does the reconstruction and triangulation. The result is saved as an OpenInventor file (see [SGI (1992a)] and [SGI (1992b)] for the file format specifications) which can be read by the simulation programs developed in this project. Figure 2.2 shows an enlarged part of the resulting triangulated surface; the view in figure 2.3 additionally has its hidden lines removed. Finally, figure 2.4 shows a sample view of the outcome as a 3D shaded surface.

¹http://www.imaris.com

 $^{^{2}}$ two parallel lines on successive slices could for example stem from two parallel planes or from a tube intersected twice parallel to its axis

³The 3D analogue of pixels



Figure 2.1: Sample micrograph stack of stained ER used for 3D reconstruction. The slices are taken at vertical distances of $\Delta z = 0.1 \,\mu\text{m}$ as described in [Sbalzarini (2001)]. (Courtesy of Anna Mezzacasa)



Figure 2.2: Surface triangulation



Figure 2.3: Surface triangulation with hidden lines removed



Figure 2.4: Shaded surface view of 3D reconstruction

Chapter 3

Validating the reconstruction

Now having a triangulated 3D reconstruction of the ER surface, the question arises whether this triangulation meets all the requirements stated in [Sbalzarini (2001)]. Therefore, a subroutine is written that performs some checks on the geometry. In order to be able to describe its algorithms, the requirements are first restated in a slightly different but topologically equivalent form.

3.1 Restatement of requirements

A triangulation is given by a set \aleph of triangles Δ_i . Each triangle is described by the location vectors of its three vertices P_1, P_2, P_3 . Since the ER membrane encloses a closed, connected and contiguous subspace of \mathbb{R}^3 , the triangulation has to meet a minimal set of requirements. We call a triangulation \aleph topologically valid if and only if it meets all of the following:

- 1. \aleph describes a closed, connected 1 surface in \mathbb{R}^3
- 2. No Δ_i must contain any edge of length ≤ 0
- 3. No $P_i(\Delta_i)$ must be inside any other triangle
- 4. Two neighboring triangles must have *exactly* two vertices P_j in common

We call it syntactically valid if it is topologically valid and additionally satisfies:

1. The vector $n = \overrightarrow{P1P2} \times \overrightarrow{P1P3}$, where " \times " is the vector cross product, must point outward (i.e. out of the domain enclosed by \aleph) for every Δ_i .

The first point of the topological validity is particularly difficult to check. The code therefore uses the following equivalent set of criteria:

- A) All edges are of length > 0 (exclude degenerate cases).
- B) No vertex lies inside any other triangle (only allow the minimum set).
- C) Every triangle has exactly 3 neighbors sharing 2 vertices with each (ensure continuity of surface).
- D) No edge of any triangle intersects the face of any other triangle (exclude overlaps).

 $^{^1}meaning$ that any point enclosed by \aleph can be connected to any other point enclosed by \aleph with a continuous curve that never intersects \aleph

The check for syntactical validity is straightforward and the criterion is not reformulated. However it can only be done for sets that are at least topologically valid which means that the two checks are split into different subroutines whereas the second can only be called when the first terminates without error. The two routines are briefly presented in turn below. It is worth mentioning that all floating point comparisons are performed using a global geometry tolerance TOL. Thus an edge is for example considered to be of length 0 if it is of length < TOL and two numbers are considered equal if their difference is less than TOL. Typically, TOL is about 10^{-10} for double precision numbers on a 32bit machine which ensures sufficient robustness against round-off errors while still maintaining good accuracy. The geometry checking routines give us a tool to validate the triangulated 3D reconstructions made by Imaris (cf. chapter 2) in order to ensure proper functioning of the simulation codes that rely on certain properties of the triangulation. If any reconstruction fails the test, it is redone using different filter settings in Imaris until a valid reconstruction results.

3.2 Checking the topological validity

The algorithms described hereafter are implemented in the function chk_topology in chktriang.f90. See appendix C.5.4.7 for a complete listing of the source code. Marginal numbers in the following paragraphs conveniently refer to statement numbers of the listing in appendix C.5.4.7.

The check of requirement A) is trivial, the program just loops over all triangles and calculates the lengths of all the edges. For requirement B), one uses the fact that for every point P inside a triangle, the following holds:

$$P_1 + \alpha a + \beta b = P$$
 for $\alpha > 0, \beta > 0, \alpha + \beta < 1$

where $a = P_2 - P_1$ and $b = P_3 - P_1$ are two edge vectors of the triangle. The corresponding linear set of equations

$$\left[\begin{array}{cc} | & | \\ a & b \\ | & | \end{array}\right] \left[\begin{array}{c} \alpha \\ \beta \end{array}\right] = \left[\begin{array}{c} | \\ P - P_1 \\ | \end{array}\right]$$

is over-defined and will not have a solution for general P. To check whether the system has a solution for a given P, the 2×2 subsystem

$$\begin{bmatrix} & | & | \\ a_{(1:2)} & b_{(1:2)} \\ | & | \end{bmatrix} \begin{bmatrix} \alpha \\ \beta \end{bmatrix} = \begin{bmatrix} & | \\ P_{(1:2)} - P_{1,(1:2)} \\ | \end{bmatrix}$$

is solved for α and β (the code directly contains the symbolic solution for performance reasons) and then one checks if this solution also satisfies the omitted third equation. If this is the case, point P lies inside the triangle if $\alpha > 0$, $\beta > 0$ and $\alpha + \beta < 1$. This is checked for every vertex of every triangle against every other triangle which makes this part of the algorithm $\mathcal{O}(3M^2)$ when M is the total number of triangles. This is acceptable since the number of triangles is typically much smaller than the number of points and rarely exceeds 100'000. Moreover, the algorithm is only run once per triangulation and optimization is often a source of programming errors that should be avoided for a "checking" code.

Requirement C) is again easy to check. For every triangle, all the other triangles aresearched and the number of common vertices is counted. Then the number of its

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Figure 3.1: Geometry of a line intersecting a triangle

neighbors (defined as triangles that share exactly 2 vertices) is determined. Every triangle should have exactly 3 neighbors.

Condition D) is checked by explicitly calculating the intersection point of every edge of every triangle with every other triangle. If such an intersection point is found, the triangulation cannot be topologically valid. However, this method calls for a fast triangle/line intersection subroutine. Such a general routine is implemented in **intersect.f90** (see appendix C.5.4.8). The task is to intersect a line $P + \lambda v$ with a triangle, given by its first vertex P_1 and the two edge vectors a and b according to figure 3.1. The intersection point is given by the following equality:

$$P + \lambda v = \alpha a + \beta b + P_1$$

for some α and β . This corresponds to the linear system of equations:

$$\underbrace{\left[\begin{array}{ccc} | & | & | \\ a & b & v \\ | & | & | \end{array}\right]}_{M} \left[\begin{array}{c} \alpha \\ \beta \\ -\lambda \end{array}\right] = \left[\begin{array}{c} | \\ P - P_1 \\ | \end{array}\right]$$

If det(M) = 0 (up to the global tolerance TOL) no intersection point exists. Else its position is given by $P + \lambda v$ and it lies in the triangle if $\alpha \ge 0$ and $\beta \ge 0$ and $\alpha + \beta \le 1$. The routine returns -1 if no intersection point in the triangle exists and $|\lambda|$ if one is found.

3.3 Checking the syntactical validity

This check is implemented in the function $chk_orientation$ in chktriang.f90 (see appendix C.5.4.7). First, the normal $n = a \times b$ of each triangle is calculated or read from memory if this has been done earlier. The code then checks whether the point that is reached by following the normal a "small distance" starting at the triangle's centroid c is inside the domain or not. "Small distance" is chosen to be 5 times the geometrical tolerance, thus $c + 5TOL \cdot n$ has to lie outside of the domain for the triangle to be properly oriented. This is checked using the function point_in_domain which is discussed in section 7.2.5, algorithm 7.8. $chk_orientation$ returns .TRUE. if all the triangles are oriented correctly, .FALSE. otherwise. 31-32

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Chapter 4

Background on measures and dimensions

4.1 Measures and distributions

Measures and distributions are a central tool in the mathematics of fractals. Since some theoretical studies involving fractal sets will be conducted in later chapters, a few basic concepts are introduced here. However, the basics of fractal geometry and dimensionality as introduced in [Sbalzarini (2001)] will not be repeated but assumed to be known. Only measures and distributions on the Euclidean \mathbb{R}^n are needed for the rest of this text. Note that what will be called "measure" hereafter is often termed "outer measure" in general texts on measure theory.

Simply put, a measure is a way of assigning a scalar numerical value to a set such that the principle "the whole is the sum of its parts" applies¹. Readers not yet familiar with the concept of measures may find it helpful to imagine them as physical "mass distributions" or "charge distributions". By restricting this section to measures on \mathbb{R}^n , many of the awkward features that can occur in more general spaces are avoided.

Definition 4.1 (Measure). Let $X \subset \mathbb{R}^n$. We call $\mu : X \mapsto \mathbb{R}^+$ a *measure* on X if μ assigns a non-negative number (possibly ∞) to each subset of X such that:

- (a) $\mu(\emptyset) = 0$
- (b) $A \subset B \Rightarrow \mu(A) \leq \mu(B)$
- (c) if A_1, A_2, \ldots is a countable sequence of sets, then:

$$\mu\left(\bigcup_{i=1}^{\infty} A_i\right) \leqslant \sum_{i=1}^{\infty} \mu(A_i)$$

Thus (a) requires the empty set to have measure zero and (b) states that the bigger set has the larger measure. Property (c) ensures that the measure of any set is no more than the sum of the measures of the pieces in any countable decomposition. Exact equality holds for "nice" disjoint sets A_i , in particular (but not only) if A_i are disjoint *Borel sets*².

 $^{^1\}mathrm{Strictly}$ speaking this is only true if the set can reasonably be decomposed into a finite or countable number of disjoint subsets.

 $^{^{2}}$ Every set that can be constructed starting with open or closed sets and taking countable unions or intersections a finite number of times is a Borel set. All sets encountered in this text are Borel sets.

This leads to the concept of *measurability*. Given a measure μ there is a family of subsets of X on which μ behaves in a nice additive way such that property (c) in definition 4.1 holds for equality. This family is called μ -measurable and termed \mathcal{M} . The exact definition is:

Definition 4.2 (Measurable). A set $A \subset X$ is called μ -measurable if

$$\mu(E) = \mu(E \cap A) + \mu(E \setminus A) \qquad \forall \quad E \subset X$$

 μ is termed a *Borel measure* on $X \subset \mathbb{R}^n$ if all Borel subsets of X are μ -measurable. It may be shown that μ is a Borel measure if and only if $\mu(A \cup B) = \mu(A) + \mu(B)$ whenever $A, B \subset X$. A Borel measure is termed *Borel regular* if every subset of X is contained in a Borel set of the same measure. All measures that will be encountered in this text are Borel regular on \mathbb{R}^n or the pertinent subset thereof. Thus, to simplify notation, the term "measure" will mean "Borel regular outer measure" and "set" will mean "Borel set" throughout this report.

A measure μ on X with $\mu(X) < \infty$ is called *finite*; if $\mu(A) < \infty$ for every bounded set A, it is *locally finite*. The *support* of μ is the set on which the "mass" of μ is concentrated, thus the smallest closed set X such that $\mu(\mathbb{R}^n \setminus X) = 0$. The support of any measure is always closed and termed $spt(\mu)$. Its definition is:

Definition 4.3 (Support). The *support* of a measure μ is the smallest closed set with complement of measure zero, thus:

$$spt(\mu) = X \setminus \bigcup \{U : U \text{ is open and } \mu(U) = 0\}$$

A measure on a bounded subset X of \mathbb{R}^n for which $0 < \mu(\mathbb{R}^n) < \infty$ is called a *distribution*. Imagine for example a finite mass spread over the set X in some way. Then, the mass contained in any subset A of X is given by $\mu(A \subset X)$. As distributions are a subclass of measures, all properties for measures also hold for distributions.

Sometimes, a restriction of a measure to a set E is used. Let μ be a measure on \mathbb{R}^n and E a Borel subset of \mathbb{R}^n . We can define a new measure ν on \mathbb{R}^n through $\nu(A) = \mu(E \cap A)$ for all A. This new measure ν is called restriction of μ to E and it is a measure on \mathbb{R}^n whose support is entirely enclosed in E. The restriction is thus a method of defining a derived measure with support in a certain region. Every μ -measurable set is also ν -measurable and, provided E is measurable and $\mu(E) < \infty$, then ν is a Borel regular measure.

Some examples of frequently used used measures are given hereafter. The corresponding proofs that the described objects really are measures in the sense of definition 4.1 are omitted to keep this section as compact as possible. They are well known and contained in every book on measure theory.

Example 4.1 (Counting measure). For each subset A of \mathbb{R}^n , let $\mu(A)$ be the number of points in A if A is finite and ∞ otherwise. μ is a measure on \mathbb{R}^n .

Example 4.2 (Point measure). Let *a* be a point in \mathbb{R}^n . Define $\mu(A) = 1$ if $a \in A$ and $\mu(A) = 0$ if $a \notin A$. Then, μ is a measure with support $\{a\}$ that we think of as unit point mass concentrated at *a*.

Example 4.3 (Lebesgue measure). The Lebesgue measure on \mathbb{R}^n is the natural extension to the terms "length", "surface area" and "volume". It is based on the "*n*-dimensional Volume" of the coordinate parallelepiped $A = \{(x_1, \ldots, x_n) \in \mathbb{R}^n : a_i \leq x_i \leq b_i\}$ defined by:

$$\operatorname{Vol}^{n}(A) = \prod_{i=1}^{n} (b_{i} - a_{i})$$

Then the *n*-dimensional Lebesgue measure \mathcal{L}^n is defined by:

$$\mathcal{L}^{n}(A) = \inf\left\{\sum_{i=1}^{\infty} \operatorname{Vol}^{n}(A_{i}) : A \subset \bigcup_{i=1}^{\infty} A_{i}\right\}$$

where the infimum is over all coverings of A by countable collections of parallelepipeds. With some effort it may be shown that \mathcal{L}^n is indeed a Borel regular measure on \mathbb{R}^n such that $\mathcal{L}^n(A)$ equals the *n*-dimensional volume of A if A is a parallelepiped or any other set for which the volume can be calculated using the usual rules of mensuration.

Later, integration with respect to a measure μ will also be used. Let $f : X \mapsto \mathbb{R}^+ \cup \{0\}$ be a non-negative *simple function*, i.e. a function that only takes a finite number of different values $a_1, \ldots, a_k \in \mathbb{R}^+ \cup \{0\}$. Then integration with respect to the measure μ on X is defined as:

$$\int f \, d\mu = \sum_{i=1}^k a_i \mu \{ x : f(x) = a_i \}$$

Integration of more general functions is defined using approximation by simple functions. We term $f : X \mapsto \mathbb{R}$ a measurable function if for all $c \in \mathbb{R}$ the set $\{x \in X : f(x) < c\}$ is a measurable set. In particular for a Borel measure μ all continuous functions are measurable. The integral of any non-negative measurable function $f : X \mapsto \mathbb{R}^+ \cup \{0\}$ is defined by:

$$\int f \, d\mu = \sup\left\{\int g \, d\mu : g \text{ is simple, } 0 \leqslant g \leqslant f\right\}$$

This value may also be infinite. Finally, for a measurable function $f : X \to \mathbb{R}$ that can also take negative values, the definition is completed writing $f_+(x) = \max\{f(x), 0\}$ and $f_-(x) = \max\{-f(x), 0\}$, so that $f(x) = f_+(x) - f_-(x)$, and defining:

$$\int f \, d\mu = \int f_+ \, d\mu - \int f_- \, d\mu$$

provided both $\int f_+ d\mu$ and $\int f_- d\mu$ are finite. For integrals over measures, the usual properties of integrals including their linearity and convergence theorems hold unchanged. If A is a Borel subset of X, the definite integral over $A \subset \mathbb{R}^n$ is defined by:

$$\int_A f \, d\mu = \int f \chi_A \, d\mu$$

where $\chi_A : \mathbb{R}^n \to \mathbb{R}$ is the *indicator function* of A with $\chi_A(x) = 1$ if $x \in A$ and $\chi_A(x) = 0$ otherwise.

4.2 Hausdorff's measure and dimension

Having now the basic knowledge of measures and distributions at hand, the special Hausdorff measure can be introduced. This will lead to the definition of Hausdorff's dimension ([Hausdorff (1919)]) which is the oldest and most important metric dimension when studying fractal sets. This is due to the fact that it is defined for any set and is mathematically practical as it is based on a comparatively simple measure. Although [Sbalzarini (2001)] already gave a "free hanging" definition of Hausdorff's dimension, it is restated here based on the more fundamental and solid background of measure theory since it is important to be familiar with Hausdorff's measure and dimension to understand the mathematics of fractals.

First, recall the following two definitions from basic set theory:

Definition 4.4 (Diameter). The *diameter* of a non-empty subset U of the ndimensional Euclidean space \mathbb{R}^n is defined as:

$$|U| = \sup \{|x - y| : x, y \in U\}$$

This means that the diameter is the largest occurring Euclidean distance between any two points in U.

Definition 4.5 (\delta-cover). A finite or countable collection of subsets $\{U_i\}$ of \mathbb{R}^n is called a δ -cover of a set $E \subset \mathbb{R}^n$ if $0 < |U_i| \leq \delta \forall i$ and $E \subset \bigcup_{i=1}^{\infty} U_i$.

This means that E is completely covered by a countable collection of subsets of diameter less or equal δ . Now let E be a subset of \mathbb{R}^n and s a non-negative number. For all δ we define the quantity:

$$\mathcal{H}^{s}_{\delta}(E) = \inf\left\{\sum_{i=1}^{\infty} |U_{i}|^{s} : \{U_{i}\} \text{ is a } \delta\text{-cover of } E\right\}$$
(4.1)

As δ decreases, the class of valid δ -covers of E is reduced. The infimum therefore increases and approaches a limit as $\delta \downarrow 0$. We write:

$$\mathcal{H}^{s}(E) = \lim_{\delta \to 0} \mathcal{H}^{s}_{\delta}(E) \tag{4.2}$$

This limit exists for all $E \subset \mathbb{R}^n$, although its value can be 0 or ∞ . We term $\mathcal{H}^s(E)$ the *s*-dimensional Hausdorff measure of E. With considerable effort is has been shown that the Hausdorff measure is a measure in the sense of definition 4.1. It is even a Borel regular measure on \mathbb{R}^n .

Hausdorff's measure generalizes Lebesgue's measure, so that $\mathcal{H}^0(E)$ gives the number of points in E, $\mathcal{H}^1(E)$ the length of a smooth curve E, $\mathcal{H}^2(E)$ the (normalized) surface area of a region E, etc. In general, for every Borel subset E of \mathbb{R}^n , is holds:

$$\mathcal{H}^{n}(E) = c_{n}\mathcal{L}^{n}(E)$$
 where $c_{n} = \frac{\pi^{\frac{n}{2}}}{2^{n}\left(\frac{n}{2}\right)!}$

Often, the Hausdorff measure of the image of a set under a Hölder mapping is needed. Let $E \subset \mathbb{R}^n$ and $f: E \mapsto \mathbb{R}^m$ a Hölder continuous³ function of order α , so that

$$|f(x) - f(y)| \le c |x - y|^{\alpha}$$

for all $x, y \in E$ and some $c, \alpha > 0$.

³A function is called *Hölder continuous* of order α if $|f(x) - f(y)| \leq c |x - y|^{\alpha} \quad \forall x, y$
Then

$$\mathcal{H}^{\frac{s}{\alpha}}(f(E)) \leqslant c^{\frac{s}{\alpha}}\mathcal{H}^{s}(E)$$

for all s. The proof of this can for example be found in [Falconer (1997)]. The special case of a *Lipschitz mapping* f (i.e. $\alpha = 1$) is particularly interesting. It also comprises the case when f is a regular similarity transformation of ratio λ . Then:

$$\mathcal{H}^s(\lambda E) = \lambda^s \mathcal{H}^s(E)$$

which generalizes the familiar scaling properties of length, area, volume, etc. to the fractal case. Also notice that Hausdorff's measure is invariant under isomeries⁴, translation and rotation as could be expected for a generalized "volume".

It it easy to show from equations 4.1 and 4.2 that for all sets $E \subset \mathbb{R}^n$ there is a critical number D_H such that $\mathcal{H}^s(E) = \infty$ if $s < D_H$ and $\mathcal{H}^s(E) = 0$ if $s > D_H$. This number D_H is called *Hausdorff dimension* of the set E and can be compactly defined as follows:

Definition 4.6 (Hausdorff dimension). The number:

$$D_H = \inf \left\{ s : \mathcal{H}^s(E) = 0 \right\} = \sup \left\{ s : \mathcal{H}^s(E) = \infty \right\}$$

is called Hausdorff dimension of E.

Thus D_H is the value of s for which the Hausdorff measure jumps from ∞ to 0 (see figure 4.1). When $s = D_H$, the Hausdorff measure can be either 0, ∞ or, in the nicest situation, $0 < \mathcal{H}^{D_H}(E) < \infty$.



Figure 4.1: Behavior of the Hausdorff measure in the vicinity of D_H

Definition 4.7 (s-set). A set $E \subset \mathbb{R}^n$ of Hausdorff dimension D_H with $0 < \mathcal{H}^{D_H}(E) < \infty$ is called an *s-set*.

As an example, imagine a flat circular disk in \mathbb{R}^3 with radius 1. From known properties of length, area and volume, we get: $\mathcal{H}^1 = \text{length} = \infty$, $0 < \mathcal{H}^2 = \frac{\pi}{4} \text{area} < \infty$ and $\mathcal{H}^3 = \frac{\pi}{6}$ volume = 0. Thus the disk is an *s*-set with Hausdorff dimension $D_H = 2$ which is consistent with the usual Euclidean definition of dimension.

⁴Mappings for which $|f(x) - f(y)| = |x - y| \quad \forall x, y$

The transformation properties of Hausdorff's dimension follow from the ones of Hausdorff's measure. In particular, for a Hölder mapping $f : E \mapsto \mathbb{R}^m$ of order α we get:

$$D_H(f(E)) \leqslant \frac{1}{\alpha} D_H(E)$$

A fundamental property of Hausdorff's dimension is its invariance under Bi-Lipschitz mappings⁵ as can easily be seen from above transformation property. This means that if two sets have different Hausdorff dimensions, there is no Bi-Lipschitz mapping between them. Analogous to topology theory where two homeomorphic sets are considered equivalent, two fractals are considered equivalent if they have the same Hausdorff dimension, i.e. there exists a Bi-Lipschitz continuous mapping between them.

Moreover, Hausdorff's dimension fulfills a whole set of properties that are natural to be expected from any meaningful definition of a dimension. [Falconer (1990)] for example contains an exhaustive list of them. However, their details are not relevant for this work and will not be stated here.

4.3 Other fractal dimensions

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In addition to the very important Hausdorff dimension, there is an infinite number of other possible definitions of metric dimensions, the most important of whitch is the *box counting dimension* for this work. Although Hausdorff's dimension comes in very handy for theoretical considerations, it has a fundamental drawback: it is not measurable or computable in most cases. The box counting dimension however is quite easy to be estimated in computer simulations what makes it a valuable tool for the investigations to come.

It is however important to keep in mind that fractals are dimensionally discordant sets (cf. [Sbalzarini (2001)]), so that different meaningful definitions of dimension can yield different numerical values for the same set and even "very similar" definitions can have grossly different properties. A common property of all dimensions is that they are based on "measurements with respect to a scale δ " in a metric space. For a large class of dimensions, a power law between the measured length in multiples of δ (N_{δ}) and the metric size of δ is then assumed:

$$N_{\delta}(E) \sim c\delta^{-s}$$

The exponent s is called dimension of E and the constant c is referred to as sdimensional length of E. Using log-log-plots or statistical fitting allows to estimate s from a number of measurements with different length scales δ . This way of defining dimension is not based on measure theory which makes it awkward for theoretical usage. It is however consistent with measure-based dimensions in the sense that the numerical values for dimensionally concordant sets are equal. It happens however, that the assumption of a power law is in fact wrong, in which case the points in a plot of log(N_{δ}) vs. log(δ) will not be on a straight line. Dimensions that are based on the power law assumption are then meaningless for the given set and should not be used. This emphasizes the need to not blindly apply linear regression fitting to such a data set but look at the distribution of points as well.

Moreover, when defining new dimensions, one should make sure that they meet all the requirements for "meaningful" dimensions such as monotony, stability, Bi-Lipschitz invariance, etc. as for example listed in [Falconer (1990)].

⁵A function $f: E \mapsto \mathbb{R}^m$ is called *Bi-Lipschitz* if $c_1|x-y| \leq |f(x) - f(y)| \leq c_2|x-y|$ for all $x, y \in E$ and certain $0 < c_1 \leq c_2 < \infty$.

The power law dimensions that will be used in this work are the box counting dimension and the interior Minkovski dimension that will be described hereafter. Another useful dimension is the *packing dimension* which is related to Hausdorff's dimension but instead of covering the set with δ -subsets, it is packed with disjoint δ -balls. It will however not be used in this text and therefore its description is omitted.

4.3.1 The box counting dimension

The box counting dimension has first been used in the 1930s. Due to its popularity, a large number of different names for it exist, among them Kolmogorov entropy, Pontrjagin-Kolmogorov dimension (as used in [Sbalzarini (2001)]), Entropy dimension, capacity dimension, logarithmic density or information dimension. To support intuition, the term "box counting dimension" will be used in this text.

For E being a non-empty bounded subset of \mathbb{R}^n let $N_{\delta}(E)$ be the smallest number of sets of diameter $\leq \delta$ that can cover E. The *lower* and *upper box counting dimensions* of E are defined as:

$$\underline{D_B} = \liminf_{\delta \to 0} \frac{\log N_\delta(E)}{-\log \delta}$$

and

$$\overline{D_B} = \limsup_{\delta \to 0} \frac{\log N_{\delta}(E)}{-\log \delta}$$

respectively. If these are equal, the common value is simply called *box counting* dimension of E:

$$D_B = \lim_{\delta \to 0} \frac{\log N_{\delta}(E)}{-\log \delta} \tag{4.3}$$

Fortunately, there are a number of equivalent forms of this definition which are practically useful. The values of above limits namely remain unaltered if $N_{\delta}(E)$ is taken to be any of the following:

- 1. the smallest number of sets of diameter $\leqslant \delta$ that can cover E
- 2. the smallest number of closed balls of radius δ that can cover E
- 3. the smallest number of cubes of edge length δ that can cover E
- 4. the largest number of disjoint balls of radius δ with centers in E
- 5. the number of δ -mesh cubes⁶ that intersect E

Variant 5 is the one that has been used in [Sbalzarini (2001)] and will also be used in the present work since the mesh needed for it is naturally given by the pixels in a 2D image or their 3D equivalent, the voxels, in a 3D volume.

⁶a δ -mesh cube is a cube of the form $[m_1\delta, (m_1+1)\delta) \times \ldots \times [m_n\delta, (m_n+1)\delta)$ where m_1, \ldots, m_n are integers.

Although the box counting dimension is defined in a completely different way as the Hausdorff dimension, there is a deterministic relationship between them as stated in the following inequality:

$$D_H \leqslant \underline{D_B} \leqslant \overline{D_B} \tag{4.4}$$

This is due to the fact that the box counting dimension is (in variants 2 to 5) based on covering the set with elementary subsets (balls, cubes, etc.) of equal and fixed size whereas the Hausdorff dimension relies on a δ -cover with subsets of different size smaller than δ . Therefore, Hausdorff's dimension employs a finer covering with the biggest subset having the size of the ones used for the box counting.

In practice, most definitions of dimensions take values between the Hausdorff and upper box counting dimension, so if it can be shown that $\overline{D_B} = D_H$ then all common definitions of dimension will take this common value.

For a complete discussion of advantages and disadvantages of the box counting dimension, see [Falconer (1990)] and [Falconer (1997)]. For our purposes it is sufficient to state that the box counting dimension is unusable for sets that only consist of a countable number of isolated points⁷. Since the ER and its triangulated description is a continuous set with no isolated points, the box counting dimension can be used without danger.

4.3.2 The Minkovski dimension

A variant of the box counting dimension that is sometimes used for one-sided closed sets is the Minkovski dimension. It is based on the *n*-dimensional volume of the δ -neighborhood or δ -parallel body E_{δ} of E given by:

$$E_{\delta} = \{ x \in \mathbb{R}^n : |x - y| \leq \delta \text{ for all } y \in \partial E \}$$

Then for $E \subset \mathbb{R}^n$:

$$\underline{D_M} = \underline{D_B} = n - \limsup_{\delta \to 0} \frac{\log \mathcal{L}^n(E_\delta)}{\log \delta}$$
$$\overline{D_M} = \overline{D_B} = n - \liminf_{\delta \to 0} \frac{\log \mathcal{L}^n(E_\delta)}{\log \delta}$$

and if they coincide:

$$D_M = D_B = n - \lim_{\delta \to 0} \frac{\log \mathcal{L}^n(E_\delta)}{\log \delta}$$

These definitions are actually equivalent to the ones of the box counting dimension. In the context of using the Lebesgue measure they are however referred to as *Minkovski dimensions*.

Important variants are the *inner* and *outer Minkovski dimension* for which only the inner or outer δ -neighborhood of a closed set is taken. This is for example useful to study the influence of the fractality of the boundary of a domain on the solution of a PDE solved in its interior (cf. section 6.2.7).

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⁷this is due to the fact that the closure of a set E has the same box counting dimension as the set E itself

Chapter 5

Measuring the fractal dimension

Now having the necessary preliminaries on the notion of dimensions, the fractal dimension of the ER surface in \mathbb{R}^3 will be measured. As stated in [Sbalzarini (2001)], this dimension could be an important parameter for any geometry model. It will be shown later in this report that equality of the dimension is indeed a necessary condition for a geometrical model to be valid. In principle, one would wish to measure Hausdorff's dimension. However, this is not possible for arbitrary geometries (cf. section 4.3) and therefore the box counting dimension as defined in section 4.3.1 will be used instead.

A box counting algorithm similar to the one described in [Sbalzarini (2001)] requires the data to be in some sort of "box oriented form". For 2D images, these "box units" are naturally given by the image's pixels. The triangulated 3D objects dealt with now however pose some problems. According to section 4.3.1 there are various equivalent ways of defining the box counting dimension. One possibility would for example be to successively discard vertices of the triangulation at hand and retriangulate the surface at every step. Due to topological limitations (recall that the ER's surface has to be contiguous) this would be not only cumbersome but also dangerous and geometry checking and possibly fixing would be needed at every step. It is therefore easier to use, again, the "pixel" oriented approach. The 3D analogue of a pixel is called a *voxel* and can be defined as a finite parallelepiped in \mathbb{R}^3 , thus:

$$V_{i,j,k} = \{ (x, y, z) : (i-1)\delta x \leq x - x_l < i\delta x, \ (j-1)\delta y \leq y - y_l < j\delta y, \\ (k-1)\delta z \leq z - z_l < k\delta z \}$$

where (x_l, y_l, z_l) are the coordinates of the lower left corner of the bounding box. This is an extension of what has been called a δ -mesh cube in section 4.3.1 to an anisotropic mesh. Hence the first step will be a conversion of the triangulated surface to a voxel set.

5.1 Generating a voxel representation

Recall that variant 5 of the definition of the box counting dimension in section 4.3.1 states that all voxels that intersect the geometry in question are to be counted. The ER surface is given as a triangulated set \aleph as defined in [Sbalzarini (2001)], meeting

all the requirements stated in section 3.1. The space \mathbb{R}^3 in which this triangulation is embedded is covered with a regular cartesian grid with grid spacings δx , δy and δz defining the voxels. All voxels that are intersected by \aleph are then set to one, the others to zero thus yielding the required binary voxel representation. The algorithm proceeds as follows:

Algorithm 5.1 (Voxelize).

- Step 1: Generate cartesian grid with spacings δx , δy , δz and determine the bounding box $[x_l, x_h] \times [y_l, y_h] \times [z_l, z_h]$ of the geometry as: $\{x, y, z\}_l = \min_p \{x_p, y_p, z_p\}, \{x, y, z\}_h = \max_p \{x_p, y_p, z_p\}$ where p loops over all triangle vertices in \aleph .
 - Step 2: Define voxels $V_{i,j,k} = [x_l + (i-1)\delta x, x_l + i\delta x) \times [y_l + (j-1)\delta y, y_l + j\delta y) \times [z_l + (k-1)\delta z, z_l + k\delta z)$ for i, j, k = 1, 2, 3, ...
- **33** Step 3: Initialize all voxels to zero: $V_{i,j,k} = 0 \quad \forall (i, j, k)$
 - Step 4: Loop over all triangles $\Delta \in \aleph$
 - 4.1: For each vertex P of Δ determine the voxel it is contained in: (i, j, k) = $int((P_x - x_l)/\delta x, (P_y - y_l)/\delta y, (P_z - z_l)/\delta z) + (1, 1, 1)$ and set voxel (i, j, k) to one¹
 - 4.2: Determine the bounding box of voxels for Δ : $(i_u, j_u, k_u) = \max(i, j, k)$ and $(i_l, j_l, k_l) = \min(i, j, k)$
 - 4.3: For each voxel (i, j, k) in the bounding box $\{i_l \leq i \leq i_u, j_l \leq j \leq j_u, k_l \leq k \leq k_u\}$ solve:

	$\left[\begin{array}{c} \\ a \\ \end{array}\right]$	 b 	$\left[\begin{array}{c} \\ n \\ \end{array} \right]$	$\begin{bmatrix} \alpha \\ \beta \\ \gamma \end{bmatrix}$	=	$\begin{bmatrix} & \\ c_V - P_1 \\ & \end{bmatrix}$	
78-90	_		_				

for α , β and λ where $a = P_2 - P_1$ and $b = P_3 - P_1$ are two edge vectors of the triangle Δ , n is its outer normal and P_1 the location vector of vertex 1 of triangle Δ . c_V is the location vector of the centroid of voxel (i, j, k), defined as: $c_V = [x_l + (i - 1/2)\delta x, y_l + (j - 1/2)\delta y, z_l + (k - 1/2)\delta z]^\top$

4.4: Set voxel (i, j, k) to one if $\alpha \ge 0$ and $\beta \ge 0$ and $\alpha + \beta \le 1$ and $Y = c_V + \lambda n \in V_{i,j,k}$ which means that the centroid of voxel (i, j, k) is over the triangle Δ and the intersection point of the normal through the centroid of the voxel with the triangle is inside the voxel. (see also figure 5.1)

An implementation of this algorithm can be found in the subroutine Voxelize.f90 in appendix C.5.4.12. A few remarks are in place. The margin numbers refer to the corresponding statements in the Fortran code. In step 4.3, a linear system of equations is to be solved. To speed up the algorithm, this has been done symbolically in advance and the general solution is programmed directly.. The condition in step 4.4 can also be taken to be $\lambda \leq \frac{1}{2} \max \{\delta x, \delta y, \delta z\}$. This however turned out to be problematic if the voxels are elongated parallelepipeds rather than cubes. For near cubic voxels and visualization purposes this condition is nevertheless sometimes preferable since it yields nicer images containing less voxels. For the box counting algorithm it is however useless. Figure 5.1 illustrates the condition of step 4.4.

22

19-36

59-69

70-75

91-93

48-58 88-90

¹Special index treatment is needed for the rightmost voxel.



Figure 5.1: Geometric situation of a voxel and a triangle

Validation

The algorithm has been tested for various simple shapes. As an example, figure 5.3 shows a voxel representation of the pyramid depicted in figure 5.2. Figure 5.4 shows a sample picture of a triangulated ER surface and figure 5.5 its voxel representation. The visualizations are made using the freely available OpenDX package (former IBM DataExplorer)².



Figure 5.2: Pyramid wireframe model



Figure 5.3: Voxel representation

5.2 A 3D box counting algorithm

Now having a voxel set representation $\{(i, j, k) : V_{i,j,k} = 1\}$ for the triangulated set \aleph , the determination of the box counting dimension according to section 4.3.1 is straightforward. Algorithms for 2D images exist for quite some time (see e.g. [Liebovitch & Toth (1989)]) however the first application to 3D systems has been

²http://www.opendx.org



Figure 5.4: Triangulated surface



done by [Stoll, Stern & Stucki (1996)], who unfortunately did not publish the algorithm as such. To ensure comparability of results and since we are dealing with triangulated sets rather than point sets, the same procedure as in [Sbalzarini (2001)] is used here but generalized to three dimensions. For details on why this method works or why outlining and low-pass filtering are needed, the reader is referred to that report. The algorithm is as follows:

Algorithm 5.2 (Box counting).

37	Step 1: Convert the triangulation \aleph to a voxel representation with voxels of size δx , δy , δz using algorithm 5.1.
42	Step 2: Determine the number of resolution reduction steps until the smallest dimension has less than 16 voxels: $n = \text{ceiling} (\log (\min \{\delta x, \delta y, \delta z\}/16) / \log 2).$
	Step 3: For each reduction step $i = 1$ to n
51-78	3.1: Outline the voxel set: set all voxels to zero that are completely surrounded by voxels of value 1.
79-80	3.2: Count the number of voxels of value one: $\eta_i = \#\{(i, j, k) : V_{i,j,k} = 1\}$ and store the edge length scale of the current voxels: $\sigma_i = \delta x/2^{i-1}$
100-115	3.3: Discrete low-pass filter: replace each voxel by the average of its 26 neighbors and itself.
120	3.4: Binarize set: set all voxels with value $<$ threshold to zero, all others to one.
97-123	3.5: Resolution reduction: delete every second row, column and z-plane from the 3D voxel set.
124-146	3.6: Reallocate voxel array to new size and update all variables.
147-148	Step 4: Get final values: $\eta_{n+1} = \#\{(i, j, k) : V_{i,j,k} = 1\}$ and $\sigma_{n+1} = \delta x/2^n$
149-152	Step 5: Let $y_i = \log \eta_i$ and $s_i = \log \left(\frac{1}{\sigma_i}\right) \forall i = 1, 2, \dots, n+1$
158-169	Step 6: Do a linear least squares fit of y_i vs. s_i to get the slope:
	$\hat{a} = \sum_{i=1}^{n+1} \left((n+1)s_i - \sum_{i=1}^{n+1} s_i \right) y_i$

Step 7: Deallocate all memory and return the estimated box counting dimension: $D_B = \hat{a}$

This algorithm is implemented as a Fortran function in BCdim.f90 (see appendix C.5.4.13, to which the statement numbers in above algorithm refer) which directly returns the numerical value of the box counting dimension. With considerable effort it could be optimized according to [Hou, Gilmore, Mindlin & Solari (1990)]. However it turned out that its performance is sufficient for the present application so that one can do without further optimization. A few remarks to the algorithm are useful: In step 2, the number of reduction steps is calculated until the smallest occurring number of voxels in a row is less than 16. This is needed since for smaller numbers the accuracy of the voxel representation deteriorates rapidly. Moreover, the definition of the box counting dimension in equation 4.3 contains the limit of $\delta \rightarrow 0$ which means that too large voxels are not significant anyway. The outline operation in step 3.1 removes all voxels that are not part of the hull. Since we are considering a surface in space, this is needed to keep the geometry from "thickening" towards a space-filling body which would grossly distort the value of its dimension. The way this outlining is done is a 3D generalization of MATLAB's bwperim command. In step 3.4 the set is binarized again. The threshold used here is a user parameter and should be chosen as high as possible such that the resulting dimensions are valid (i.e. smaller than the topological dimension of the space the object is embedded in). Some authors (e.g. [Liebovitch & Toth (1989)]) state that it is advisable to exclude the first few measurement points (i.e. the ones at the finest resolution level) from the least squares fit since they may contain artifacts and measurement noise. In our case however this turned out to be not necessary since the initial voxel representation is generated in the computer (thus free of noise) and already is quite coarse (due to memory limitations). Moreover, the measurement points all lie on a straight line pretty well. This means that there are no inaccurate points that should be excluded and therefore they are kept in order to have more measurement points available.

The least squares fit in step 6 is done according to the following calculation. Let $Y = [y_i]^\top \in \mathbb{R}^{n+1}$ be the vector of the logarithms of the measured voxel numbers and $\Theta = [a \ b]^\top \in \mathbb{R}^2$ the vector of unknown parameters. According to equation 4.3 we now assume a linear relation between the logarithm of the voxel number and the logarithm of the corresponding inverse voxel size: $y_i = as_i + b$. Since this will not be exactly true for all i, we define the residual $\varepsilon_i = as_i + b - y_i$ and require the sum over all i of its squares to be minimum. Defining the matrix

$$\Phi = \left[\begin{array}{ccc} s_1 & s_2 & \dots & s_{n+1} \\ 1 & 1 & \dots & 1 \end{array} \right]^\top$$

the whole problem can be compactly formulated for all i using matrix vector notation:

$$Y = \Phi\Theta$$
$$\varepsilon = \Phi\Theta - Y$$
$$\varepsilon^{\top}\varepsilon \stackrel{!}{=} \min$$

Inserting the definition of ε into the last equation gives:

$$\varepsilon^{\top}\varepsilon = \left(\Phi\Theta - Y\right)^{\top}\left(\Phi\Theta - Y\right) = \Theta^{\top}\Phi^{\top}\Phi\Theta - \Theta^{\top}\Phi^{\top}Y - Y^{\top}\Phi\Theta + Y^{\top}Y$$

The minimization condition requires the derivative with respect to Θ to vanish, thus:

$$\frac{\partial}{\partial \Theta} \left(\varepsilon^{\top} \varepsilon \right) = 2 \Phi^{\top} \Phi \Theta - 2 \Phi^{\top} Y \stackrel{!}{=} 0$$

Solving this equation for Θ yields the least squares estimate:

$$\widehat{\Theta}_{LS} = \left(\Phi^{\top}\Phi\right)^{-1}\Phi^{\top}Y$$

This is however still not very useful for our purpose. Explicitly inserting the definition of Φ gives:

$$\Phi^{\top} \Phi = \left[\begin{array}{cc} \sum s_i^2 & \sum s_i \\ \sum s_i & n+1 \end{array} \right] \in \mathbb{R}^{2 \times 2}$$

where all the sums are taken over i = 1, ..., n + 1. Its inverse is then given by:

$$\left(\Phi^{\top}\Phi\right)^{-1} = \frac{1}{\Delta} \left[\begin{array}{cc} n+1 & -\sum s_i \\ -\sum s_i & \sum s_i^2 \end{array} \right] \in \mathbb{R}^{2 \times 2}$$

where the determinant Δ is given by:

$$\Delta = (n+1)\sum s_i^2 - \left(\sum s_i\right)^2$$

Using this allows to explicitly calculate the least squares estimates of the unknown line parameters a and b as:

$$\hat{a} = \sum_{i=1}^{n+1} \left((n+1)s_i - \sum_{i=1}^{n+1} s_i \right) y_i$$
$$\hat{b} = \sum_{i=1}^{n+1} \left(\sum_{i=1}^{n+1} s_i^2 - s_i \sum_{i=1}^{n+1} s_i \right) y_i$$

The slope \hat{a} approximates the box counting dimension D_B according to its definition in equation 4.3.

Validation

The algorithm as presented above is now validated using some simple Euclidean shapes. The first case that is presented here consists of two plane parallel triangles³ (see figure 5.6) and the second of a cubic box of edge length 4 (see figure 5.7).



Figure 5.6: Two parallel triangles

Figure 5.7: Cubic box

³because a single triangle would not be a 3D object



Figure 5.8: $\log \eta$ vs. $\log(1/\sigma)$ (red crosses) for the triangles and its least squares fit (black dotted)



Figure 5.9: $\log \eta$ vs. $\log(1/\sigma)$ (red crosses) for the box and its least squares fit (black dotted)

Figure 5.8 shows the result for the triangles. The measurement points of the box counting algorithm are indicated as red crosses and the least squares fit as a dotted black line. As expected for a dimensionally concordant set, the crosses lie on a straight line since the power law assumption is exactly valid (cf. section 4.3). The dimension of the triangles is measured to be 1.9827. For the box we get 2.0163 and the points in figure 5.9 are no longer exactly on the straight line due to the higher complexity of the shape. (but they are still very close to it.)

As it is known that both shapes are of dimension 2.0 for all meaningful definitions of dimension (recall that we just consider their surface, not the volume), it can be stated that the box counting algorithm yields the correct results up to one decimal digit.

5.3 Results for the ER

The box counting algorithm described so far is now used to estimate the dimension of the surface of the ER. It will then be possible to compare it to the theoretical prediction made in [Sbalzarini (2001)] on page 64 that the dimension should be between 2.8 and 3.0 with a higher probability for the lower bound.

Nine different triangulated 3D reconstructions of ER shapes are considered which allows some statistical averaging. The following figures show views of these nine samples together with their names.



bip2



erp573_1



Figure 5.10: Shaded surface views of all reconstructed ER samples.

The threshold for the binarization step (step 3.4 in algorithm 5.2) is chosen according to the rules in [Sbalzarini (2001)]. The largest value still yielding meaningful results turns out to be 0.3 for most samples, 0.29 for erp573_2 and erp574_2 and 0.27 for erp573_3. Choosing only slightly larger values causes the measured dimension to jump up to values larger than 3 which is impossible due to the Szpilrajn inequality. The exceptions are caused by variations in the reconstruction of the micrographs. Thus the right threshold seems to be around 0.3. The following values are obtained for the different samples:

	0.3	0.29	0.27
bip2	2.8166	2.5655	
clx	2.7195	2.4906	
erp57	2.4839	2.3379	
erp572	2.9876	2.6684	
$erp573_1$	2.9554	2.5391	
$erp573_2$	3.4788	2.9704	
erp573_3	>4	3.2033	2.9960
$erp574_1$	2.8836	2.5648	
$erp574_2$		2.5628	

Table 5.1: Measured box counting dimensions

The numbers in bold are the ones for the highest threshold still giving meaningful results. These are the estimated dimensions. Their average value is 2.8195 which is indeed between 2.8 and 3.0 and closer to 2.8, as predicted by the fractal projection theorem in [Sbalzarini (2001)]. The final result to one significant decimal is thus: $D_B = 2.8$.

Figure 5.11 finally shows the plots of $\log \eta$ vs. $\log (1/\sigma)$ for all samples. The red crosses mark the measurement points of the reduction steps and the dotted lines the least squares fit. The corresponding numerical values of the box counting dimension are given along with the captions of the sub-figures. It can be seen that the crosses lie on the lines pretty well such that the power law assumption seems to be valid. The box counting dimension thus is a meaningful dimension for the ER and a good estimate of Hausdorff's dimension.



Figure 5.11: $\log \eta$ vs. $\log(1/\sigma)$ (red crosses) for all ER samples and their least squares fits (black dotted). The numerical value of the box counting dimension is given in the caption of each sub-figure.

Chapter 6

Diffusion on fractal sets

6.1 Preliminaries and definitions

In [Sbalzarini (2001)] it was suspected that the fractal dimension of the ER could be a defining geometrical property in the sense that a diffusion process on a model geometry of equal dimension would exhibit "the same" behavior. The quotation marks are there due to the fact that not point-wise identity is required but rather equal time behavior of diffusion, since this is the quantity that is actually measured in a FRAP curve. Model geometry and "real" ER should therefore yield the same FRAP curve when starting from the same initial condition.

In mathematical terms this statement means that two equivalent (in the sense of section 4.2) fractals are suspected to exhibit the same time constant of the solution of the diffusion equation or, in other words, the same eigenvalues of the Laplacian found as exponents in their heat kernel.

To investigate this assertion, some knowledge of the theory of partial differential equations on fractal sets is needed. The PDE to be considered is the heat equation:

$$\frac{\partial u}{\partial t} = D\nabla^2 u$$

which also governs diffusion. Moreover, one must define which dimension the above statement refers to as there is a multitude of different definitions for fractal dimensions. Here, Hausdorff's dimension D_H as defined in section 4.2 is used since a wealth of literature exists for it and its corresponding measure.

All fractal dimensions quantify some sort of "space filling ratio" for the geometrical object in question. For a surface in \mathbb{R}^3 it gives an intuitive measure for how much of the volume is filled up by the surface. It is therefore intuitive to assume that this dimension determines diffusion properties of the surface as the diffusion flux increases with increasing surface area per volume. A very porous hot cube with a rugged surface looses its heat faster than a solid cube of same size with smooth surface. Equality of the fractal dimension in general and Hausdorff's dimension in particular is therefore a necessary condition for the equality of time constants. But is it also sufficient?

To address this question, two preliminary definitions are needed. Fractal objects can be subdivided into two main classes. Imagine a random walk on the fractal. If the set of possible destination points for every step of Brownian motion is finite, the fractal is called *finitely ramified*. If the random walker can choose among an infinite number of places to visit within the next step, it is called *infinitely ramified*. These loose definitions can be made mathematically rigorous as follows:

Definition 6.1 (Finitely ramified). A set for which the number of possible different paths between any two points of the set is finite or countable is called *finitely ramified*.

Definition 6.2 (Infinitely ramified). A set for which the number of possible different paths between any two points of the set is infinite is called *infinitely ramified*

6.2 Sufficiency of Hausdorff's dimension

Consider two different fractal sets of equal Hausdorff dimension, e.g. Koch's curve (see figure 6.1) and the cartesian product of two Cantor sets (figure 6.2) that both have $D_H = \log 4/\log 3$. Looking at the two sets, it is hard to believe that diffusion processes on them should have equal time constants. Already the fact that Koch's curve is completely connected whereas the 2D Cantor set consists of individual disconnected areas gives rise to the suspicion that equality of Hausdorff's dimension may not be sufficient.



Figure 6.1: Koch curve at generation 2



Figure 6.2: 2D Cantor set at generation 2

To disprove that Hausdorff's dimension is sufficient, it is enough to find a single counter example. For reasons of simplicity, the following will be shown for finitely ramified fractals. The results obtained can then be generalized to the more complex class of infinitely ramified fractals that the ER belongs to. The Sierpinski gasket has good enough regularity and connectivity properties to allow reasonable progress. Therefore, it will serve as example.

First, the definition of Brownian motion on the Sierpinski gasket is considered, which then leads to the heat kernel of the diffusion equation, the Laplacian and its eigenvalues. The following arguments closely follow the ones in [Falconer (1997)].

6.2.1 Brownian motion on the Sierpinski gasket

First, some notation needs to be introduced which is best done by recalling the definition of Brownian motion on \mathbb{R} . One way to define Brownian motion is as a limit of suitably scaled random walks. Let $X_k(t)$ be the random walk on the real set $\{j2^{-k} : j \in \mathbb{N}_0^+\}$ starting with $X_k(0) = 0$ and taking steps at time intervals of 4^{-k} . Given the position $X_k(m4^{-k})$ of the walker at time $m4^{-k}$, his position $X_k((m+1)4^{-k})$ at time $(m+1)4^{-k}$ is equally likely to be $X_k(m4^{-k}) - 2^{-k}$ or $X_k(m4^{-k}) + 2^{-k}$. It may be shown that as $k \to \infty$ (infinitely small steps in infinitely small time intervals; cf. also [Sbalzarini (2001)], p. 21), the sequence of random

walks $X_k(t)$ converges to a continuous process X(t) on \mathbb{R} named one-dimensional Brownian motion. Notice that scaling the base of the time scale as two times the base of the length scale is essential for convergence to a non-degenerate process. Thus if k is sufficiently large, $X_k(t)$ and X(t) look very similar on all but the finest scales and the increments of Brownian motion, $X(t + \delta t) - X(t)$, are normally distributed with mean 0 and variance $2\delta t$ for all t and $\delta t > 0$. Typically, the motion travels a distance of $\sqrt{\delta t}$ in a time interval of duration δt . Moreover, Brownian motion has independent increments, that is, there is no historical memory of the path¹. More generally, Brownian motion on \mathbb{R}^n may be constructed as the limit of scaled random walks on n-dimensional cubic lattices.

We now attempt to mimic this construction of Brownian motion on the Sierpinski gasket E following the findings of [Barlow & Perkins (1988)]. The Sierpinski gasket is a fractal of Hausdorff dimension $D_H = \log 3/\log 2 \approx 1.585$ embedded in the Euclidean \mathbb{R}^2 (see figure 6.3). To avoid the need of regarding the three corner vertices of the bounded Sierpinski gasket as exceptional, consider the *extended Sierpinski gasket*, thus E extends outward to infinity using self-similarity as shown in figure 6.4. There is a natural sequence of Euclidean graphs E_0, E_1, \ldots that approximate the extended Sierpinski gasket (see figures 6.4 to 6.6). These graphs are called *pre-Sierpinski gaskets* and it holds: $E_0 \subset E_1 \subset E_2 \subset \ldots$ and $E = \bigcup_{k=0}^{\infty} E_k$.



Waclaw Sierpinski, 1882-1969



Figure 6.3: Sierpinski gasket E



Figure 6.5: Pre-Sierpinski gasket of order 1



Figure 6.4: Pre-Sierpinski gasket of order 0



Figure 6.6: Pre-Sierpinski gasket of order 2

¹Such a random process is called a *Markov chain*.

Let V_k be the set of vertices of E_k . The graph of E_k has edges of length 2^{-k} and each vertex is adjacent to four others. For k = 0, 1, 2, ... the vertices of V_{k+1} are obtained by augmenting V_k by additional vertices at the midpoints of the edges of E_k , with appropriate additional edges added to form E_{k+1} .

Now define random walks $X_k(t)$ on V_k by traveling along the edges of E_k , taking steps at time intervals of α_k (where α_k is to be specified) and starting at $X_k(0) = 0$. Thus if $X_k(m\alpha_k)$ is the vertex of V_k occupied by the random walker at time $m\alpha_k$, then $X_k((m+1)\alpha_k)$ is one of the four vertices of V_k adjacent to $X_k(m\alpha_k)$ in E_k chosen with equal probability $\frac{1}{4}$, independent of all previous steps.

For $k \ge 1$, this random walk $X_k(t)$ on E_k induces a random walk on E_{k-1}^2 . By the symmetry of E_k there is equal probability of moving to each of the four adjacent vertices of V_{k-1} , so this induced random walk is just $X_{k-1}(t)$ undertaken with steps of varying time interval. For the random walks $X_k(t)$ to have a chance of converging to a reasonable limiting process, we should ideally choose the time intervals α_k so that, for each k, the time for $X_k(t)$ to move from a vertex of V_{k-1} to a neighboring vertex of V_{k-1} is α_{k-1} . This can however only be achieved on average by ensuring that the expected time of such a step is α_{k-1} .



Figure 6.7: Situation of the random walk on E_k from x

Consider a portion of E_k near vertex x. By the symmetry of E_k , this is always equivalent to that shown in figure 6.7, where $A = \{a_1, a_2, a_3, a_4\}$ and b_1, b_2 and c are as indicated. Write $\mathsf{E}(p, A)$ for the expected number of steps in a random walk on E_k to get from a point $p \in V_k$ to any point of A. Given that the random walker starts at x, symmetry allows us to assume without loss of generality that the first step is to b_1 when determining $\mathsf{E}(x, A)$. Examining the probabilities of possible steps from b_1, b_2 and c, one finds:

$$E(x, A) = 1 + E(b_1, A)$$

$$E(b_2, A) = E(b_1, A) = 1 + \frac{1}{4}E(x, A) + \frac{1}{4}E(b_2, A) + \frac{1}{4}E(c, A) + \frac{1}{4} \times 0$$

$$E(c, A) = 1 + \frac{1}{2} \times 0 + \frac{1}{4}E(b_1, A) + \frac{1}{4}E(b_2, A)$$

Solving these equations for $\mathsf{E}(x, A)$ gives that $\mathsf{E}(x, A) = 5$. Thus if the random walk $X_k(t)$ on E_k is undertaken with time intervals α_k between each step, the induced random walk on E_{k-1} will have a mean time interval of $5\alpha_k$ between each step.

²Simply note the sequence of vertices of V_{k-1} visited by $X_k(t)$ ignoring consecutive occurrences of the same vertex and regard these as the sequence of vertices visited by a random walk on E_{k-1} .

Whilst a coefficient of 4 might be expected since each vertex is adjacent to four others, the topology of the Sierpinski gasket makes 5 the right number. To take a non-degenerate limit as $k \to \infty$, the random walk on E_{k-1} induced by $X_k(t)$ must, at large scales, be close to the random walk X_{k-1} . To achieve this we need $\alpha_{k-1} = 5\alpha_k$ for each k.

Hence to ensure scaling compatibility, we set $\alpha_k = 5^{-k}$ for k = 0, 1, 2, ... Then it maybe shown, analogously to standard Brownian motion on \mathbb{R} or \mathbb{R}^n , that the sequence of random walks $X_k(t)$ converges for $k \to \infty$ to a random process X(t)which we call Brownian motion on the extended Sierpinski gasket E.

While the basic properties of Brownian motion on \mathbb{R} are mirrored in Brownian motion on E, the exponents are different. As $X_k(t)$ (and thus in the limit X(t)) takes an average time of $5\alpha_k = 5^{-k+1}$ to move between adjacent vertices of V_{k-1} which are distance 2^{-k+1} apart, the Brownian motion X(t) typically moves a distance of $\delta t^{\log 2/\log 5}$ in a time interval δt . This should be compared to $\delta t^{1/2}$ for standard Brownian motion on \mathbb{R}^2 . This leads to defining the dimension of the walk for the extended Sierpinski gasket as $d_w = \log 5/\log 2 \approx 2.322$. It is at least plausible that the mean square of the increments satisfies:

$$\mathsf{E}\left(\left|X(t+\delta t) - X(t)\right|^{2}\right) \asymp \delta t^{2/d_{w}}$$

for $\delta t > 0$. \asymp indicates that this strictly is only asymptotically true, an inevitable consequence of the fractality of the domain. Again, this should be compare to standard Brownian motion on \mathbb{R}^2 with $\mathsf{E}\left(|X(t+\delta t)-X(t)|^2\right) = \delta t$.

Moreover it can be shown (see [Barlow & Perkins (1988)]) that the path of Brownian motion on E is Hölder continuous of order γ , thus:

$$|X(t_1) - X(t_2)| \leq c |t_1 - t_2|^{\gamma}$$

for all $\gamma < 1/d_w$ and $0 \leq t_1 \leq t_2 \leq T$ where the constant *c* depends on γ and *T*. The Hausdorff dimension of the path of Brownian motion is $\log 3/\log 2$ which is equal to the Hausdorff dimension of *E* itself, so the path "fills" the set *E* completely.

6.2.2 Heat kernel and transition density

Having now some knowledge of the basic process of Brownian motion and the definition of the dimension of the walk, the *transition density* $p_t(x, y)$ can be derived. It determines the probability density for position $y \in E$ to be reached after time t by Brownian motion starting at $x \in E$, thus for any measurable set A:

$$P(X(t+\delta t) \in A | X(t) = x) = \int_A p_{\delta t}(x, y) \, d\mu(y)$$

where μ is the restriction of the D_H -dimensional Hausdorff measure to E (thus μ is the natural locally finite measure on E). A knowledge of $p_t(x, y)$ allows the statistics of the motion to be studied. In particular, the transition density is the stochastic analogue to the heat kernel. Using a procedure similar to the one in [Sbalzarini (2001)], pp. 17-21, [Barlow & Perkins (1988)] derived bounds for the transition density by careful analysis of the underlying random walk:

$$c_{1}t^{-D_{H}/d_{w}}\exp\left\{-c_{2}\left(|x-y|t^{-1/d_{w}}\right)^{d_{w}/(d_{w}-1)}\right\} \leq p_{t}(x,y)$$
$$\leq c_{3}t^{-D_{H}/d_{w}}\exp\left\{-c_{4}\left(|x-y|t^{-1/d_{w}}\right)^{d_{w}/(d_{w}-1)}\right\}$$
(6.1)

for certain constants c_1 , c_2 , c_3 , $c_4 > 0$. For standard Brownian motion on \mathbb{R}^n the dimension of the walk is $d_w = 2$ and $\mu = \mathcal{L}^n$ (the Lebesgue measure, which is equal to the *n*-dimensional "volume" for any Borel set) and the transition density is given by:

$$p_t(x,y) = (4\pi t)^{-n/2} \exp\left\{-\left|x-y\right|^2/(4t)\right\}$$
(6.2)

which is the familiar Gaussian kernel of diffusion on an infinite domain. As illustrated in [Sbalzarini (2001)], Brownian motion is intimately connected with solutions of the heat equation. Diffusion of heat on \mathbb{R}^n may be thought of as the aggregate effect of a large number of "heat particles" following independent Brownian paths. Let ν be the heat distribution on \mathbb{R}^n at time t = 0. Then, the temperature at point x and time t is given by:

$$u(x,t) = \int_{\mathbb{R}^n} p_t(x,y) \, d\nu(y) \tag{6.3}$$

where $p_t(x, y)$ is the standard transition density or Green's function (i.e. the heat kernel evaluated at time t) given by equation 6.2. It may be checked by differentiation that

$$\frac{\partial p_t}{\partial t} = \nabla_x^2 p_t$$

as required for Green's function. Thus 6.3 satisfies the heat equation

$$\frac{\partial u}{\partial t} = \nabla^2 u$$

on \mathbb{R}^n with $\int_A u(x,t) dx \to \nu(A)$ as $t \to 0$. In a similar way, Brownian motion on the extended Sierpinski gasket E may be regarded as modeling diffusion on E. Thus an initial heat distribution ν on E would yield the temperature distribution at time t but with $p_t(x,y)$ now being the transition density given by 6.1.

6.2.3 The Laplacian on the Sierpinski gasket

To obtain a meaningful analogue of the heat equation on E, we must define the Laplacian ∇^2 on E. Again, we start with a discrete approximation on the pre-Sierpinski gaskets E_k , which will in the limit give a continuous operator on E. Recall that the Laplacian on \mathbb{R} is the limit of differences:

$$\frac{d^2f}{dx^2} = \lim_{h \to 0} h^{-2} \left[(f(x+h) - f(x)) + (f(x-h) - f(x)) \right]$$
$$= \lim_{h \to 0} h^{-2} \sum_{y=x\pm h} (f(y) - f(x))$$

For a continuous $f : E \to \mathbb{R}$ we use discrete approximations via the geometric graphs E_k (see figures 6.4 to 6.6). Writing C(E) for the set of all continuous functions on E, [Falconer (1997)] defines $\nabla^2 f \in C(E)$ by the requirement that, for every bounded set A:

$$\lim_{k \to \infty} \sup_{x \in A \cap V_k} \left| 5^k \cdot \sum_{w \in V_k(x)} \left(f(w) - f(x) \right) - \nabla^2 f(x) \right| = 0$$
(6.4)

where $V_k(x)$ comprises the four vertices of V_k adjacent to x and the number "5" is exactly what is needed for this definition to be meaningful. This is a consequence of the dimension of the walk of E being $\log 5/\log 2$. We are now able to show that the transition density function on E indeed satisfies the heat equation. For given k, let $x, y \in V_k$. On the assumption that Brownian motion on E starting at x at time 0 reaches one of the four adjacent vertices of V_k at time $\delta t = 5^{-k}$ (we now know that this is true on average), consideration of transition densities to y gives:

$$p_{t+\delta t}(x,y) \simeq \sum_{w \in V_k(x)} \frac{1}{4} \cdot p_t(w,y)$$

where $x_1, \ldots, x_4 \in V_k(x)$ are adjacent to x. This is true on average, due to the symmetry of E. Thus by simple arithmetic:

$$(p_{t+\delta t}(x,y) - p_t(x,y)) / \delta t \simeq \frac{1}{4} \cdot 5^k \cdot \sum_{w \in V_k(x)} (p_t(w,y) - p_t(x,y))$$

Letting $\delta t \to 0$ and inserting this into equation 6.4 gives:

$$\frac{\partial p_t}{\partial t}(x,y) = \frac{1}{4} \nabla^2 p_t(x,y)$$

Moreover, $p_t(x, \cdot)$ is concentrated around x when t is small since X(t) will not have moved far. Therefore:

$$\int_{A} u(x,t) \, d\mu(x) \to \nu(A) \qquad \text{ as } \quad t \to 0$$

The last two equations mean that the transition density p_t meets all the requirements for a Green's function (see e.g. [Sbalzarini (2001)]) and therefore u, as given by equation 6.3 with p_t now being the transition density for Brownian motion on E, satisfies the diffusion equation on the Sierpinski gasket E:

$$\frac{\partial u}{\partial t} = \frac{1}{4} \nabla^2 u$$

Also notice the changed apparent diffusivity as compared to the Euclidean case. With considerable effort, these arguments can be made more rigorous (see e.g. [Barlow & Perkins (1988)]).

6.2.4 Eigenvalues of the Laplacian

We move on to the related problem of finding the eigenvalues of the Laplacian on a fractal domain. Here we need the domain to be bounded, so from now on we take E to be the non-extended usual Sierpinski gasket and adapt the previous notation in an obvious way to the bounded setting. Thus E_k is the graph with a finite vertex set V_k and edges of length 2^{-k} that approximates E. The definition of the Laplacian given by equation 6.4 is slightly modified to require $\nabla^2 f \in C(E)$ to satisfy:

$$\lim_{k \to \infty} \sup_{x \in V_k \setminus V_0} \left| 5^k \cdot \sum_{w \in V_k(x)} \left(f(w) - f(x) \right) - \nabla^2 f(x) \right| = 0$$

where $V_k(x)$ is the set of vertices in V_k adjacent to x, other than the vertices of V_0 . We are interested in eigenfunctions of the Neumann problem in this context, that is, for functions with vanishing normal derivative on V_0 , the three corners of E. The eigenvalues of Helmholtz's equation

$$\nabla^2 u + \lambda u = 0$$
 with $u \in C(E)$ and $\nabla u = 0$ for $x \in V_0$ (6.5)

may be shown to be real and non-negative. We now seek estimates for the *eigenvalue* distribution function:

$$N(\lambda) := \#\{k : \lambda_k \leqslant \lambda\} \tag{6.6}$$

where λ_k are the eigenvalues of problem 6.5. Define the spectral dimension as:

$$d_s = \frac{2D_H}{d_w}$$

where D_H is Hausdorff's dimension and d_w is the dimension of the walk as defined in subsection 6.2.1. Whereas d_w indicates the scaling behavior of the rate of diffusion through E, d_s is defined in terms of the "density of states". d_s therefore can be interpreted as the asymptotic frequency of the large eigenvalues of the Laplacian on a bounded region E, which is related to is harmonic composition (hence the name *spectral* dimension). [Barlow & Perkins (1988)] have shown that the eigenvalue distribution function satisfies:

$$N(\lambda) \simeq \lambda^{d_s/2}$$

This should be compared to Weyl's theorem³ for bounded open domains in \mathbb{R}^n where d_s is replaced by n. Again, one notices the fact that the basic laws remain unchanged when going from Euclidean to fractal sets but the exponents change.

Much more precise information may be obtained on the asymptotics of $N(\lambda)$ by applying the renewal theorem (see e.g. [Falconer (1997)], section 7.2). This theorem gives a periodic positive function $p(\cdot)$ with period log 5 such that:

$$N(\lambda) \sim p(\log \lambda) \lambda^{d_s/2}$$

6.2.5 Extension to infinitely ramified fractals

Up to this point, we have shown that Hausdorff's dimension is not a sufficient parameter to describe the time constants of diffusion on finitely ramified fractals. The Sierpinski gasket has been used as a counter example to disprove the proposition made in [Sbalzarini (2001)]. The ER however belongs to the class of infinitely ramified fractals.

The Sierpinski carpet as depicted in figure 6.8 is a well suited example of this class and will be taken as a model here. Its Hausdorff dimension is $D_H = \log 8 / \log 3 \approx$ 1.893. As shown in [Barlow & Bass (1992)], Brownian motion on the Sierpinski carpet is a strong Markov process⁴ with continuous paths and state space. The infinitesimal generator of the limiting process is again called a "Laplacian". Using similar techniques as in the previous subsections, the fundamental solution of the diffusion equation on the Sierpinski carpet can be derived as the transition density of the underlying Brownian motion. [Barlow & Bass (1992)] found:

$$c_{1}t^{-D_{H}/d_{w}}\exp\left\{-c_{2}\left(|x-y|t^{-1/d_{w}}\right)^{d_{w}/(d_{w}-1)}\right\} \leqslant p_{t}(x,y)$$
$$\leqslant c_{3}t^{-D_{H}/d_{w}}\exp\left\{-c_{4}\left(|x-y|t^{-1/d_{w}}\right)^{d_{w}/(d_{w}-1)}\right\}$$

for certain constants c_1, \ldots, c_4 . This is exactly the same as equation 6.1 for the Sierpinski gasket. The only difference is that in the present case, the value of d_w

³Weyl's theorem states that if $\partial\Omega$ is sufficiently smooth, the eigenvalue distribution function on Ω is $N(\lambda) \simeq c_n \mathcal{L}^n(\Omega) \lambda^{n/2}$ where \mathcal{L}^n means the *n*-dimensional "volume" and $c_n = (2\pi)^{-n} \mathcal{L}^n(B)$ with *B* being the unit ball in \mathbb{R}^n .

 $^{^{4}}$ meaning that all steps are independent and there is no historical memory on the paths



Figure 6.8: Pre-Sierpinski carpet of order 3

is unknown. For infinitely ramified fractals, it is not possible to get numerical values for the dimension of the walk since Brownian motion is now a continuous process. All that can be done is to derive limiting resistances for diffusion (see e.g. [Barlow & Bass (1992)]).

Nevertheless, the identity of the two transition densities shows that the solution of the diffusion equation is topologically equivalent to the one on finitely ramified fractals but with possibly different numerical values for some constants. The basic finding for finitely ramified fractals that Hausdorff's dimension is generally not sufficient to capture the time scales of the solution is therefore true for infinitely ramified fractals as well. This can be intuitively seen if one imagines an infinitely ramified fractal as the limiting case of a partially space-filling curve. Since the curve itself is finitely ramified, it is natural that its infinitely ramified limiting case will inherit the topological properties of the diffusion solution.

6.2.6 Extension to anisotropic diffusion

As real diffusion in the ER might be anisotropic, this case deserves some attention as well. As shown by [Barlow, Hattori, Hattori & Watanabe (1997)], the most interesting aspects of the behavior of diffusion (e.g. the spectral dimensions) are embodied in the asymptotic behaviors of *effective resistances*. These resistances are defined as the H_1 norm of a potential which in turn is the solution to a Laplace equation with appropriate Neumann or Dirichlet boundary conditions. It is thus at least plausible that resistance and diffusion are closely related.

Let $R_n^x(r)$ and $R_n^y(r)$ be the effective resistances of the pre-Sierpinski carpet at the n^{th} stage of its construction in x and y direction, respectively. The anisotropy of the "material" it is made of is parametrized by the ratio of resistances for a unit square: $r = R_0^y/R_0^x$. [Barlow, Hattori, Hattori & Watanabe (1997)] have proven that for sufficiently large n, the ratio $R_n^y(r)/R_n^x(r)$ is bounded by a positive constant independent of r. Furthermore, the ratio decays exponentially fast if $r \gg 1$. This means that for the true Sierpinski carpet with $n \to \infty$ and large ratios of anisotropy, isotropic conditions are weakly restored. Furthermore it is proven that found for the isotropic case. Therefore, the topological behavior of the diffusion solution is the same and all deviations are bounded by a positive constant. The role of the spectral dimension is obvious as it is embodied in the effective resistances. What has

been said in the previous sections remains therefore valid (at least asymptotically for large r and n) for the anisotropic case as well.

6.2.7 Diffusion on domains with fractal boundary

The last missing piece in this survey of diffusion on fractal sets is the connection to Euclidean domains. Recall that the ER is not a true fractal in reality but there is a lower bound to the geometrical scales, given by the molecular structure of the ER membrane. At larger scales though, the ER can be *modeled* using *concepts* from fractal theory. The limiting case between the two scales corresponds to a Euclidean domain with fractal boundary.

Then the question is how and whether diffusion *inside* the domain is influenced by the fractality of its boundary. Again we will consider the eigenvalues of the Laplacian as characteristic for the time scales of the solution. Let $D \subset \mathbb{R}^n$ $(n \ge 1)$ be a bounded open region with boundary ∂D . Here we do not need D to be connected (although the ER is). The eigenvalue problem is again given by Helmholtz's equation:

$$\nabla^2 u + \lambda u = 0 \qquad \text{in} \qquad D$$

with Neumann boundary condition $\nabla u(x) \cdot n = 0$ for $x \in \partial D$. The eigenvalues are those λ_k for which there is a non-trivial solution. Again we are interested in the eigenvalue distribution function as defined by equation 6.6. In particular how its behavior reflects the nature of the domain boundary ∂D . A classical result of Weyl states that is ∂D is sufficiently smooth then:

$$N(\lambda) \sim c_n \mathcal{L}^n(D) \lambda^{n/2}$$

as $n \to \infty$, where $c_n = (2\pi)^{-n} \mathcal{L}^n(B)$, B is the unit ball in \mathbb{R}^n and \mathcal{L}^n is the *n*-dimensional Lebesgue measure (*n*-dimensional "volume"). This asymptotic expansion can be continued as:

$$N(\lambda) = c_n \mathcal{L}^n(D) \lambda^{n/2} + b_n \mathcal{L}^{n-1}(\partial D) \lambda^{(n-1)/2} + o\left(\lambda^{(n-1)/2}\right)$$

for a constant b_n depending only on n. Thus the "surface area" of the boundary ∂D determines the second term in the expansion of $N(\lambda)$. Notice that the exponent (n-1)/2 is half the dimension of the boundary. For fractal boundaries, it has been shown by [Falconer (1997)], that there is a connection between the fractal dimension of the boundary and the second term in the expansion of $N(\lambda)$. One can thus "hear" the dimension of an oscillating fractal membrane. It can be shown that $\lambda^{s/2}$ is an upper bound for the second term where s is the *interior Minkowski dimension*, a one-sided variant of the box counting dimension as defined in section 4.3.2.

6.3 Conclusions

Summarizing the previous sections it can be said that a model geometry for the ER that has the same box counting or Hausdorff dimension will not necessarily exhibit the same FRAP behavior. Equality in Hausdorff's dimension is **necessary** but not sufficient for the model to be valid. A necessary and sufficient set of equalities would comprise Hausdorff's dimension and the spectral dimension (see section 6.2.4), since only then, the two objects will have the same dimension of the walk as defined in section 6.2.1. It is this dimension of the walk that captures the properties of Brownian motion and thus diffusion on a given geometry. For Euclidean objects it is always equal to 2, for fractals it will be different. The Sierpinski gasket for example has been shown to have $d_w = 2.322...$ The greater value for the fractal indicates that diffusion on it is slower. The expected mean square increment is $\delta t^{2/d_w}$. This makes sense since the topological structure of the fractal set constrains the Brownian motion. On the other hand diffusion through a fractal boundary is faster than through a Euclidean one since its effective surface is larger (e.g. porous media). Moreover, the exponents in most physical laws, the coefficients in the definition of the Laplacian and its eigenvalue distribution function as well as the apparent diffusivity will be different for fractal and Euclidean domains. Only if all of this is taken into account, models that capture the basic integral time behavior can be derived. Unfortunately, the spectral dimension of any given object cannot be easily measured. Indeed it is still unknown how to determine it (or the dimension of the walk) for infinitely ramified fractals. This leads to the conclusion that the idea of a fractal model geometry for the ER as proposed in [Sbalzarini (2001)] is impracticable. Instead, the fact that the dimension of the walk completely captures the influence of geometry (or the restriction of Brownian motion to it) on the time behavior of diffusion will be used to directly infer FRAP data models that appropriately account for the complexity of the ER geometry. The results of this chapter show that this is possible using d_w as a single numerical parameter.

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Chapter 7

Simulation techniques

7.1 Random walk

For all random walk simulations of diffusion, the 3D variant of the code developed and described in [Sbalzarini (2001)] is used. However, the following improvements have been made in the meantime:

- 1. The geometry description input files now contain a file header telling the program whether it is an old-style ***.surf** file (cf. [Sbalzarini (2001)]) or an OpenInventor 3D file according to [SGI (1992a)] and [SGI (1992b)]. Therefore the input file read routines have been changed to implement this.
- 2. The global parameter file bdiff.dat has been changed to a more flexible format now allowing comments and directives of the form KEYWORD = value in arbitrary order (cf. appendix C.5.2). The new subroutine ReadParams (given in appendix C.5.4.5) implements this. For all parameters not being explicitly specified in the input file, default values are now set in Defaults.f90 (similar to the one in appendix C.5.4.3).
- 3. The numerical value of π is computed to machine precision as $\pi = 4 \arctan(1)$ rather than being hard-coded to a certain (fixed) precision.
- 4. Position and size of the bleached box are now read from the parameter file bdiff.dat as bleachbox = <Xmin>, <Ymin>, <Xmax>, <Ymax> rather than being hard-coded in the program.
- 5. A global tolerance parameter TOL has been added to make real value comparisons more robust against round-off errors. Two real numbers are now considered equal if their difference is less than TOL.
- 6. The subroutine point_in_domain no longer needs the triangle's normals but explicitly counts the number of triangles intersected along a ray in x-direction starting at the point in question. See algorithm 7.8 for further details on this.
- 7. Finally, the directions of the random walk steps are no longer chosen on the unit sphere around the current position but on the unit semi-sphere. Consequently, the step length is no longer of fixed sign but is now also allowed to be negative to make the whole space reachable again. Thus, for space particles, the algorithm now proceeds as follows:

Algorithm 7.1 (Space random walk).

Step 1: Choose a random point on the unit semi-sphere centered at the current position $\{(r, \varphi, \vartheta) : r = 1, 0 \le \varphi < \pi, -\pi/2 \le \vartheta \le \pi/2\}$ according to:

$$\varphi = \pi \cdot \mathcal{U}(0,1)$$
 $\vartheta = \arcsin\left(2 \cdot \mathcal{U}(0,1) - 1\right) + \frac{\pi}{2}$

where $\mathcal{U}(0,1)$ is a uniformly distributed random number between 0 and 1.

- Step 2: Choose the step variance according to $\sigma^2 = 2nD\delta t$ where D is the diffusion constant (interpolated as described in [Sbalzarini (2001)]), δt the simulation time step size and n the number of spatial dimensions (thus n = 3).
- Step 3: Now the random walk step is given by:

$$s = \mathcal{N}(0, \sigma^2) \cdot \left[\begin{array}{c} \sin \vartheta \cos \varphi \\ \sin \vartheta \sin \varphi \\ \cos \vartheta \end{array} \right]$$

where $\mathcal{N}(0, \sigma^2)$ means a Gaussian (normal) deviate with zero mean and variance σ^2 .

Step 4: Advance to the new position: $x_{n+1} = x_n + s$

For surface particles the corresponding algorithm becomes:

Algorithm 7.2 (Surface random walk).

- Step 1: Choose a random point on the unit semi-circle around the current position $\{(r, \vartheta) : r = 1, 0 \leq \vartheta < \pi\}$ according to: $\vartheta = \pi \cdot \mathcal{U}(0, 1)$
- Step 2: The step variance is again given by: $\sigma^2 = 2nD\delta t$ with *n* now being 2 and *D* the surface diffusion coefficient on the tube in question.
- Step 3: The polar angle (φ) and axial coordinate (ζ) of the new position are now given by:

$$\varphi_{n+1} = \varphi_n + \frac{1}{R} \mathcal{N}(0, \sigma^2) \cos \vartheta \mod 2\pi$$
$$\zeta_{n+1} = \zeta_n + \mathcal{N}(0, \sigma^2) \sin \vartheta$$

where R is the radius of the current tube.

Besides these changes, the program code as well as its input and output facilities are the same as in [Sbalzarini (2001)]. A complete reference manual as well as a printed version of the source code is included in Appendix C of [Sbalzarini (2001)] and will not be reproduced in this report.

7.2 Particle Strength Exchange

In addition to the random walk code, the method of particle strength exchange (PSE) is used. It is a Lagrangian particle method suited to solve diffusion or the diffusive (viscous) part of other equations such as the Navier-Stokes equations in complex geometries. For this work, the PSE method is preferred to other alternatives such as embedded boundary methods ([McCorquodale, Colella & Johansen (2001)]) or the diffusion velocity method ([Beaudoin, Huberson & Rivoalen (2001)]). Advantages over the random walk method are that it converges much faster (as will be shown in section 8.4) and yields smooth solutions. Being a particle method, it is suited for simulations in arbitrarily shaped domains, just as the random walk is. With prospect to future work, the fully anisotropic implementation of [Zimmermann, Koumoutsakos & Kinzelbach (2001)] is used in this project.

7.2.1 The principles of particle methods

The method of particle strength exchange (PSE) is part of a larger family of particle methods widely used in various areas of computational physics such as fluid mechanics, electrostatics, plasma physics, etc. A complete introduction and survey of particle methods can for example be found in [Hockney & Eastwood (1988)], for applications in fluid mechanics [Cottet & Koumoutsakos (2000)] is recommended.



Figure 7.1: Ways to solve a differential equation numerically.

The basic idea of all of these methods consists of finding an integral approximation of the differential operator under consideration, which is then solved by numerical quadrature using the irregularly distributed particle locations as quadrature points. Figure 7.1 depicts the general philosophy. Starting from the governing differential equation (ODE or PDE) of the problem, two ways of solution exist. One is to discretize the equation using a grid-based method such as finite differences, finite elements or finite volumes. This discrete equation can then be solved numerically. The issues to be concerned about include consistency and accuracy for the discretization step as well as stability and accuracy for the solution step. If either consistency or stability criteria are violated, the discretized equation will not represent the original physical problem any more. This is possible since it is - in principle – a completely new, artificial formula for which we are not a priori guaranteed that is has anything to do with the governing equation we started from. The second way consists of writing the analytical solution of the equation (in integral form using Green's function or some kernel) even though we may not be able to explicitly solve it. This step is exact and no errors and constraints need to be worried about. One then proceeds by numerical integration (sometimes called quadrature) of this integral solution. The only issue of this step is accuracy. In fact, this way of solving the equation is *always stable*, provided that the actual solution of the governing equation is stable itself, and consistency is guaranteed by construction of the method. The drawbacks however are that (i) quadrature converges slower than numerical solutions of the discretized equation and (ii) the resulting system of equations is an N-body problem making it potentially scaling of $\mathcal{O}(N^2)$. It was

this high computational cost that long prevented the use of particle methods in computational science. Fortunately, fast N-body solvers such as cell-list algorithms or multipole expansions which make the methods $\mathcal{O}(N \log N)$ or even $\mathcal{O}(N)$ are available nowadays. See for example [Hockney & Eastwood (1988)] and appendix B of [Cottet & Koumoutsakos (2000)] for detailed descriptions of those fast solvers.

Generally speaking, particle methods are a way of translating the governing PDE of a physical problem to a set of N ODEs where N is the total number of particles involved. These ODEs are called *equations of motion* since they describe the motion of the particles as:

$$\frac{dx_k}{dt} = F(x, t, \omega) \qquad k = 1, \dots, N$$

As we are interested in representing some smooth function (e.g. charge, concentration, probability, etc) by a particle distribution, the particles also carry a physical quantity generally referred to as their *strength* ω . The strength is also allowed to change in time as characterized by a second set of ODEs:

$$\frac{d\omega_k}{dt} = G(x, t, \omega) \qquad k = 1, \dots, N$$

Thus, we define a particle as follows:

Definition 7.1 (Particle). A particle P is a mathematical object described by its location x_p and the strength ω_p it is carrying: $P = (x_p, \omega_p)$.



Figure 7.2: Two particles of strength ω_1 and ω_2 carrying mollifier functions ζ_{ϵ}

In order to be able to reconstruct a smooth function from the particles, they are assigned *mollifier functions* ζ_{ϵ} . A mollifier is a local (but not necessarily compact) function centered at the particle's position (see figure 7.2). It can be thought of as a cloud of charge, mass, etc. that is carried around by the particle and has to meet certain requirements (see [Cottet & Koumoutsakos (2000)] for details). The *core size* ϵ of this mollifier determines the size of the particles. Introducing such finitely sized particles is equivalent to fixing the particle interaction potential such that is goes to zero if the distance between two particles vanishes¹. The smaller the particles are, the more accurate the representation of the smooth function gets.

 $^{^{1}}$ This unphysical fix is needed to prevent the particle interactions from becoming of infinite magnitude. It will however not introduce a large error as the particles will not come that close in real simulations anyway.

For the approximation to converge, it is however required that the distance h between any two particles is always less than their size², thus:

$$\frac{h}{\epsilon} < 1$$

This means that more accurate function approximations can only be obtained by increasing the number of particles when decreasing their size.

7.2.2 The isotropic PSE method

As for diffusion, the governing equation is given by

$$\frac{\partial c}{\partial t} = D\nabla^2 c(x, t) \tag{7.1}$$

the PSE method as introduced by [Degond & Mas-Gallic (1989a)] starts by finding an integral formulation of the Laplacian that allows consistent evaluation on the particle locations. To start with, consider the solution at a location y and expand it into a Taylor series around x:

$$c(y) = c(x) + \sum_{i=1}^{r+1} \frac{1}{i!} (y-x)^i \nabla^i c(x) + O\left(|y-x|^{r+2} \|c\|_{\infty}\right)$$

Subtracting c(x) on both sides, multiplying the whole equation by a *regularized* kernel function η_{ϵ} and integrating over y yields:

$$\int (c(y) - c(x)) \eta_{\epsilon}(y - x) dy = \sum_{i=1}^{r+1} \frac{1}{i!} \int (y - x)^i \nabla^i c(x) \eta_{\epsilon}(y - x) dy$$
$$+ \|c\|_{\infty} O\left(\int |y - x|^{r+2} \eta_{\epsilon}(y - x) dy\right)$$

To get an approximation of the Laplacian, we have to ask the following general requirement for the kernel function η (according to [Degond & Mas-Gallic (1989a)]):

$$\int_{\mathbb{R}^n} \prod_{i=1}^n x_i^{\alpha_i} \eta(x) \, dx = \begin{cases} 0, & \forall \alpha \in \mathbb{N}^n, \ \alpha \neq 2e_i, \ 1 \leq \sum_{i=1}^n \alpha_i \leq r+1 \\ 2, & \text{if } \alpha = 2e_i, \ i \in [1, \dots, n] \end{cases}$$

where *n* is the dimension of the space, *r* is the *order* of the regularized kernel function and $x = (x_1, \ldots, x_n)$ is a position in \mathbb{R}^n . $\alpha = (\alpha_1, \ldots, \alpha_n)$ is an *n*-dimensional index and (e_1, \ldots, e_n) is the canonical base of \mathbb{R}^n . In the 3D case, this requirement may be expressed as:

$$\int x_i x_j \eta(x) \, dx = 2\delta_{ij} \qquad \text{for } i, j = 1, 2, 3 \tag{7.2}$$

$$\int x_1^{i_1} x_2^{i_2} x_3^{i_3} \eta(x) \, dx = 0 \quad \text{if } i_1 + i_2 + i_3 = 1 \text{ or } 3 \leqslant i_1 + i_2 + i_3 \leqslant r + 1 \tag{7.3}$$

$$\int |x|^{r+2} |\eta(x)| \, dx < \infty \tag{7.4}$$

for $i_1, i_2, i_3 \in \mathbb{N}_0^+$. The first condition is to normalize the kernel function. The second one requires all moments up to order r+1 to vanish and the third one is

²This is known as "particles must overlap".

required to have the truncation error bounded. The regularized kernel is derived from this kernel function as follows:

$$\eta_{\epsilon}(x) = \epsilon^{-n} \eta\left(\frac{x}{\epsilon}\right) \qquad , \epsilon > 0$$

where n is the topological dimension of the space we wish to solve the diffusion equation in. Using requirements 7.2 and 7.3, one notices that all terms of the form $\int (y-x)^i \eta_{\epsilon}(y-x) \, dy$ in above series expansion vanish. The only term remaining is:

$$\int (c(y) - c(x)) \eta_{\epsilon}(y - x) dy = \nabla^2 c(x) \int (y - x)^2 \eta_{\epsilon}(y - x) dy$$
$$+ \|c\|_{\infty} O\left(\int |y - x|^{r+2} \eta_{\epsilon}(y - x) dy\right)$$

The term $\nabla^2 c(x)$ can be taken out of the integral as it does not depend on y. After change of variables $z = (y - x)/\epsilon$ in the integrals, one obtains:

$$\epsilon^{-n} \int \left(c(y) - c(x) \right) \eta_{\epsilon}(y - x) \, dy = \nabla^2 c(x) + O\left(\epsilon^r\right)$$

due to the requirement 7.2 for the kernel function η . Thus the integral operator that approximates the Laplacian is found to be:

$$\nabla_{\epsilon}^{2} c(x) = \epsilon^{-n} \int \left(c(y) - c(x) \right) \eta_{\epsilon}(y - x) \, dy \tag{7.5}$$

and the approximation error is $O(\epsilon^r)$ with r being the largest integer for which conditions 7.3 and 7.4 are fulfilled (see [Cottet & Koumoutsakos (2000)] for a rigorous error treatment). Using the particle locations as quadrature points leads to the discrete version of the operator:

$$\nabla_{\epsilon,h}^2 c^h(x_p^h) = \epsilon^{-n} \sum_{q \neq p} (v_q c_q^h - v_p c_p^h) \eta_\epsilon(x_q^h - x_p^h)$$
(7.6)

where v_q and v_p are the particle's volumes such that $v_q c_q^h = c(x_q^h) dy$ is the strength (mass is this context). It is noteworthy that this operator is not the only possibility of discretizing the Laplacian onto particles. Compared to another method (called Fishelov's scheme) it has however the big advantage of being conservative. As above operator implicitly incorporates the conservation of mass, it is considered the better choice for solving the diffusion equation.

The approximation c^h to the continuous concentration c at any location and time can be reconstructed from the strengths $v_p c_p^h$ of the particles using:

$$c^{h}(x,t) = \sum_{p} v_{p}c_{p}^{h}(t)\zeta_{\epsilon}(x-x_{p}^{h})$$

where $\zeta_{\epsilon}(x) = \epsilon^{-n} \zeta(x/\epsilon)$ is the mollifier function that the particles "carry around" (not to be confused with the PSE kernel η_{ϵ}). In the simplest case of point particles it is identical to the Dirac delta distribution: $\zeta(x) = \delta(x)$. The final PSE scheme is now easily obtained by inserting equation 7.6 into equation 7.1:

$$\frac{\partial c_p^h}{\partial t} = D\epsilon^{-n} \sum_{q \neq p} (v_q c_q^h - v_p c_p^h) \eta_\epsilon (x_q^h - x_p^h) \qquad \forall \ p \in [1, \dots, N]$$

This is an N-body problem as for each particle it involves a sum over all other particles. However, since the kernel η_{ϵ} is chosen to be local, only the nearest neighbors of each particle significantly contribute to its sum. The simulation code therefore implements a cell-list algorithm for nearest neighbor search and interactions are only calculated between particles that are closer than a cut-off of 10ϵ . It can also be seen from this equation, that in order to simulate diffusion, the strengths of all the particles change (i.e. they exchange mass) while their locations remain the same, i.e. they do not move. This means that all the geometry and boundary condition handling only needs to be done once when initializing the particles (cf. section 7.2.5). Moreover, it allows regular spacing of the particles inside the computational domain in which case the particle volumes simply become: $v_p = h_1 h_2 h_3$ where $h_{1,2,3}$ are the (constant) inter-particle spacings in all three spatial directions.

The anisotropic advection-dispersion equation

It is however possible to also take convection into account. In this case the particles move as governed by the convection and they exchange strength according to diffusion. It is this appealing physical problem separation character that makes part of the beauty of particle methods. The simulation code used in this work directly implements a solver for the advection-dispersion equation:

$$\frac{\partial c(x,t)}{\partial t} + \nabla (u(x,t) \cdot c(x,t)) - \sigma(x,t) = \nabla (D(x,t) \cdot \nabla c(x,t))$$
(7.7)

where u is a given mean velocity field, c stands for the local concentration and σ is a term to account for sources and sinks. Due to the anisotropy, D now is a symmetric tensor of second rank that contains diffusion and dispersion coefficients.

Using the principles of particle methods, this equation is now transformed into the following set of equations:

$$\frac{dx_p}{dt} = u(x_p, t) \tag{7.8}$$

$$\frac{\partial c(x,t)}{\partial t} = \nabla (D(x,t) \cdot \nabla c(x,t)) + \sigma(x,t)$$
(7.9)

with u(x,t) being a given external mean velocity field which moves the particles along their characteristics. The first equation describes the movement of the particles due to convection, the second equation describes the changes in their strength due to diffusion and is solved using the anisotropic PSE method as described below.

7.2.3 The anisotropic extension of the PSE method

In order to have a PSE simulation that is functionally equivalent to the random walk code, an extended anisotropic version of it is used. In this case, the diffusion coefficient is no longer a scalar but a symmetric tensor of second rank that is allowed to vary in space (in 3D):

$$D(x) = \begin{bmatrix} D_{xx}(x) & D_{xy}(x) & D_{xz}(x) \\ D_{xy}(x) & D_{yy}(x) & D_{yz}(x) \\ D_{xz}(x) & D_{yz}(x) & D_{zz}(x) \end{bmatrix}$$
(7.10)

Similar to the derivations for the isotropic case in section 7.2.2, it has been shown by [Degond & Mas-Gallic (1989b)] by means of Taylor series expansions that the following integral operator Q_{ϵ} is an approximation of the differential diffusion operator:

$$\nabla(D \ \nabla c(x,t)) \approx Q_{\epsilon}(t) \ c(x,t) = \int_{\mathbb{R}^n} \sigma_{\epsilon}(x,y,t) \ [c(y) - c(x)] \ dy \tag{7.11}$$

The regularized kernel $\sigma_{\epsilon} = \epsilon^{-n} \sigma(x/\epsilon)$ again satisfies certain moment conditions. The discretized particle approximation $\overline{Q}^{h}_{\epsilon}$ is obtained by applying a quadrature rule to the integral operator $Q_{\epsilon}(t)$ using the particles as quadrature points:

$$\overline{Q}^{h}_{\epsilon}(t) c^{h}_{k}(t) = \sum_{l} \sigma_{\epsilon}(x_{k}(t), x_{l}(t), t) \left[c_{l}(t) - c_{k}(t)\right] \cdot h^{n}$$
(7.12)

where h is the inter-particle spacing and n is the space dimension. The regularized kernel is defined as:

$$\sigma_{\epsilon}(x_k, x_l, t) = \epsilon^{-n} \sum_{i,j=1}^{n} M_{ij}(x_k, x_l, t) \psi_{ij}^{\epsilon}(x_l - x_k)$$
(7.13)

where ϵ is again the core size of the particles and $M_{ij}(x, y, t)$ is a function of the diffusion tensor D_{ij} . With the spherically symmetric matrix cut-off (smoothing) function:

$$\psi_{ij}^{\epsilon}(x) = \epsilon^{-n} \psi_{ij}\left(\frac{x}{\epsilon}\right)$$

According to [Degond & Mas-Gallic (1989b)], the components of ψ are chosen to be:

$$\psi_{ij} = \epsilon^{-(n+2)}\overline{\Theta}\left(\frac{|x_k - x_l|}{\epsilon}\right) \cdot \sum_{i,j=1}^n (x - y)_i (x - y)_j$$

with a scalar cut-off function $\overline{\Theta}(x)$. Substituting everything into equation 7.13 yields the regularized anisotropic PSE kernel function:

$$\sigma_{\epsilon}(x_k, x_l, t) = \frac{1}{\epsilon^{n+4}} \overline{\Theta}\left(\frac{|x_k - x_l|}{\epsilon}\right) \sum_{i,j=1}^n M_{ij}(x_k, x_l, t)(x_k - x_l)_i(x_k - x_l)_j$$

Introducing the following constraint according to [Degond & Mas-Gallic (1989b)]:

$$\frac{4\pi}{15} \int_0^\infty r^6 \ \overline{\Theta}(r) \, dr \stackrel{!}{=} 1$$

with $r = (x_i - x_j)$, $\overline{\Theta}(r) = a_0 \cdot e^{-\beta r^2}$ and $\beta = 1$, we obtain the following normalized 2^{nd} order cut-off function:

$$\overline{\Theta}\left(\frac{|x_k - x_l|}{\epsilon}\right) = \frac{4}{\pi\sqrt{\pi}} e^{-\frac{(x_k - x_l)^2}{\epsilon^2}}$$
(7.14)

[Degond & Mas-Gallic (1989b)] suggest $M(x_k, x_l, t)$ to be of the form

$$M(x_k, x_l) = \frac{1}{2} (m(x_k) + m(x_l))$$

where

$$m = D - \frac{1}{n+2} Tr(D) \cdot \mathbb{I}$$

with $Tr(\cdot)$ standing for the trace of a matrix, and \mathbb{I} for the identity matrix. In 3D, the diffusion tensor is given by equation 7.10 and thus we get:

$$m = \begin{bmatrix} \frac{1}{5}(4D_{xx} - D_{yy} - D_{zz}) & D_{xy} & D_{xz} \\ D_{xy} & \frac{1}{5}(4D_{yy} - D_{xx} - D_{zz}) & D_{yz} \\ D_{xz} & D_{yz} & \frac{1}{5}(4D_{zz} - D_{yy} - D_{xx}) \end{bmatrix}$$

where all coefficients can be functions of the location x. Finally, using a simple Euler time discretization with time step size δt , the discretized form of equation 7.9 in 3D may be expressed as:

$$c_k^{n+1} = c_k^n + \frac{\delta t \cdot h_1 h_2 h_3}{\epsilon^7} \sum_{l=1}^N (c_l^n - c_k^n) \cdot \frac{4}{\pi \sqrt{\pi}} e^{-\frac{(x_k - x_l)^2}{\epsilon^2}} \left(M_{11} (x_k - x_l)_1^2 + 2 \cdot M_{12} (x_k - x_l)_1 (x_k - x_l)_2 + 2 \cdot M_{13} (x_k - x_l)_1 (x_k - x_l)_3 + M_{22} (x_k - x_l)_2^2 + 2 \cdot M_{23} (x_k - x_l)_2 (x_k - x_l)_3 + M_{33} (x_k - x_l)_3^2 \right)$$

As in the isotropic case, the sum is only to be taken over the nearest neighbors within a distance of 10ϵ due to the local character of the interaction kernel (negative exponential). This is efficiently done using a cell-list algorithm. The algorithm has been implemented in a parallel simulation code using MPI by Dr. Jens Walther and Stephanie Zimmermann (see [Zimmermann, Koumoutsakos & Kinzelbach (2001)] for a reference application). The changes and additions needed to handle finite domains with arbitrary triangulated boundaries and to make the code restartable are however products of the present project.

7.2.4 Boundary condition handling

Since the PSE algorithm as described above strictly only applies to infinite domains, some kind of boundary condition handling needs to be included. The easiest way to do so consists of placing mirror particles in a r_c -neighborhood outside of the simulation domain Ω . The same technique is also used for analytically solving partial differential equations using Green's function. It is exact for domains whith boundaries that can be described as a set of infinite straight planes (e.g. cubes, boxes, etc.). This clearly is not the case for the ER, giving rise to some boundary errors. Since such errors violate the no flux boundary condition, they will change the total mass inside the domain Ω . Checking the conservation of mass therefore is a way to monitor those errors.

The PSE code is extended to handle three different kinds of particles: real particles, mirror particles for the boundary condition and ghost particles for inter-processor communication, each kind is identified by a particle attribute which is positive for real particles, negative for mirror particles and zero for ghost particles. To enforce the boundary condition, the strength of all mirror particles is set to the one of their corresponding real particle every time step directly after the PSE solver.

7.2.5 Geometry processing and initialization

Since no convection is present (we are simulating pure diffusion), the PSE algorithm does not move the particles but only changes their strengths. This is an important difference to the random walk technique and eliminates the need for geometry processing (checking whether a particle is inside or outside the domain) at every time step. Instead, the geometry handling needs to be done only once when initializing the particles. Therefore, a preprocessor code is written that writes an input file for the PSE solver containing all initial particle positions, strengths and attributes. This preprocessor is the only algorithm that needs to know the triangulated surface description of the domain Ω . Once the particles are initialized, the information about the shape of the domain are contained in their positions since no particles with positive attribute will be placed outside of Ω .

All geometry processing and particle initialization is concentrated in the program init_part. The particles are initialized on a regular cartesian lattice placing a

particle at the center of each grid cell³. The complete source code is given in appendix C.5.4. The top-level algorithm proceeds as follows, marginal numbers refer to statement labels in the main program (appendix C.5.4.2):

Algorithm 7.3 (Preprocessing).

- 36 Step 1: Set default values for all variables.
- 110 Step 2: Read triangulation set and problem parameters.
- Step 3: Determine the bounding box of the triangulation $[x_l, x_h] \times [y_l, y_h] \times [z_l, z_h]$ as: $\{x, y, z\}_l = \min_p \{x_p, y_p, z_p\}, \{x, y, z\}_h = \max_p \{x_p, y_p, z_p\}$ where p loops over all triangle vertices in \aleph .
- 231–240 Step 4: Calculate centroids and normals of all triangles.
- 241 Step 5: Set up triangle lists and sort triangles.
- 256-264 Step 6: Check topological and syntactical validity of the triangulation according to the algorithms presented in chapter 3.
- Step 7: Initialize particles on regular lattice by placing a particle at the center of each grid cell that is inside the domain Ω .
- Step 8: Initialize strengths of all particles by setting the ones inside the bleached box to zero, all others to c_0 . Initialize all attributes to any value > 0 (typically a running index).
- 390–394 Step 9: Add mirror particles for boundary condition handling.
 - Step 10: Calculate and output some diagnostics.
- 417–424 Step 11: Write particle file to be read in by the PSE solver.
- 448-459 Step 12: Deallocate memory and terminate.

Some explanations to selected steps are in order. Step 5 is the core part of two cell-list algorithms included for performance reasons. Without any sorting, step 7 would have to check every triangle for each grid cell to determine whether the cell's center is inside the domain or not, making it $\mathcal{O}(NM)$. This is computationally not feasible for any surfaces containing more than just a few dozens of triangles. Therefore, step 5 sets up two different lists of triangles: the first list sorts the triangles into (y, z)-bins B of size $[x_l, x_h] \times b_y \times b_z$:

$$B_{i,j} = \{ (x, y, z) : x_l \leq x \leq x_h, \ (i-1)b_y \leq y - y_l < ib_y, \ (j-1)b_z \leq z - z_l < jb_z \}$$

It proceeds as follows (see appendix C.5.4.9 for the source code, marginal numbers refer to the corresponding statement labels):

Algorithm 7.4 (Make bin lists).

- Step 1: Subdivide the bounding box into $N_y \times N_z$ cylindrical bins of size $b_x = x_h x_l$, $b_y = (y_h - y_l)/N_y$, $b_z = (z_h - z_l)/N_z$.
- Step 2: Associate a triangle list $L_{i,j}(k)$ with each bin (i, j)
- Step 3: Loop over all bins and assign all triangles that intersect the parallelepiped $\{(x, y, z) : x_l \leq x \leq x_h, y_l + (i-1)b_y \leq y < y_l + ib_y, z_l + (j-1)b_z \leq z < z_l + jb_z\}$ to the triangle list $L_{i,j}$ of bin (i, j).

21-27

28-143

 $^{^{3}}$ This staggered arrangement is needed to avoid particle loss due to round-off errors if particles sit exactly on the domain boundaries
Since a triangle typically belongs to several bins, regular lists have to be used instead of linked lists⁴ (see section 7.2.6 for details on the list structure). The code efficiently determines all bins that are intersected by a certain triangle Δ_k using the following algorithm:

Algorithm 7.5 (Sort triangles to bin lists).

- Step 1: For each vertex $P_{1,2,3}$ of Δ_k , determine the index of the bin it is in: $(i, j) = \operatorname{ceiling}((P_y y_l)/b_y, (P_z z_l)/b_z)$, where $\operatorname{ceiling}(f)$ is the closest integer larger or equal to f. If Δ_k is not yet contained in the list $L_{i,j}$, add it: $L_{i,j}(\operatorname{length}(L_{i,j}) + 1) = k$.
- Step 2: Determine the minimum and maximum bin indices for Δ_k : $\{i, j\}_l = \min_{i,j}\{(i, j) : \Delta_k \in L_{i,j}\}, \{i, j\}_h = \max_{i,j}\{(i, j) : \Delta_k \in L_{i,j}\}$ 64-67
- Step 3: Loop over all bins $\{(i, j) : i \in [i_l, i_h], j \in [j_l, j_h]\}$
 - 3.1: If any corner point of the bin's projection onto the (y, z)-plane is inside the triangle's projection, add Δ_k to $L_{i,j}$ if it is not already in this bin's triangle list. Go to step 3.
 - 3.2: If any edge of the triangle's projection onto the (y, z)-plane intersects any edge of the bin's projection, add Δ_k to $L_{i,j}$ if it is not already in this bin's triangle list.

Steps 3.1 and 3.2 use 2×2 linear systems of equations to determine the intersection points, analogous to the ones in chapter 3. Figure 7.3 depicts the geometrical situation, taken from a real run with the erp572 ER sample. The point in question is P = (401.2188..., 387.6003..., 5.2666...). It is marked with a large light-blue asterisk. The corresponding bin this point belongs to has indices (70,38) and is shown as a red box. All 49 triangles that intersect this bin (thus are elements of its triangle list) are drawn in green. The 4 triangles that actually are intersected by the ray starting at P and proceeding parallel to the x-axis are shown in pink, the ray itself and all ray-triangle intersection points are colored blue. Due to the fact that there are 4 intersected triangles, the point P is outside of the domain Ω .

The second list sorts the triangles into (x, y, z)-cells C of size $b_x \times b_y \times b_z$:

$$C_{i,j,k} = \{ (x, y, z) : (i-1)b_x \leq x - x_l < ib_x, \ (j-1)b_y \leq y - y_l < jb_y, \\ (k-1)b_z \leq z - z_l < kb_z \}$$

It is only used for the creation of the boundary condition mirror particles as it simplifies the search for points closer than r_c to the surface and – for such points – the search for the triangle closest to it (to mirror at). The creation of the lists is done in AllocateLL (see appendix C.5.4.10) as:

Algorithm 7.6 (Make cell lists).

Step 1: Determine the number of cells needed: $N_x = \operatorname{int}((x_h - x_l)/r_c), N_y = \operatorname{int}((y_h - y_l)/r_c), N_z = \operatorname{int}((z_h - z_l)/r_c)$ where $\operatorname{int}(f)$ means the closest integer $\leq f$. 35-37

Step 1: Subdivide the bounding box into cubic cells of edge length $b_x = (x_h - x_l)/N_x$, $b_y = (y_h - y_l)/N_y$, $b_z = (z_h - z_l)/N_z$. 41-43

Step 2: Associate a triangle list $L_{i,j,k}(p)$ with each cell (i, j, k) 45-66

Step 3: Loop over all cells and assign triangles to cell lists according to algorithm 7.7

41-63

90-106

107-138



Figure 7.3: Geometric situation for the bin list algorithm. See text for explanations

It is necessary to set up completely separate lists for bins and cells since neither the size of the cells nor the criterion when a triangle is element of a certain cell's list are identical. Since the bins discussed above serve to determine whether a point is inside or outside of the domain, all triangles that intersect the bin must be in its lists. The cells discussed now however serve to determine whether a point is closer than r_c to a triangle. Therefore their size is not arbitrary but must be set to r_c . Furthermore, not all triangles that intersect a certain cell must be in its triangle list but only those being closer than r_c to the cell's centroid (thus in the in-sphere of the cell). To keep the lists as short as possible, the following algorithm is implemented in SortT (appendix C.5.4.9) to sort the triangles Δ_n to the cell lists:

Algorithm 7.7 (Sort triangles to cell lists).

Step 1: For each vertex $P_{1,2,3}$ of Δ_n , determine the index of the cell it is in: $(i, j, k) = \operatorname{ceiling}((P_x - x_l)/b_x, (P_y - y_l)/b_y, (P_z - z_l)/b_z)$, where $\operatorname{ceiling}(f)$ is the closest integer larger or equal to f. If Δ_n is not yet contained in the list $L_{i,j,k}$, add it: $L_{i,j,k}(\operatorname{length}(L_{i,j,k}) + 1) = n$.

Step 2: Determine the minimum and maximum cell indices for triangle Δ_n : $\{i, j, k\}_l = \min_{i,j,k} \{(i, j, k) : \Delta_n \in L_{i,j,k}\}, \{i, j, k\}_h = \max_{i,j,k} \{(i, j, k) : \Delta_n \in L_{i,j,k}\}$

- Step 3: Loop over all cells $\{(i, j, k) : i \in [i_l, i_h], j \in [j_l, j_h], k \in [k_l, k_h]\}$
 - 3.1: If triangle Δ_n is not already in $L_{i,j,k}$, calculate the intersection point of the triangle's normal through the cell's centroid with the triangle (see

172-197

198-203

 $^{^{4}}$ In a linked list of objects, each object can only occur once

figure 5.1 for the geometrical situation). If this intersection point is inside the triangle and inside the cell, add Δ_n to $L_{i,j,k}$.

Notice that the second algorithm (the cell lists) is only needed for boundary mirror particles. It is therefore not executed when boundary condition handling is switched off.

Using the described bin and cell list algorithms has the effect of reducing the execution time from several dozens of days to about one minute for a typical ER sample consisting of about $8 \cdot 10^4$ triangles and $2 \cdot 10^6$ grid points.

Now having completed step 5 of the top-level algorithm 7.3, some comments on steps 7 and 9 follow. These steps are concerned with the actual initialization of the particles in the regular lattice cells. The core algorithm needed is the one to determine whether a given point is inside or outside of the domain Ω . According to the solution of the topological container problem given in [Sbalzarini (2001)], this is done by counting the number of intersections of an arbitrary ray with the domain boundary $\partial\Omega$ starting at the point in question. If this count is odd, the point is inside Ω , else outside. Using the bin lists described above, this is implemented for any given point Q in $\mathcal{O}(N \log M)$. The code is given in appendix C.5.4.11 to which the marginal statement labels refer.

Algorithm 7.8 (Point in domain).

- Step 1: Determine indices of bin Q is in: $(i, j) = \operatorname{ceiling} ((Q_y y_l)/b_y, (Q_z z_l)/b_z)$ and set the intersection counter to zero: $N_i = 0$. 18-26
- Step 2: Loop over all triangles in this bin's triangle list, thus $\{\Delta_k : \Delta_k \in L_{i,j}\}$
 - 2.1: Intersect the ray $Q + \lambda[1, 0, 0]$ with triangle Δ_k by solving the following linear system of equations:

$$\begin{bmatrix} | & | & 1 \\ a & b & 0 \\ | & | & 0 \end{bmatrix} \begin{bmatrix} \alpha \\ \beta \\ -\lambda \end{bmatrix} = \begin{bmatrix} | \\ Q - P_1 \\ | \end{bmatrix}$$

where $a = P_2 - P_1$ and $b = P_3 - P_1$ are edge vectors of Δ_k and $P_{1,2,3}$ are its vertices. Since the direction can be chosen arbitrarily, it is chosen to yield a simple system of equations that can be solved at minimum cost. For this reason, the general triangle-ray intersection routine described in section 3.2 is not used but above system's solution is hard-coded.

2.2: If $\alpha \ge 0$ and $\beta > 0$ and $\alpha + \beta \le 1$ and $\lambda > 0$ (up to a tolerance TOL), increment the intersection counter by 1: $N_i = N_i + 1$. 87,105

Step 3: If N_i is odd, the point Q is inside the domain Ω , else outside.

Finally, step 9 of algorithm 7.3 adds the mirror particles for the boundary condition handling according to section 7.2.4. It is only invoked if needed and proceeds for any given point $Q \in \Omega$ as follows:

Algorithm 7.9 (Add mirror particles).

Step 1: Determine the indices of the cell
$$Q$$
 belongs to:
 $(i, j, k) = \operatorname{ceiling} ((Q_x - x_l)/b_x, (Q_y - y_l)/b_y, (Q_z - z_l)/b_z).$
119-127

Step 2: Loop over all triangles Δ_n in this cell and its 26 nearest neighbors $(i \pm 1, j \pm 1, k \pm 1)$ (if they exist). 128-131

61

73-86

109-113

2.1: Intersect the normal n_n of Δ_n through Q with Δ_n using the general intersection routine described in section 3.2: $Q + \lambda_n n_n = P_1 + \alpha a_n + \beta b_n \Rightarrow \lambda_n$.

Step 3: Find the closest triangle by selecting the one with the minimum λ : $\lambda_{min} = \min_n \lambda_n$ and $n_{min} = \operatorname{argmin}_n \lambda_n$.

Step 4: If $\lambda_{min} \leq r_c$, mirror point Q according to: $Q_{mirror} = Q + 2\lambda_{min}n_{n_{min}}$ where $n_{n_{min}}$ means the normal of the the closest triangle. Set the strength (mass) of Q_{mirror} to strength of Q and the attribute of Q_{mirror} to the negative attribute of Q.

This algorithm makes use of the cell list as described in algorithms 7.6 and 7.7 to reduce the computational cost from $\mathcal{O}(MN)$ to $\mathcal{O}(N \log M)$.

Now the preprocessing is completed and the particle locations, strengths and attributes are written to a binary file which can be read by the PSE solver as an initial condition.

7.2.6 Implementation notes

Some details on the technical implementation of the algorithms described in the preceding section seem noteworthy. They mainly concern the data structures involved and the methods of parallelization using the MPI message passing interface. Readers not interested in such details can easily skip this subsection without loss of essential information. Nevertheless, the following is included in this report for the sake of completeness since the ideas presented hereafter are neither obvious nor straightforward.

Dynamic list data structures

The algorithms presented in section 7.2.5 often involve lists of triangles associated to cells or bins. The general structure is that the bounding box of the computational domain is subdivided into a number of disjoint box-shaped sub-spaces. These subspaces are called *cells* if the subdivision takes place in all three spatial dimensions and *bins* if the space is only subdivided along two dimensions. Obviously, bins are just cells with one index removed. Therefore they can be treated using the same data structures. It was however necessary to distinguish between them in the previous section since the criteria for a triangle to be a member of a bin or cell are different. The following will only deal with cells, the application to bins (2-index cells) is straightforward.

Assume the bounding box of the domain is subdivided into $N_x \times N_y \times N_z$ cells that are addressed using integer indices $\{(i, j, k) : 1 \leq i \leq N_x, 1 \leq j \leq N_y, 1 \leq k \leq N_z\}$. To each cell, we associate a list $L_{i,j,k}$ that contains the indices p of all the triangles that "belong" to the cell (i.e. meet all the criteria of the cell): $L_{i,j,k}(p)$. The obvious way to implement such an object in Fortran 90 would be an array of rank 4. Its size in the fourth dimension would have to be equal to the number of triangles in the longest list of all cells. This is however not favorable for the following two reasons:

- 1. One triangle will typically belong to several lists (i.e. it intersects several cells), making the length of the longest list large.
- 2. The triangles are very inhomogeneously distributed in the bounding box (as they are concentrated on the surface of the domain).

135-140

142 main 391 The latter point means that in fact most cells will not contain any triangles. Just the ones intersecting the domain boundary will have lists of non-zero length. It would thus be a huge waste of memory to allocate a cubic four dimensional array.

Instead, a new data structure is defined which can be seen as an array of pointers to lists. A list is implemented as a 1D array of variable length. This allows the lengths of all the lists to be individually set and it brought down the main memory usage of the algorithm from about 500 MB to 1 MB. The following Fortran statements define the new data type as implemented in globals.f90 (see appendix C.5.4.1, statements 5 to 7 and 27, 29):

```
TYPE ptr_to_list
    REAL(MK), DIMENSION(:), POINTER :: list
END TYPE
TYPE(ptr_to_list), DIMENSION(:,:,:), ALLOCATABLE :: cell
```

The following statements allocate an array of $N_x \times N_y \times N_z$ empty triangle lists:

```
ALLOCATE(cell(Nx,Ny,Nz))
nullify(cell)
```

The current lengths of all the lists are stored in an additional array of integers: ncell(Nx,Ny,Nz). Every time we wish to add an element to a list, we first check whether the list is long enough to hold the new element. If this is not the case, it is dynamically enlarged. The process of adding an element m to the end of a list therefore is described by:

```
p = ncell(i,j,k)+1
if(ASSOCIATED(cell(i,j,k)%list)) then
    if(p > size(cell(i,j,k)%list)) then
        cell(i,j,k)%list=>reallocate(cell(i,j,k)%list,p)
    end if
else
    ALLOCATE(cell(i,j,k)%list(p))
end if
cell(i,j,k)%list(p) = m
```

The function **reallocate** encapsulates the work of dynamically resizing an object while preserving its contents. It takes a pointer to an array and returns a pointer to the new, resized array. According to [Numerical Recipes in Fortran 90 (1996)] it is multiply overloaded for different array dimensions and data types and implemented as a Fortran 90 module (see appendix C.5.4.15):

```
INTERFACE reallocate
MODULE PROCEDURE reallocate_rv,reallocate_rm,reallocate_iv, &
    reallocate_im,reallocate_hv
END INTERFACE
```

Since all triangle lists are 1D arrays of integers, only the version for integer vectors is used. It reads as follows (statements 22 to 35):

```
function reallocate_iv(p,n)
INTEGER(I4B), DIMENSION(:), POINTER
INTEGER(I4B), INTENT(IN)
INTEGER(I4B)
IN
```

```
ALLOCATE(reallocate_iv(n), STAT=ierr)
if(ierr .NE. 0) then
    WRITE(*,'(A)') 'Error allocating memory in reallocate_iv'
    return
end if
if(.NOT. ASSOCIATED(p)) return
nold = size(p)
reallocate_iv(1:min(nold,n))=p(1:min(nold,n))
DEALLOCATE(p)
END function reallocate_iv
```

To simplify the usage of these list types, the subroutine AllocateLL has been implemented (see appendix C.5.4.10). It takes 5 integer arguments: a flag, the three indices of the cell the operation should affect and an error status. Table 7.1 summarizes possible ways of using it. After the subroutine returns, the value of the fifth variable is zero if no error occurred, else it is non-zero.

Call	Purpose
AllocateLL(0,0,0,0,istat)	Initialize all lists and allocate memory
AllocateLL(1,0,i,j,istat)	Enlarge bin list (i, j)
AllocateLL(1,i,j,k,istat)	Enlarge cell list (i, j, k)
AllocateLL(2,0,0,0,istat)	Clear all lists and free all memory

Table 7.1: Calling modi of AllocateLL

Paradigms of parallelism

For reasons of speed and to be prepared for larger problems, the particle preprocessing algorithm 7.3 is parallelized using MPI. Examination of algorithm 7.3 reveals that the most time consuming part is step 7 as it consists of invoking algorithm 7.8 for every grid point. Two different ways of parallelizing this step can be imagined:

- 1. Distribute grid points: all processors know all triangles but only check a fraction of the grid points.
- 2. Distribute triangles: one processor loops over all the grid points but the triangles to be intersected with are distributed among the other processors.

For the following reasons, the first possibility is better than the second:

- The number of points is typically more than an order of magnitude larger than the number of triangles.
- Communication only happens once at the very beginning of step 7 and once at its end (for possibility 2 one would have to collect the numbers of intersection after every point).

The amount of data that has to be communicated is 9M + G for the first case and $G(\log M + n)$ for the second with G being the number of grid points, M the number of triangles and n the number of processors. Even though the amount of data that has to be transmitted is larger in the first case, only two communication channels have to be opened as opposed to $G \approx 10^6$ is the second case. Since for MPI, the latency of opening a channel is crucial rather than the bulk amount of data transmitted once the channel is open, the first alternative will be faster.

The parallel algorithm therefore proceeds as follows (statement numbers refer to the listing in appendix C.5.4.2):

```
58
```

Algorithm 7.10 (MPI parallelization).

Step 1:	Start MPI and probe all channels and processors.	37-105
Step 2:	Master processor reads all triangles from file.	110
Step 3:	Master broadcasts all triangles to slaves.	127-147
Step 4:	Master does cartesian domain decomposition and assigns a set of grid points to check to each processor.	176-191
Step 5:	Each processor loops over all grid points it got assigned and checks whether they are inside or outside the domain (i.e. invokes algorithm 7.8). The results are stored in a local flag array.	294-316
Step 6:	Slaves communicate their results (flag) to master.	319-330
Step 7:	Master concatenates flag arrays.	327

The array **flag** temporarily holds the local result of algorithm 7.8 for all grid points on a processor. This is recommended to avoid the need of communicating them back to the master after each grid point.

7.3 Post-processing and visualization

The PSE simulation code following the algorithms described in sections 7.2.2 and following can write output files at certain time steps containing the positions and strengths of all the particles. The files are in ASCII format and each line contains the information of one particle as:

(4E16.8) xp(1,i), xp(2,i), xp(3,i), c(i)

where xp(1..3,i) are the x, y and z coordinates of the position of particle i and c(i) is its strength (i.e. the local concentration).

In order to be able to use these output files to visualize concentration fields in OpenDX⁵, they have to be post-processed. Just as other visualization tools, OpenDX can only handle data when its *connection structure* is known. This means that the data points either have to be on a grid or a triangulation of the scattered data locations has to be performed, which is $\mathcal{O}(N^4)$ however. Since the preprocessor described in section 7.2.5 initializes the particles on a regular lattice and they do not move during the PSE simulation (there is no convection), it is best to convert the output files to a grid representation again. The grid will extend throughout the bounding box of the particles. However, not all grid points will be occupied by a particle since they have only been placed inside the ER geometry. The conversion therefore consists of determining which grid points are associated to which particle and which are empty.

The program **res2dx** as given in appendix C.3.2.1 performs this task for all the input files passed to it as command line arguments. Before being able to translate the PSE result files to OpenDX grid data, it has to recover the grid geometry from the available data. Since the grid is known to be cartesian and regular, this only includes the numbers of grid points N_x , N_y , N_z in all three spatial directions, the bounding box of the data and the corresponding grid spacings δx , δy , δz . Using this information, the OpenDX header file (called a *general* file) can then be written. This file contains all the information about the data's structure and organization.

⁵OpenDX is a freely available data visualization tool formerly known as IBM Visualization Data Explorer. See http://www.opendx.org for further information.

For regular grid data, its syntax is:

```
file = <name of grid data file>
grid = <Nx> x <Ny> x <Nz>
format = ascii
interleaving = record
majority = row
field = c
structure = scalar
type = float
dependency = positions
positions = regular, regular, regular, <xl>, <dx>, <yl>, <dy>, <zl>, <dz>
end
```

where xl, yl, zl are the lower boundaries for x, y and z (i.e. the coordinates of the lower left corner of the bounding box). Then, the strengths of all the particles are written to a separate grid data file that simply contains one floating point number per line:

c(i)

indicating the strength of the particle at grid point *i*. Hereby, the grid points are listed in row-majority order, as declared in the header file. This means that for each grid point, the conversion program has to find the particle that belongs to it in order to write its strength to the output file. Since particles are only placed where there is ER geometry, not every grid point will have a particle, so this check is actually needed. Without further tricks, this would be $\mathcal{O}(NN_xN_yN_z)$ and thus unfeasible. A chaining mesh of cells is therefore established in the bounding box of the data which is thus subdivided into $Nc_x \times Nc_y \times Nc_z$ mesh cubes *C* of size $b_x \times b_y \times b_z$:

$$C_{i,j,k} = \{ (x, y, z) : (i-1)b_x \leq x - x_l < ib_x, \ (j-1)b_y \leq y - y_l < jb_y, \\ (k-1)b_z \leq z - z_l < kb_z \}$$

Unlike for all the triangle lists discussed so far, a particle belongs to exactly one such cell $C_{i,j,k}$. It is therefore possible to use a linked list of particles instead of the dynamic list data structures presented in section 7.2.6. To do so, a *head of chain* particle $H_{i,j,k}$ is assigned to each cell $C_{i,j,k}$ and a linked list of particles $\ell(p)$ is set up for $p = 1, \ldots, N$ such that each entry points to the next particle in the same cell. The end of the list is marked by a zero entry. This cell-list algorithm brings the computational cost for the conversion down from $\mathcal{O}(NN_xN_yN_z)$ to $\mathcal{O}(N_xN_yN_z)$. Figure 7.5 illustrates this structure.

The following algorithm describes the conversion process in detail. The marginal numbers refer to the statement labels in the source code as given in appendix C.3.2.1.

Algorithm 7.11 (Conversion of particles to grid data).

- Step 1: Determine number of particles in input file: N
 - Step 2: Determine bounding box of data $[x_l, x_h] \times [y_l, y_h] \times [z_l, z_h]$ as: $\{x, y, z\}_l = \min_p \{x_p, y_p, z_p\}, \{x, y, z\}_l = \min_p \{x_p, y_p, z_p\}$ where p loops over all input particles: $p = 1, \ldots, N$
 - Step 3: Determine grid spacings as the minimum distances of all possible pairwise distinct particles in all three spatial directions: $\delta x = \min_{(p,q), p \neq q} |x_p x_q|, \delta y = \min_{(p,q), p \neq q} |y_p y_q|, \delta z = \min_{(p,q), p \neq q} |z_p z_q|$

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61-63

55-60

53



Figure 7.4: OpenDX visual program

Figure 7.5: A linked list data structure

Step 4:	Calculate the number of grid points as: $(N_x, N_y, N_z) = \text{floor}((x_h - x_l)/\delta x, (y_h - y_l)/\delta y, (z_h - z_l)/\delta z)$	67-69
Step 5:	Write data header file according to above syntax	74-87
Step 6:	Nullify the linked list: $H_{i,j,k} = 0 \forall (1 \leq i \leq Nc_x, 1 \leq j \leq Nc_y, 1 \leq k \leq Nc_z), \ell(p) = 0 \forall p \in [1, N]$	88-99
Step 7:	Determine the size of the cells: $b_x = (x_h - x_l)/Nc_x$, $b_y = (y_h - y_l)/Nc_y$, $b_z = (z_h - z_l)/Nc_z$	107-109
Step 8:	For each particle $x_p, p = 1, \ldots, N$:	
	 8.1: Determine the indices of the cell the particle is in: (i, j, k) = ceiling((x_p - x_l)/b_x, (y_p - y_l)/b_y, (z_p - z_l)/b_z) 8.2: Add particle p to the linked list of cell C_{i j k}: ℓ(p) = H_{i j k}, H_{i j k} = p 	111-119 120-121
Step 9:	For each grid point (x, y, z) :	
	9.1: Determine the indices of the cell it is in: $(i, j, k) = \text{ceiling}((x-x_l)/b_x, (y-y_l)/b_y, (z-z_l)/b_z)$	133-141
	9.2: Loop over all particles in the same cell by following the linked list. If a particle is found that is closer to the grid point (x, y, z) than $r = \frac{1}{2} \min\{\delta x, \delta y, \delta z\}$, write its strength to the output grid data file. If no particle is found at all, output 0.	142-152

Step 9 loops over all grid points in row majority. This means that the innermost loop is over z and the outermost is over x. Traversing the linked list in step 9.2 for a cell $C_{i,j,k}$ is done as follows:

```
\begin{array}{l} f=0\\ p=H_{i,j,k}\\ \text{do while }p\neq 0 \text{ and }f=0\\ \text{ if }\left((x_p,y_p,z_p)-(x,y,z)\right)^2 < r^2 \text{ then }\\ \text{ write }c(p)\\ f=1\\ \text{ end if }\\ p=\ell(p)\\ \text{ end do}\\ \text{ if }f=0 \text{ then }\\ \text{ write }0\\ \text{ end if } \end{array}
```

Using the header and grid files produced by this algorithm, the visual program shown in figure 7.4 can be used in OpenDX to produce 3D visualizations of the concentration field. Figure 7.6 shows an example taken from a simulation run of section 9.1.



Figure 7.6: Example visualization of a simulated concentration field

Chapter 8

Tests and validation

8.1 Test case description

In order to be able to validate the random walk and PSE simulation codes, a simple test case for which an analytic solution exists is sought. To maintain the link to FRAP analysis, the cubic box used throughout [Sbalzarini (2001)] is chosen, thus the problem domain consists of the cube:

$$\Omega = [0, L] \times [0, L] \times [0, L]$$

with boundary $\partial\Omega$ and interior $int(\Omega) = \{\Omega \setminus \partial\Omega\}$. Since ER lumen proteins do not cross the ER membrane, the biologically correct boundary condition is zero flux at all walls, thus:

$$\frac{\partial(\cdot)}{\partial n} = \nabla(\cdot) \cdot n = 0 \qquad \text{on } \partial\Omega$$

where n is the outer normal on $\partial \Omega$.

This cubic domain is intersected by the square cylinder:

$$B = [a, b] \times [c, d] \times [0, L]$$

which can be thought of as the area where photobleaching takes place at time $t = 0^{-}$ (i.e. just before time zero). Therefore, the cylinder B will be called "bleached box" hereafter. See figure 8.1 for an overview of this geometrical setting. The initial condition that corresponds to a FRAP setting is:

$$u_0(x) = \begin{cases} 0 & \text{if } x \in B\\ c_0 & \text{if } x \in int(\Omega) \setminus B \end{cases}$$
(8.1)

The governing equation to be solved on this geometry with above initial and boundary conditions is the diffusion equation:

$$\begin{cases} \frac{\partial u}{\partial t} = D\nabla^2 u + q(x,t) & \text{for } x \in int(\Omega), \ t > 0\\ \alpha \frac{\partial u}{\partial n} + \beta u = r(x,t) & \text{for } x \text{ on } \partial\Omega\\ u(x,t=0) = u_0(x) & \text{for } x \in int(\Omega) \end{cases}$$
(8.2)

where ∇^2 is the Laplacian operator and u(x,t) is the unknown function to be solved for. Since we have zero flux boundary conditions and there are no sources to be considered, the simplifications $q \equiv 0$, $r \equiv 0$, $\alpha = 1$ and $\beta = 0$ are readily made.

8.2 The analytic solution

Eigenfunctions and Eigenvalues

The analytic solution of the test case as defined in the previous section can be derived using the general approach of eigenfunction series expansions. The basic domain Ω is the cartesian product of three identical intervals. Therefore, the normalized eigenfunctions $\varphi(x)$ and eigenvalues λ of Helmholtz's equation on the 1D domain [0, L] are needed. The eigenproblem to be solved is:

$$\begin{cases} \nabla^2 \varphi + \lambda \varphi = 0 & \text{if } x \in (0, L) \\ \frac{d\varphi}{dx} = 0 & \text{for } x = 0 \text{ and } x = L \end{cases}$$
(8.3)

Writing φ' for $\frac{d\varphi}{dx}$, this becomes:

$$\varphi''(x) + \lambda\varphi(x) = 0$$

which has the general solution

$$\varphi(x) = A\cos\left(\sqrt{\lambda}x\right) + B\sin\left(\sqrt{\lambda}x\right)$$

using the boundary condition at x = 0 yields:

$$\varphi'(0) = B\sqrt{\lambda} \stackrel{!}{=} 0 \qquad \Rightarrow \quad B = 0$$

since λ must not vanish in all cases. The second boundary condition at x = L leads to:

$$\varphi'(L) = -A\sqrt{\lambda}\sin\left(\sqrt{\lambda}L\right) \stackrel{!}{=} 0 \qquad \Rightarrow \quad \sin\left(\sqrt{\lambda}L\right) = 0$$
$$\Rightarrow \quad \sqrt{\lambda}L = k\pi \qquad \forall k \in \mathbb{N}_0^+$$

since $\lambda \equiv 0$ is, again, not a valid solution. These eigenfunctions are now normalized to unity by requiring:

$$\int_0^L \varphi^2(x) \, dx \stackrel{!}{=} 1$$

which is equal to setting their scalar products to one¹. For k > 0, this corresponds to:

$$A^{2} \int_{0}^{L} \cos^{2} \left(\sqrt{\lambda}x\right) dx \stackrel{!}{=} 1$$

$$\Rightarrow A^{2} \left[\frac{1}{2\sqrt{\lambda}} \left(x + \sin\left(\sqrt{\lambda}x\right)\cos\left(\sqrt{\lambda}x\right)\right)\right]_{0}^{L} = \frac{1}{2}A^{2}L \stackrel{!}{=} 1$$

$$\Rightarrow A = \sqrt{\frac{2}{L}}$$

For $k = 0, \lambda$ is zero as well and above condition becomes:

$$A^2 \int_0^L dx = A^2 L \stackrel{!}{=} 1 \qquad \Rightarrow \quad A = \sqrt{\frac{1}{L}}$$

 $^{^1\}mathrm{We}$ require the eigenfunctions to be an orthonormal basis of the solution space of problem 8.2

Therefore, the required normalized eigenfunctions and eigenvalues for this case are given by:

$$2\varphi_k(x) = \begin{cases} \sqrt{\frac{2}{L}}\cos\left(\frac{k\pi}{L}x\right) & \text{if } k > 0\\ \sqrt{\frac{1}{L}} & \text{if } k = 0 \end{cases}$$
(8.4)

$$\lambda_k = \left(\frac{k\pi}{L}\right)^2 \qquad \qquad k \in \mathbb{N}_0^+ \tag{8.5}$$

It is easy to verify that the $\varphi(x)$ fulfill

$$\int_0^L \varphi_i(x)\varphi_j(x)\,dx = \delta_{ij}$$

which is necessary and sufficient for an orthonormal function base.





Figure 8.1: Geometry of test case

Figure 8.2: Slice through domain with integral areas

Solution by Eigenfunction series expansion

The general solution of problem 8.2 for $\alpha = 1$ and $\beta = 0$ is given by:

$$\begin{aligned} u(x,t) &= \int_{\Omega} K(x,\xi,t) u_0(\xi) \, d\xi + \int_0^t \int_{\Omega} K(x,\xi,t-\tau) q(\xi,\tau) \, d\xi \, d\tau \\ &+ D \int_0^t \int_{\partial\Omega} K(x,\xi,t-\tau) r(\xi,\tau) \, ds_\xi \, d\tau \end{aligned}$$

where ds_{ξ} means an infinitesimal segment of $\partial\Omega$ and $K(x, \xi, t)$ is the heat kernel (see also [Sbalzarini (2001)] pp. 20-21). According to the product space theorem, the heat kernel of a product space is equal to the product of the heat kernels of its cartesian factors. Since Ω is the cartesian product of three intervals in x, y and z, one writes:

$$K(x, y, z, \xi, \eta, \zeta, t) = K_x(x, \xi, t) \cdot K_y(y, \eta, t) \cdot K_z(z, \zeta, t)$$

$$(8.6)$$

where the individual factor kernels $K_{x,y,z}$ are given by:

$$K(x,\xi,t) = \sum_{k=0}^{\infty} \varphi_k(x)\varphi_k(\xi)e^{-D\lambda_k t}$$
(8.7)

$$K(y,\eta,t) = \sum_{k=0}^{\infty} \varphi_k(y) \varphi_k(\eta) e^{-D\lambda_k t}$$
(8.8)

$$K(z,\zeta,t) = \sum_{k=0}^{\infty} \varphi_k(z)\varphi_k(\zeta)e^{-D\lambda_k t}$$
(8.9)

with φ_k and λ_k being above eigenfunctions and eigenvalues of Helmholtz's equation 8.3. Using $q \equiv 0$ and $r \equiv 0$, the solution can formally be written as:

$$u(x, y, z, t) = \int_{\Omega} K(x, y, z, \xi, \eta, \zeta, t) u_0(\xi, \eta, \zeta) \, d\xi \, d\eta \, d\zeta$$

Taking into account the special topology of the initial condition, this is equal to:

$$u(x, y, z, t) = c_0 \int_{\Omega \setminus B} K(x, y, z, \xi, \eta, \zeta, t) \, d\xi \, d\eta \, d\zeta$$

and basically consists of integrating the heat kernel over $\Omega \setminus B$. This is done by splitting the integral into four parts according to figure 8.2:

$$u(x, y, z, t) = c_0 \underbrace{\int_0^L \int_0^c \int_0^L K \, d\xi \, d\eta \, d\zeta}_{I} + c_0 \underbrace{\int_0^a \int_c^d \int_0^L K \, d\xi \, d\eta \, d\zeta}_{III} + c_0 \underbrace{\int_b^L \int_c^d \int_0^L K \, d\xi \, d\eta \, d\zeta}_{III} + c_0 \underbrace{\int_b^L \int_c^d \int_0^L K \, d\xi \, d\eta \, d\zeta}_{III} + c_0 \underbrace{\int_0^L \int_c^L \int_0^L K \, d\xi \, d\eta \, d\zeta}_{III} + c_0 \underbrace{\int_0^L \int_c^L \int_0^L K \, d\xi \, d\eta \, d\zeta}_{III} + c_0 \underbrace{\int_0^L \int_c^L \int_0^L K \, d\xi \, d\eta \, d\zeta}_{III} + c_0 \underbrace{\int_0^L \int_c^L \int_0^L K \, d\xi \, d\eta \, d\zeta}_{III} + c_0 \underbrace{\int_0^L \int_c^L \int_0^L K \, d\xi \, d\eta \, d\zeta}_{III} + c_0 \underbrace{\int_0^L \int_c^L \int_0^L K \, d\xi \, d\eta \, d\zeta}_{III} + c_0 \underbrace{\int_0^L \int_c^L \int_0^L K \, d\xi \, d\eta \, d\zeta}_{III} + c_0 \underbrace{\int_0^L \int_c^L \int_0^L K \, d\xi \, d\eta \, d\zeta}_{III} + c_0 \underbrace{\int_0^L \int_c^L \int_0^L K \, d\xi \, d\eta \, d\zeta}_{III} + c_0 \underbrace{\int_0^L \int_c^L \int_0^L K \, d\xi \, d\eta \, d\zeta}_{III} + c_0 \underbrace{\int_0^L \int_c^L \int_0^L K \, d\xi \, d\eta \, d\zeta}_{III} + c_0 \underbrace{\int_0^L \int_c^L \int_0^L K \, d\xi \, d\eta \, d\zeta}_{III} + c_0 \underbrace{\int_0^L \int_c^L \int_0^L K \, d\xi \, d\eta \, d\zeta}_{III} + c_0 \underbrace{\int_0^L \int_c^L \int_0^L K \, d\xi \, d\eta \, d\zeta}_{III} + c_0 \underbrace{\int_0^L \int_0^L \int_0^L \int_0^L \int_0^L \int_0^L K \, d\xi \, d\eta \, d\zeta}_{III} + c_0 \underbrace{\int_0^L \int_0^L \int_0^L \int_0^L \int_0^L \int_0^L \int_0^L K \, d\xi \, d\eta \, d\zeta}_{III} + c_0 \underbrace{\int_0^L \int_0^L \int_0^L \int_0^L \int_0^L \int_0^L K \, d\xi \, d\eta \, d\zeta}_{III} + c_0 \underbrace{\int_0^L \int_0^L \int_0^L \int_0^L \int_0^L \int_0^L K \, d\xi \, d\eta \, d\zeta}_{III} + c_0 \underbrace{\int_0^L \int_0^L \int_0^L \int_0^L \int_0^L \int_0^L \int_0^L K \, d\xi \, d\eta \, d\zeta}_{III} + c_0 \underbrace{\int_0^L \int_0^L \int_0^L \int_0^L \int_0^L \int_0^L K \, d\xi \, d\eta \, d\zeta}_{III} + c_0 \underbrace{\int_0^L \int_0^L \int_0^L \int_0^L \int_0^L \int_0^L \int_0^L K \, d\xi \, d\eta \, d\zeta}_{III} + c_0 \underbrace{\int_0^L \int_0^L \int_0^L \int_0^L \int_0^L \int_0^L \int_0^L K \, d\xi \, d\eta \, d\zeta}_{III} + c_0 \underbrace{\int_0^L \int_0^L \int_0^$$

The only differences between these integrals are their limits. It is therefore sufficient to know the general definite integral of the k-th Eigenfunction:

$$\int_{a}^{b} \varphi_{k}(x) dx = \begin{cases} \sqrt{\frac{2}{L}} \frac{L}{k\pi} \left[\sin\left(\frac{k\pi}{L}b\right) - \sin\left(\frac{k\pi}{L}a\right) \right] & \text{if } k > 0\\ \sqrt{\frac{1}{L}} \left(b - a\right) & \text{if } k = 0 \end{cases}$$
(8.11)

The four integral terms are now treated in turn.

Integral I

Substituting the heat kernel from equations 8.6 to 8.9 into integral I gives:

$$\int_0^L \int_0^c \int_0^L \sum_{k=0}^\infty \sum_{l=0}^\infty \sum_{m=0}^\infty \varphi_k(x)\varphi_l(y)\varphi_m(z)\varphi_k(\xi)\varphi_l(\eta)\varphi_m(\zeta)e^{-D(\lambda_k+\lambda_l+\lambda_m)t}\,d\xi\,d\eta\,d\zeta$$

Switching the order of integration and summation and rearranging some terms this becomes:

$$\sum_{k=0}^{\infty}\sum_{l=0}^{\infty}\sum_{m=0}^{\infty}\varphi_k(x)\varphi_l(y)\varphi_m(z)\int_0^L\varphi_k(\xi)\,d\xi\int_0^c\varphi_l(\eta)\,d\eta\int_0^L\varphi_m(\zeta)\,d\zeta\,e^{-D(\lambda_k+\lambda_l+\lambda_m)t}$$

This can now be integrated using 8.11 three times with appropriate substitutions. Then, since $\sin(k\pi) \equiv 0 \ \forall k$, only terms with k = m = 0 contribute to the sums. Taking this into account, substituting explicit expressions for the eigenfunctions using equation 8.4 and taking all terms for l = 0 out of the sum to get rid of the case distinctions leads to the final result for integral I:

$$I = \frac{c}{L} + \sum_{l=1}^{\infty} \frac{2}{l\pi} \sin\left(\frac{l\pi}{L}c\right) \cos\left(\frac{l\pi}{L}y\right) e^{-D\lambda_l t}$$
(8.12)

Integral II

Substituting the heat kernel from equations 8.6 to 8.9 into integral II, switching the order of integration and summation and doing some rearrangements of terms gives:

$$\sum_{k=0}^{\infty}\sum_{l=0}^{\infty}\sum_{m=0}^{\infty}\varphi_k(x)\varphi_l(y)\varphi_m(z)\int_0^a\varphi_k(\xi)\,d\xi\int_c^d\varphi_l(\eta)\,d\eta\int_0^L\varphi_m(\zeta)\,d\zeta\,e^{-D(\lambda_k+\lambda_l+\lambda_m)t}$$

Using 8.11 for each integral in this expression and noticing that only terms with m = 0 are non-zero leads to the following final expression for integral II:

$$II = \sum_{k=0}^{\infty} \sum_{l=0}^{\infty} \varphi_k(x) \varphi_l(y) \underbrace{\left\{\begin{array}{c} \frac{\sqrt{2L}}{k\pi} \sin\left(\frac{k\pi}{L}a\right)}{\sqrt{\frac{1}{L}a}}\right\}}_{t_1} \underbrace{\left\{\begin{array}{c} \frac{\sqrt{2L}}{l\pi} \sin\left(\frac{l\pi}{L}d\right)}{\sqrt{\frac{1}{L}} (d-c)}\right\}}_{t_3} e^{-D(\lambda_k + \lambda_l)t}$$
(8.13)

where the stacked notation in curly braces means that the upper expression is to be used if the corresponding index variable (k or l) is > 0, the lower one if it is zero.

Integral III

Substituting the heat kernel from equations 8.6 to 8.9 into integral III, again switching the order of integration and summation and doing some rearrangements, one finds:

$$\sum_{k=0}^{\infty}\sum_{l=0}^{\infty}\sum_{m=0}^{\infty}\varphi_k(x)\varphi_l(y)\varphi_m(z)\int_b^L\varphi_k(\xi)\,d\xi\int_c^d\varphi_l(\eta)\,d\eta\int_0^L\varphi_m(\zeta)\,d\zeta\,e^{-D(\lambda_k+\lambda_l+\lambda_m)t}$$

Using 8.11 for each integral in this expression and noticing that only terms with m = 0 contribute to the sum leads to the following final expression for integral III:

$$III = \sum_{k=0}^{\infty} \sum_{l=0}^{\infty} \varphi_k(x) \varphi_l(y) \underbrace{\left\{\begin{array}{c} -\frac{\sqrt{2L}}{k\pi} \sin\left(\frac{k\pi}{L}b\right) \\ \sqrt{\frac{1}{L}}\left(L-b\right) \end{array}\right\}}_{t_2} \underbrace{\left\{\begin{array}{c} \frac{\sqrt{2L}}{l\pi} \sin\left(\frac{l\pi}{L}d\right) \\ \sqrt{\frac{1}{L}}\left(d-c\right) \end{array}\right\}}_{t_3} e^{-D(\lambda_k + \lambda_l)t}$$

$$(8.14)$$

Integral IV

Substituting the heat kernel from equations 8.6 to 8.9 into integral IV and doing the usual rearrangements gives:

$$\sum_{k=0}^{\infty}\sum_{l=0}^{\infty}\sum_{m=0}^{\infty}\varphi_k(x)\varphi_l(y)\varphi_m(z)\int_0^L\varphi_k(\xi)\,d\xi\int_d^L\varphi_l(\eta)\,d\eta\int_0^L\varphi_m(\zeta)\,d\zeta\,e^{-D(\lambda_k+\lambda_l+\lambda_m)t}$$

Inserting explicit expressions for the eigenfunctions from 8.4 and again using equation 8.11 in its appropriate form and the fact that only for k = m = 0 all of the three factors are non-zero, the final result for integral IV becomes:

$$IV = \frac{L-d}{L} - \sum_{l=1}^{\infty} \frac{2}{l\pi} \sin\left(\frac{l\pi}{L}d\right) \cos\left(\frac{l\pi}{L}y\right) e^{-D\lambda_l t}$$
(8.15)

Notice that the sum is to be taken starting from l = 1 rather than l = 0 since the l = 0 terms have been explicitly taken out of the sum to get rid of case distinctions (cf. also integral I).

Final solution and discussion

Now the final solution for u(x,t) is known from equation 8.10 by substituting all integral expressions 8.12, 8.13, 8.14 and 8.15:

$$u(x, y, t) = c_0 \left(I(y) + II(x, y) + III(x, y) + IV(y) \right)$$
(8.16)

If is noteworthy that the exact solution to the test problem turns out to be independent of z. This is due to the zero flux Neumann boundary condition as it allows free slip in directions tangential to the domain boundary $\partial\Omega$. Dirichlet (no slip) boundary conditions would in fact induce a z-dependence since the solution would have to skew towards the two ends of the bleached box in order to maintain u = 0 on $\partial\Omega$. Using Neumann boundary conditions on the other hand, the concentration front diffuses in orthogonal to $\partial\Omega$ and the solution will look the same on all z-planes. This is no contradiction to the findings in [Sbalzarini (2001)] as the apparent D values will be different in 2D and 3D since all moves parallel to the z axis are missed by the observer. Nevertheless, the shape of the solution in 3D looks the same on all z-planes. It will however not look the same as the 2D solution at the same time. At different times though, the 2D and 3D solutions will in fact be equal and it is exactly due to this time shift that the apparent diffusion constant is different²

Numerical results and visualizations

The exact solution as given by equation 8.16 consists of an infinite series expansion in a function space whose base functions are the eigenfunctions of Helmholtz's equation on the given domain Ω . In order to get useful results, a numerical evaluation of it for certain locations in space and certain times is made. First implementations in MATLAB turned out to be impracticable as the solution consists of double sums which makes it $\mathcal{O}(M^2)$ to be evaluated using M terms in the expansion. Therefore, a Fortran program has been written to do the evaluation fast. The corresponding source code is given in appendix C.1.2.1. The code is written to minimize its execution time by pre-calculating as many terms as possible and using the fact that the same terms appear in different places of the solution, e.g. t_1, t_2 and t_3 as labeled in equations 8.13 and 8.14. Moreover, it contains the compiler directives needed to parallelize it on a vector processor.

 $^{^{2}}$ recall that the diffusion constant basically reflects the solution's time scales for fixed lengths.

The solution is approximated on a regular cartesian grid with spacing δ in each direction up to a certain number of terms in the series expansion and for discrete time steps, thus:

$$u(x, y, t) \approx \widetilde{u}(x_i, y_j, t_n)$$

where

$$x_i = (i-1)\delta \qquad \qquad y_i = (j-1)\delta$$

for i, j = 1, 2, 3, ... It turns out that taking 300 terms of the expansion is sufficient to obtain a converged solution to machine precision. The program has been tested to recover the initial condition for $t \downarrow 0$. To be mathematically correct and to avoid Gibbs oscillations (sharp corners in the initial condition !), the final code however directly outputs the initial condition for t = 0 and only employs the series expansion for t > 0. This is the mathematically correct way since to solution of a PDE is only defined for t > 0 anyway.

To be able to check the boundary condition, the conservation of mass in the domain Ω is considered. The total mass is hereby calculated as the sum of the masses of all grid cells. A grid cell is defined as the parallelepiped between 4 grid points, thus:

$$C_{i,j} = \{(i,j) : (i-1)\delta \leq x < i\delta, \ (j-1)\delta \leq y < j\delta\}$$

and its mass $m(C_{i,j})$ is given by the concentration at its center times its volume where the center concentration c_c is taken to be the mean of the concentrations at the surrounding grid vertices:

$$m(C_{i,j}, t_n) = c_c(t_n)\delta^2$$
$$c_c(t_n) = \frac{1}{4} \left(\tilde{u}(x_i, y_j, t_n) + \tilde{u}(x_i + 1, y_j, t_n) + \tilde{u}(x_i, y_j + 1, t_n) + \tilde{u}(x_i + 1, y_j + 1, t_n) \right)$$

This is the consistent way of evaluating the mass. Just taking the sum of \tilde{u} over all grid points would not be sufficient (in fact, the results would be wrong).

In order to be able to get a reference FRAP curve, a corresponding output facility is implemented in the code. The FRAP value at time t_n is taken to be the mean concentration in the bleached box B, thus:

$$F(t_n) = \frac{1}{\mathcal{L}^3(B)} \sum_{(i,j) \in B} m(C_{i,j}, t_n)$$

where $\mathcal{L}^{3}(B)$ means the volume of B. This is consistent with above way of evaluating the mass and physically meaningful since the light intensity of the fluorescence is proportional to the concentration of the fluorophore in the area in question. Plotting F vs. t now approximates the exact FRAP curve of the test case to machine precision.

For the concrete test case the numerical parameter values from [Sbalzarini (2001)] are taken. Hence Ω is the cube of edge length 4 and the bleached box is $B = [2.0, 3.0] \times [2.0, 3.0] \times [0.0, 4.0]$. The initial concentration in $\Omega \setminus B$ is chosen to be unity, the diffusion coefficient is 0.03. For the numerical evaluation, a cartesian 41×41 grid is chosen and the solution is calculated on all grid points every 0.05 s. Table 8.1 summarizes these problem parameters.

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Parameter	Value
a	2.0
b	3.0
с	2.0
d	3.0
L	4.0
c_0	1.0
D	0.03
Grid	41×41
Time step	0.05
Number of terms	300

Table 8.1: Test case parameters

Figure 8.3 shows snapshots of the exact solution at times t = 0, t = 0.3, t = 1, t = 5, t = 10 and t = 20. Figure 8.4 shows the reference FRAP curve $F(t_n)$. The dashed blue line indicates the steady state value for $t \to \infty$. Due to the no flux boundary condition this is lower than 1 as Ω is initially not completely filled. Figure 8.5 shows the total mass inside Ω . As expected from the boundary condition, it is conserved with an absolute error (due to the finite number of terms that have been taken for the series expansion) of 10^{-14} which is comparable to the machine epsilon. Moreover, a movie animation of the exact solution is available on volume 1 of the companion CD of this report (see appendix A).



Figure 8.3: Snapshots of the exact solution at different times



Figure 8.4: Analytic FRAP curve (red) with asymptote 0.9375 (blue)



Figure 8.5: Total mass for analytic solution

8.3 A finite difference code

To double-check the exact solution derived in section 8.2 and to have an additional benchmark for the PSE simulation code, a simple high resolution finite difference code is implemented for the test case presented in section 8.1.

To keep the possibility of programming errors as low as possible, the standard CRAYFISHPAK library³ version 2.0 is used. This library contains fast solvers for the non-homogeneous linear Helmholtz equation:

$$\nabla^2 u + \lambda u = f$$

Doing an implicit Euler time discretization of the diffusion equation leads to Helmholtz's equation for u_{n+1} as follows:

$$\begin{aligned} \frac{\partial u}{\partial t} &= D\nabla^2 u\\ \text{Euler:} & u_{n+1} &= u_n + D\delta t \nabla^2 u_{n+1} \\ \Rightarrow \nabla^2 u_{n+1} - \frac{1}{D\delta t} u_{n+1} &= -\frac{1}{D\delta t} u_n \end{aligned}$$

Solving Helmholtz's equation with $\lambda = -(D\delta t)^{-1}$ and $f = \lambda u_n$ for the new solution u_{n+1} in every time step thus corresponds to solving the diffusion equation with implicit time stepping.

Neumann zero flux boundary conditions are taken on all 6 walls of the cubic computational domain Ω , which is reflected by boundary condition type 3 and boundary value 0 in CRAYFISHPAK. To set the initial condition given in equation 8.1, an initialization routine (H3GCIS in this case) is called once. Inside the time loop, the fast solver H3GCSS for regular cartesian 3D grids is called repeatedly. The complete Fortran source is given in appendix C.2.2.1.

The test case is run with the geometry parameters given in table 8.1 on a regular cartesian $81 \times 81 \times 81$ mesh with time step $\delta t = 0.005$ and D = 0.03. Figure 8.6 shows the resulting FRAP curve in comparison to the one of the exact solution. It can be seen that the finite difference code stays below the exact solution due to numerical dissipation. The total mass in the domain is again calculated as the sum of the masses of all grid cells where the mass of a grid cell is defined to be the concentration at its center times its volume. The concentration at the center is

³Commercial library by Green Mountain software, Madeira Beach, FL.

taken to be the mean value of the concentrations of the 8 surrounding grid vertices (cf. section 8.2). As a check, figure 8.7 shows the total mass vs. time. It is conserved up to an absolute error of 10^{-9} which reflects the no flux boundary condition very well.





Figure 8.6: FRAP curve of FD code (red) compared to exact solution (blue)

Figure 8.7: Total mass for the FD code

The discretisation used is first order in time and second order in space and fails to recover the error introduced by the very steep gradients of the initial condition. This can be seen when running it with starting time $t_0 = 0.1$ rather than from the beginning. The initial condition given by equation 8.1 is now replaced with the exact solution at time t_0 . Since this is smooth, the finite difference code now performs very well indicating that the error observed in figure 8.6 indeed stems from the initial condition. Figure 8.8 shows the corresponding FRAP curve in comparison to the exact solution and the FRAP curve obtained when starting at $t_0 = 0$. Figure 8.9 shows the total mass vs. time, which is conserved within an absolute error of 10^{-9} again.





Figure 8.8: FRAP curve of FD starting at $t_0 = 0.1$ (red) compared to exact solution (green) and FD starting at $t_0 = 0$ (blue)

Figure 8.9: Total mass for FD code starting at $t_0 = 0.1$

As the time discretization is implicit, it is possible to use very large time steps, exceeding the CFL limit⁴ of the problem and still getting a bounded solution. To check this, the previous run is repeated with $\delta t = 1.0$ which is about two orders of magnitude larger than the CFL limit $\delta t < \frac{\delta x^2}{2D} = 0.041667...$ It can be seen from figure 8.10 that the solution indeed is bounded and does not deviate grossly from the previous one. The mass is still conserved up to an absolute error of 10^{-9} .

⁴stating that $2D\delta t/\delta x^2 \leq 1$ for $x \in \mathbb{R}^n$



Figure 8.10: FRAP curve for $\delta t = 1.0$ (red) compared to FRAP curve for $\delta t = 0.005$ (green) and exact solution (blue)



Figure 8.11: Total mass vs. time for the large time step $\delta t = 1.0$

8.4 Validation of random walk and PSE

In order to validate the PSE simulation method as described in section 7.2 as well as the random walk method outlined in section 7.1, they are applied to the test case introduced in section 8.1, again using the numerical parameter values given in table 8.1. Furthermore, comparative timings for the PSE and the random walk method are made to estimate their performance.

Recall that the FRAP value is the normalized mean concentration in the bleached box. Since the particles in a PSE simulation carry mass as their strength, it is simply the sum of the strengths of all the particles in the bleached box divided by the volume of the bleached box, thus:

$$F(t_n) = \frac{1}{\mathcal{L}^3(B)} \sum_{p \in B} v_p^h c_p^h(t_n) = \frac{1}{\#\{p : p \in B\}} \sum_{p \in B} c_p^h(t_n)$$

since all the particles have the same volume and they completely fill the bleached box. Figure 8.12 shows a comparison of the FRAP curves from the analytic solution (green) and the one obtained with the PSE algorithm (red). For the PSE, 64000 particles have been used, placed on a regular cartesian grid with grid spacing 0.1. The core size of the particles is determined from the following convergence requirement:

$$\frac{h}{\epsilon} < \text{overlap}$$

where h is the inter-particle spacing (thus 0.1) and the overlap parameter is set to 0.9, which means that the particles are required to always overlap by at least 11% of their core size. Evaluating this expression gives $\epsilon > 1/9$, so it is set to $\epsilon = 0.1118$ giving $\epsilon^2 = 0.0125$. The simulation is run for T = 600 time steps with a step size of $\delta t = 0.05$ which is smaller than the time integration stability requirement

$$\delta t < \frac{h^2}{2D} \approx 0.16667\dots$$

with the diffusion constant D = 0.03 and the inter-particle spacing h = 0.1.

A very good coincidence of the FRAP curves of the PSE algorithm and the exact solution can be observed even for small times. For times larger than 4 they hardly differ, resulting in a relative error of $-5.64 \cdot 10^{-4}$ at time 30. This is about 20 times



Figure 8.12: PSE FRAP curve (red) compared to analytic curve (green) and finite difference result (blue). The asymptotic value 0.9375 is shown in pink

better than the finite difference result with a relative error of $-1.10 \cdot 10^{-2}$ (blue curve). Both curves stay below the analytic one due to numerical dissipation.

The RMS error is defined as:

$$E_{\rm RMS} = \sqrt{\frac{1}{T} \sum_{n=1}^{T} (F(t_n) - F^*(t_n))}$$

where $F(t_n)$ is the simulated FRAP value at time $t_n = (n-1)\delta t$, $F^*(t_n)$ is the exact value from the analytic solution at the same time and T is the number of time steps. For the PSE, an RMS error value of $E_{\rm RMS} = 5.705 \cdot 10^{-3}$ results and for the finite difference code one of $E_{\rm RMS} = 2.131 \cdot 10^{-2}$, again about an order of magnitude worse. Figure 8.13 shows snapshots of the PSE simulation at different times. The visualizations have been done using the techniques described in section 7.3.

Figure 8.14 shows the comparison between the FRAP curves of the random walk code and the analytic solution. The random walk has been simulated using 10^6 particles initially placed at random locations. It is run for 4000 time steps with a step size of $\delta t = 0.005$. However, since the random walk code simply outputs the relative number of particles in the bleached box whereas the analytic curve reflects the mean concentration, some post processing is needed to make the outputs comparable. Let N be the total number of particles in a random walk simulation and p the unknown mass (or strength) per particle. Due to the fact that the volume of the bleached box is 1/16 of the total volume of the cubic domain, the initial concentration outside the bleached box is given by:

$$c_0 = \frac{16 \cdot Np}{15 \cdot V}$$





t = 0.25



Figure 8.13: Snapshots of concentration distribution for a PSE simulation of the box test case

with V being the total volume of the domain. As $c_0 = 1.0$ by construction of the problem, this can be solved for p. The instantaneous mean concentration (thus the FRAP value) in the bleached box is then given by:

$$c(t) = \frac{16 \cdot N_B(t)p}{V} = 15 \cdot c_0 \frac{N_B(t)}{N}$$

where $N_B(t)$ denotes the number of particles in the bleached box. Therefore, the output of the random walk code N_B/N has to be multiplied by 15 before comparing it to any other FRAP curve. It can be seen from figure 8.14 that the random walk curve follows the one of the analytic solution pretty well resulting in a relative error of $4.66 \cdot 10^{-3}$ at time 20 which is about 3 times less than the finite difference code with a relative error of $-1.41 \cdot 10^{-2}$. The RMS error for the random walk is $3.976 \cdot 10^{-3}$, thus of comparable accuracy as the PSE. The random walk results could be further improved by performing several runs and subsequent ensemble averaging to obtain smoother FRAP curves.

Another disadvantage of the random walk method besides the missing smoothness of the results is its slow convergence for increasing numbers of particles. According to [Cottet & Koumoutsakos (2000)], the error of a random walk simulation decreases like \sqrt{N} if the number of particles N is increasing. Thus an increase in the number of particles by a factor of 10 will only reduce the error by a factor of about 3 as can be seen from the RMS error values in table 8.2. The PSE on the other hand converges with N^2 meaning that an increase in the number of particles by a factor



Figure 8.14: Random walk FRAP curve (red) compared to analytic curve (green) and finite difference result (blue). The asymptotic value 0.9375 is shown in pink

of 10 will reduce the error by a factor of 100. Figure 8.15 shows the FRAP result for a random walk run with 10^4 particles and figure 8.16 for one with 10^5 particles. Figure 8.14 has been created using 10^6 particles. Despite their variance, the means of all curves follow the exact solution quite well, indicating that the random walk code works correctly.



Figure 8.15: Random walk FRAP curve with 10000 particles (red) compared to analytical solution (green)

Figure 8.16: Random walk FRAP curve with 100000 particles (red) compared to analytical solution (green)

To further compare the PSE and random walk methods, timings are made. Both codes are compiled (using maximum optimization -03 -fast and the Portland Group Fortran 90 compiler) and run on the same machine, a 1.4 GHz AMD Athlon

Run	RMS error
Random walk, $N = 10^4$	0.02896
Random walk, $N = 10^5$	0.01076
Random walk, $N = 10^6$	0.003976
Finite difference	0.02131
PSE, $N = 64000$	0.005708

Table 8.2: RMS error comparison

yielding 2792 MIPS⁵. Table 8.3 summarizes the results. The first two PSE runs with 64000 particles have been performed with an inter-particle spacing of 0.1 and a particle core size of 0.1118, just as before. The difference is that for the first run the code was compiled in its serial version whereas for the second run, the parallel version including MPI has been used, although run on single processor. It can be seen that the communication overhead added by MPI amounts to about 2.25 s per time step. Therefore, running in parallel is only useful for problems at least as large as this test case.

Run	CPU time per time step
PSE without MPI, $N = 64000$	$3.888\mathrm{s}$
PSE with MPI, $N = 64000$	$6.141\mathrm{s}$
PSE without MPI, $N = 10648$	$0.783\mathrm{s}$
Random walk, $N = 10^5$	$0.494\mathrm{s}$
Random walk, $N = 10^6$	$5.149\mathrm{s}$
Random walk, $N = 10^7$	$57.024\mathrm{s}$

Table 8.3: Comparative timings

The first random walk run with 10^5 particles is faster than the PSE but cannot be used without averaging, which makes more than one run necessary. The second run with 10^6 particles is comparable to the PSE method in accuracy but already slower than its serial version. The last run with 10^7 particles finally is about 15 times slower than the PSE.

In addition to this, it can be seen that the time needed for the random walk scales almost linearly with the number of particles used. The PSE even scales sub-linear due to the small number of particles compared to the computational cost of all auxiliary and bookkeeping routines.

Taking all the results of this chapter into account leads to the decision that for the simulations to come, the PSE is the algorithm of choice. The random walk will not be used any further due to the following disadvantages:

- It is slower than the PSE at comparable accuracy
- Due to its \sqrt{N} convergence, the number of particles needed to simulate diffusion in a full ER exceeds 10^{10} thus breaking all computer memory limits. (fewer particles can be used if several runs are performed and averaged but this will again cost time.)
- The random walk code is, at the current stage of development, not parallelized, excluding the (necessary) use of distributed memory machines.

 $^{^5}$ Million instructions per second. Determined using the Linux kernel time calibration loop

• The geometry handling is the random walk code is not optimized. Thus the code is $\mathcal{O}(TNM)$ where M is the number of triangles in the surface, N the number of particles and T the number of time steps. This should be compared to the PSE for which an optimized geometry preprocessor has been written making it $\mathcal{O}(TN \log N + N \log M)$.

Random walk nevertheless has been applied by [Ölveczky & Verkman (1998)] and [Sbalzarini (2001)] to the problem studied in this work and if future work makes the use of sources, sinks, tubes or separator planes necessary, one would have to revert to random walk simulations, optimizing the algorithm wherever possible.

Chapter 9

FRAP simulations

9.1 PSE simulations in all ER samples

In this chapter, the PSE simulation method as described in section 7.2 will be applied to the simulation of diffusion in reconstructed ER geometries (cf. chapter 2). The random walk code will not be used hereafter for the reasons mentioned at the end of section 8.4. The only difference to the PSE simulations conducted for the text case in section 8.4 emerges from the fact that the bleached area is no longer a box but rather the set intersection of a box with the space Ω enclosed by the ER membrane. Thus we define the bleached area as:

$$A=B\cap \Omega$$

and the FRAP value is calculated as:

$$F(t_n) = \frac{1}{\mathcal{L}^3(A)} \sum_{p \in A} v_p^h c_p^h(t_n) = \frac{1}{\#\{p : p \in A\}} \sum_{p \in A} c_p^h(t_n)$$

This definition includes the box test case in the limit A = B. To initialize the particles, all ER samples are preprocessed as described in section 7.2.5. As we wish to simulate a FRAP experiment, the proper initial condition is given by:

$$c_0^h(x) = \begin{cases} c_0 & \text{if } x \in int(\Omega) \setminus A \\ 0 & \text{if } x \in A \end{cases}$$
(9.1)

Without loss of generality, the constant initial concentration outside the bleached are is chosen to be $c_0 = 1$ as this simply corresponds to normalizing the FRAP curves. Figure 9.1 shows the initial particle distributions for all the ER samples considered. The blue points symbolize particles of initial concentration 1, the red points those of initial concentration 0 (i.e. those inside A). The assumption of a homogeneous initial concentration distribution outside the bleached area A seems feasible due to the following facts:

- 1. After transfection, the cells are incubated for 12 hours. During this time they express the green fluorescent protein which freely diffuses in the ER lumen and fills it completely. The time scale of diffusion is about 5 s. Therefore a homogeneous distribution inside the ER can be assumed after 12 hours.
- 2. The experimenter chooses "sane" cells, i.e. cells which exhibit homogeneous fluorescence and do neither under- nor over-express the protein.
- 3. The pixel intensities of the pictures taken from the microscope are normalized to the level of homogeneous concentration before any FRAP data is calculated.









Figure 9.1: Initial particle distributions for all ER samples. Blue points correspond to particles of concentration 1, red points of those of concentration 0.

Please refer to figure 5.10 in section 5.3 for shaded surface views of the geometries. The bounding boxes of all the samples are given in table 9.1 and table 9.2 states the geometries of the bleached boxes used. All boxes are of the form $[a, b] \times [c, d] \times [z_l, z_h]$ (recall that the actual bleached area is the set intersection of the bleached box and the effective ER geometry).

Sample	x_l	y_l	z_l	x_h	y_h	z_h
bip2	43.6768	171.4180	-1.7248	512.7590	421.4890	17.3869
$_{\rm clx}$	122.9590	99.3180	-2.0038	344.7140	477.0080	17.3165
erp57	142.7870	112.7570	-2.0900	513.4100	362.9470	15.5391
erp572	54.8070	102.2590	-1.8534	414.2720	512.7500	17.2687
$erp573_1$	225.6450	42.7788	-1.8912	513.1250	339.9230	18.3436
erp573_2	70.3277	131.2920	-1.9464	292.5250	512.5840	17.8942
erp573_3	160.5760	34.8480	-1.9580	512.8440	513.3540	17.9499
$erp574_1$	-1.5552	215.3650	-2.1574	419.5930	458.9890	30.0631
$erp574_2$	152.2640	90.8761	-2.0212	512.4500	270.2680	30.2559

Table 9.1: Bounding boxes of all ER sample domains

Sample	a	c	b	d
bip2	210.0	360.0	260.0	410.0
clx	175.0	125.0	225.0	175.0
erp57	190.0	260.0	240.0	310.0
erp572	250.0	125.0	300.0	175.0
erp573_1	230.0	90.0	280.0	140.0
erp573_2	225.0	450.0	275.0	500.0
erp573_3	300.0	400.0	350.0	450.0
$erp574_1$	300.0	385.0	350.0	435.0
$erp574_2$	270.0	95.0	320.0	145.0

Table 9.2: Positions and sizes of the bleached boxes B

All PSE simulations are run for the same value of the diffusion constant D = 7.5 in order to be able to study the influences of geometry. Running at a different diffusion constant would simply correspond to stretching the time axis accordingly.

This can be seen when recalling that all possible solutions to the diffusion equation are linear combinations of its elementary solution (i.e. its Green's function). This elementary solution is given by:

$$K(x,\xi,t) = \frac{1}{(4\pi Dt)^{n/2}} \exp\left[-\frac{(x-\xi)^2}{4Dt}\right]$$

and only contains Dt as a product pair. Enlarging D by a certain factor and at the same time reducing t by the same factor therefore leaves its function value unchanged. Since diffusion is a linear problem, this applies not only to the single elementary solutions but also to any linear combination of them, thus to all solutions to the diffusion equation. Performing a diffusion run at $2 \cdot D$ is therefore the same as running at D and replacing t by $1/2 \cdot t$.

The initial inter-particle spacings in x, y and z direction (termed h_1 , h_2 and h_3 , respectively) are chosen according to the geometrical resolution desired. Since the domains are quite flat compared to their x and y extensions (cf. table 9.1), the particles are placed denser in the z direction to have enough resolution. x and y spacings are chosen so as to be smaller than the smallest occurring triangle edge. This is recommended to resolve the whole geometry as described by the triangulated surface and to increase the accuracy of the boundary condition handling¹. Table 9.3 summarizes the particle grid spacings and the resulting total numbers of particles for the different simulation runs.

Sample	h_1	h_2	h_3	N
bip2	2.0	2.0	0.4	493850
clx	2.0	2.0	0.4	488054
erp57	2.0	2.0	0.4	509578
erp572	2.0	2.0	0.4	545859
erp573_1	2.0	2.0	0.4	473466
erp573_2	1.3	2.0	0.4	495373
erp573_3	2.0	2.0	0.4	335480
$erp574_1$	2.0	1.7	0.4	481922
$erp574_2$	2.0	1.6	0.4	504770

Table 9.3: Inter-particle spacings and total numbers of particles

The time step for all the simulations is chosen according to the time integration requirement requirement (see [Zimmermann, Koumoutsakos & Kinzelbach (2001)]):

$$\delta t < \frac{\min_{i=1,2,3} h_i^2}{2D}$$

where h_i stands for the inter-particle spacing in the i^{th} direction. As the smallest occurring inter-particle spacing is 0.4 for all simulations and D is constant as well, a time step of $\delta t = 0.01 < 0.01067...$ is chosen for all runs. All simulations are run for T = 2000 time steps to a final time of $t_{final} = 20.0$. The PSE kernel function's core size is determined using the requirement:

$$\epsilon > \frac{\max_{i=1,2,3} h_i}{\text{overlap}}$$

¹If the particles are too far apart, the core size ϵ becomes large and it will eventually be the case that boundary image particles of one part of the domain overlap with real particles of another part or – put differently – that particles in one ER filament directly interact with those in another one.

with overlap ≤ 1 , this is a generalized version of the requirement stated in section 7.2.1. For safety reasons, we choose overlap = 0.901 for all simulations. As the maximum occurring inter-particle distance is 2.0 for all samples, the core size is chosen to be $\epsilon = 2.222$ for all runs. Table 9.4 summarizes the parameters that are common to all simulations of this chapter.

Parameter	ϵ	δt	D	t_{final}
Value	2.222	0.01	7.50	20.0

Table 9.4: Common parameters for all PSE runs

In order to be able to compare the FRAP curves of the simulation runs amongst each other, they are normalized to their respective steady-state value. Since the bleached areas of the different ER samples contain different numbers of particles and the total number of particles also varies, the different FRAP curves will have different asymptotic levels as $t \to \infty$. This is the case due to the no-flux boundary condition² and the fact that the total mass in the domain is conserved. Initially, the total mass in the system is given by:

$$m_t = \sum_{p=1}^N v_p c_p^h = v_p (N - N_B)$$

where $N_B = \#\{p : p \in B\} = \#\{p : p \in A\}$ is the number of particles inside the bleached area. The last equality makes use of the initial condition as given by equation 9.1. The asymptotic value of the concentration is now given by homogeneously distributing this mass among all the particles, thus:

$$c_{\infty} = \frac{m_t}{Nv_p} = \frac{N - N_B}{N}$$

Table 9.5 lists the values for the ER samples under consideration. The gap between the asymptotic FRAP value and the pre-bleach level (1 in this case) is called *immobile fraction* in biology. It is observed due to the no-flux boundary condition, the conservation of mass and the finiteness of the domain. In fact, the simulations conducted in this chapter have an average immobile fraction of 2.48%, which is in nice coincidence with the experimentally observed 2 to 3% by Anna Mezzacasa.

Another interesting property to consider is the volume-filling coefficient of the ER geometry in the bleached box. This number is the fraction of the bleached box's volume that is actually taken up by the ER lumen (i.e. is filled with particles). It is formally defined as:

$$\phi_0 = \frac{\mathcal{L}^3(\Omega \cap B)}{\mathcal{L}^3(B)}$$

For the box test case considered in section 8.4, this value is $\phi_0 = 1$ by definition as the domain Ω and its bounding box are identical. For the ER geometries however, this is no longer the case and the volume-filling coefficient can be approximated as:

$$\phi_0 = \frac{N_B}{N_g} = \frac{h_1 \cdot h_2 \cdot h_3 \cdot N_B}{(b-a) (d-c) (z_h - z_l)}$$

where N_g stands for the number of grid points inside the bleached box B. The values for all the ER samples are also listed in table 9.5. It can be seen that the ER

 $^{^2{\}rm This}$ is the biologically correct boundary condition since ER lumen proteins do not spontaneously cross the ER membrane.

Sample	N	N_B	N_g	c_{∞}	ϕ_0
bip2	493850	11892	29375	0.9759	0.4048
$_{\rm clx}$	488054	12195	30000	0.9750	0.4065
erp57	509578	6396	27500	0.9874	0.2326
erp572	545859	7767	29375	0.9858	0.2644
$erp573_1$	473466	9950	31250	0.9790	0.3184
$erp573_2$	495373	12904	46550	0.9740	0.2772
erp573_3	335480	15942	30625	0.9525	0.5206
$erp574_1$	481922	12177	58000	0.9747	0.2099
$erp574_2$	504770	23284	62000	0.9539	0.3755

typically fills about 1/3 of the space in the area where the bleaching takes place. The effects of geometry on diffusion are therefore expected to be significant since diffusion is not free in space but restricted to about 1/3 of it.

Table 9.5: Asymptotic FRAP values and volume-filling coefficients

Having now all the preliminary information, the simulation results can be discussed. Figure 9.2 shows the resulting FRAP curves for all ER samples using the bleached box geometries given in table 9.2. All curves are normalized using their respective c_{∞} out of table 9.5. Therefore, they all asymptotically recover to 1.0 what makes them comparable. Figure 9.3 shows all 9 normalized FRAP curves in a single plot. As expected, they differ due to geometry influences (recall that all simulations have been done using the same diffusion constant).

As already mentioned, a way to check if the boundary condition is treated well is to consider the total mass in the system versus time. It turned out that for all runs presented in this chapter, the total mass is conserved to machine precision (plots are therefore omitted). This is possible due to the fact that (i) the particle discretization of the Laplacian chosen in section 7.2.2 is conservative and (ii) the geometrical resolution and curvature have been taken into account when choosing the inter-particle spacings and the PSE kernel core size (i.e. the cut-off).

Figure 9.5 finally shows a comparison between the FRAP curves of the box test case (drawn in red) and of the ER sample erp574_2 (in blue). Fluorescence recovery is clearly faster for the box since diffusion is free and not limited to certain geometrical structures. The half-recovery time for the ER sample is about 3 times the one for the box test case. This is consistent with above observation that the typical volume-filling coefficient is 1/3. Both curves have again been normalized by their asymptotic value, so they ultimately recover to 1.0.

Figure 9.4 shows snapshots of the concentration distribution from the simulation in the erp572 sample at times t = 0.25, 1.5, 3.0 and 5.0. They have been created according to the techniques described in section 7.3. The concentration in a slice through the center of the ER parallel to the *xy*-plane is plotted in the third direction and color coded. The area of interest around the bleached box is enlarged.





Figure 9.2: Simulated normalized FRAP curves for all ER samples.



Figure 9.3: Comparison of all normalized FRAP curves









t = 5.0 t = 5.0

Figure 9.4: Snapshots of concentration distribution for erp572 PSE simulation



Figure 9.5: Comparison of box test case (red) and erp574_2 (blue)

9.2 Influence of bleached box geometry

It is expected that not only the geometry of the domain (i.e. the ER) influences the resulting FRAP curve but also the position and size of the bleached box. Fluorescent proteins that diffuse into the bleached area have to pass certain apertures defined by the intersection of the bleached box's surface and the ER membrane. These apertures change in size and (possibly) number if the intersecting planes are shifted. To investigate this, additional simulations are performed. All of them are to use the same ER sample geometry, namely erp572. All simulation parameters except the bleached box geometry are kept the same as for the original run in this sample (cf. section 9.1). The new runs will be termed erp5722, erp5723 and erp5724. Table 9.6 contains the corresponding bleached boxes, figure 9.6 shows the initial particle distributions. All relevant information of the original run (see figure 9.6a), the first two have their bleached box placed at a different location. For the third one both the location and the size of the bleached box are changed as its edge length is 75 instead of 50.



Figure 9.6: Initial particle distributions for different bleached box geometries. Blue points correspond to particles of concentration 1, red points to those of concentration 0.

Run	a	С	b	d
erp572	250.0	125.0	300.0	175.0
erp5722	80.0	300.0	130.0	350.0
erp5723	350.0	200.0	400.0	250.0
erp5274	225.0	125.0	300.0	200.0

Table 9.6: Comparison of bleached box geometries

Table 9.7 shows the steady-state FRAP values for all runs. The FRAP curves will again be normalized by c_{∞} to be able to compare the influences of geometry. Figure 9.7 shows the resulting normalized FRAP curves from all four runs. In accordance with the theoretical predictions of [Axelrod et al. (1976)], it can be seen that position and size of the bleached box indeed do have a significant influence on the result. It seems that fluorescence recovers slower for the same value of the diffusion constant if the bleached box is larger (this would make sense as the distance a particle has to travel is larger, too). When developing new models in chapter 11, this fact will become important.

Sample	N	N_A	c_{∞}
erp572	545859	7767	0.9858
erp5722	545859	6459	0.9882
erp5723	545859	7447	0.9864
erp5724	545859	16886	0.9691

Table 9.7: Asymptotic FRAP values for the different bleached box geometries



Figure 9.7: Comparison of normalized FRAP curves for different bleached box geometries. (red: erp572, green: erp5722, blue: erp5723, black: erp5724)
9.3 Timings and parallel speed-up estimation

As the PSE simulation code is parallelized using MPI, it is insightful to conduct some benchmark runs on different parallel machines to get comparative speed-up figures. The two computers used are:

- asgard: A distributed memory machine of type Beowulf running SuSE Linux 6.3 (kernel 2.2.14-SMP) and MPICH 1.2.1. It consists of 251 nodes with 2 Intel Pentium III 500 MHz processors and 1 GB main memory each, making a total number of 502 processors, interconnected by a 100 MBit/s fast Ethernet. The theoretical peak performance is 550 MFLOPS per processor.
- prometeo: A shared memory NEC SX-5/10A vector supercomputer running SUPER-UX R11.1 and MPI/SX 6.3.0. It has 10 processors sharing 64 GB of main memory. The peak performance is 8 GFLOPS per processor.

PSE simulations of the ER sample erp572 (545859 particles) are run for 20 time steps with a step size of $\delta t = 0.01$ on both machines using different numbers of processors. The time measurement on asgard is done using the MPI_Walltime call, on prometeo, the system's tracing facility is used. When more than one processor is used, the timing for the slowest one is reported since this is the speed-limiting factor (all other processors will have to wait for the slowest one before a new time step can be started). Tables 9.8 and 9.9 give the measured elapsed time per time step for the PSE calculations, excluding all I/O and auxiliary routines. The third column lists the *speed-up* figures as defined by:

$$S(n) = \frac{t_E(1)}{t_E(n)}$$
(9.2)

where $t_E(n)$ means the execution time on *n* processors. The last column gives the parallel speed-up efficiency:

$$e(n) = \frac{S(n)}{n}$$

number of processors	seconds per time step	speed-up	parallel efficiency
1	545.928		
2	282.319	1.934	96.7%
4	241.087	2.264	56.6%
8	238.932	2.285	28.6%
16	194.855	2.802	17.5%
32	127.988	4.265	13.3%
64	63.994	8.531	13.3%

Table 9.8: Parallel scalability on asgard

number of processors	seconds per time step	speed-up	parallel efficiency
1	37.327		
2	19.542	1.910	95.5%
4	13.977	2.671	66.8%
6	11.723	3.184	53.1%

Table 9.9: Parallel scalability on prometeo

Figure 9.8 visualizes the timing results for asgard in a log-lin plot. Figure 9.9 shows a plot of the parallel speed-up as defined by equation 9.2 vs. the number of processors. The dashed blue curves indicate the ideal levels for perfect scalability. Figures 9.10 and 9.11 show the same for prometeo. It can be seen that both machines scale almost perfectly when going from 1 to 2 processors. Using a higher number of processors on asgard does not really improve things any more until going above 16 CPUs. On prometeo, a more or less steady speed-up is observed even for 4 and 6 processors. This is suspected to be caused by the fact that asgard is a distributed memory machine whereas prometeo has a shared memory architecture. MPI communication on asgard has to pass the slow and potentially congested Ethernet links, on prometeo it just consists of copying within the main memory. Therefore it can be assumed that MPI does not add any significant overhead to the program on prometeo (see tracing results below), on asgard however it does. This accounts for the different scaling behavior of the two machines. To explain the deviation from the blue optimum line even on prometeo, one has to consider the quality of load balancing.



Figure 9.8: CPU time per time step for asgard (red) vs. ideal line (blue)



Figure 9.9: Parallel speed-up on asgard (red) compared to ideal speedup (blue)



Figure 9.10: CPU time per time step for prometeo (red) vs. ideal line (blue)



Figure 9.11: Parallel speed-up on prometeo compared to ideal speedup (blue)

The goal of load balancing is to equally distribute the amount of work to be done among a certain number of processors. If one processor gets assigned more work than its colleagues, they will have to sit idle and wait for it at the end of the time step when the results are communicated.

One parameter to measure the quality of load balancing is the *load distribution coefficient*:

$$\ell = \frac{\max_i t_i}{\min_i t_i} \qquad i = 1, \dots, n$$

where t_i is the time spent on processor *i* out of *n*. Perfect load balancing would thus result in value of $\ell = 1$. Table 9.10 lists the numbers for prometeo. In fact, the load balance is almost perfect for 2 processors but then deteriorates rapidly. Already using 6 processors causes an imbalance of a factor of 10. This explains the speed-up curve shown in figure 9.11. Consider the following snapshot of total execution time (i.e. for all 20 time steps) on two neighboring processors of the 4 processor case:

processor	pse	diag
1	103.682 sec	175.214 sec
2	279.548 sec	0.248 sec

It is noticeable that the difference of the two numbers for the **pse** routine pretty much corresponds to the time spent in **diag** on the first processor. **diag** is the subroutine that is called directly after **pse** and that collects the results from the different processors to calculate the diagnostics (e.g. the FRAP value). It is therefore the first MPI communication call after the PSE calculation. Above figures show that processor 1 – which gets assigned less work and therefore completes the PSE earlier – has to wait for the results of processor 2 to arrive, thus spending a lot of idle CPU time in the subroutine **diag** which normally only takes about 0.2 s.

The reason for such a poor load balance is the use of MPI's built-in cartesian domain decomposition routines. They only take into account the edge lengths of the bounding box but not the particle densities. Since the bounding box of the present ER problem is very flat compared to its x and y dimensions (cf. table 9.1) and the particles are not homogeneously distributed inside it (but only where there is ER geometry), this is not a good way to decompose the domain. Some processors might not get any particles at all while others might have densely filled sub-domains. Load balancing thus is an issue to consider when trying to improve the code's performance. A custom load balancing algorithm would not equally distribute the volume of the bounding box but rather the number of particles.

number of processors	$\min_i t_i (i)$	$\max_i t_i$ (i)	ℓ
2	17.525(2)	19.542(1)	1.115
4	5.184(1)	13.977(2)	2.696
6	1.170(1)	11.723(4)	10.020

Table 9.10: Load balance on prometeo

Using a single processor only, prometeo is about 14.6 times faster than asgard. This is in good agreement with the fact that its peak performance is 14.5 times the one of asgard, meaning that on both machines the code sustains about the same fraction of the peak performance. To estimate this fraction, the program is traced on prometeo using the **-ftrace** compiler option. The resulting output for a representative run is:

/tmp.speed/ivo/PSEtrace ftrace -f ftrace.out.0.0

FLOW TRACE ANALYSIS LIST

Execution : Tue Feb 5 09:42:35 2002 Total CPU : 0:13'03"761

PROG.UNIT	FREQUENCY	EXCLUSIVE		AVER.TIME	MOPS	MFLOPS	V.OP	AVER.	I-CACHE	O-CACHE	BANK
		TIME[sec](%)	[msec]			RATIO	V.LEN	MISS	MISS	CONF
		700 100(0	>	00000 405							
pse	20	760.123(9)	(.0)	38006.135	3647.5	2004.6	99.10	134.9	0.0706	0.9937	32.31/1
init	1	12.147(1.5)	12146.724	85.2	0.4	1.92	5.2	1.9725	0.6545	0.0000
do_10	1	8.541(1.1)	8540.770	92.7	0.7	2.69	12.5	1.2/81	0.0673	0.0000
sortp	20	1.908((0.2)	95.408	700.4	69.3	84.18	255.9	0.0004	0.4233	0.0013
updatep	20	0.641((0.1)	32.071	16/1.5	34.0	92.67	255.9	0.0024	0.1088	0.0027
bc	21	0.183((0.0)	8.692	4631.0	378.2	99.35	255.8	0.0004	0.0002	0.0005
diag	21	0.084((0.0)	3.999	4954.6	819.3	99.22	255.5	0.0090	0.0031	0.0000
partalloc	63	0.081((0.0)	1.291	4896.5	0.0	98.70	255.9	0.0015	0.0007	0.0020
pwrite	189	0.039((0.0)	0.206	130.2	0.0	0.84	25.7	0.0053	0.0012	0.0000
readctrl	1	0.006((0.0)	5.921	127.7	0.0	1.34	40.4	0.0009	0.0002	0.0000
pse3d	1	0.002((0.0)	2.027	26.8	0.3	11.53	120.3	0.0011	0.0004	0.0000
aniso_coeff	f 1	0.002((0.0)	1.664	2641.6	0.0	99.35	255.9	0.0000	0.0000	0.0000
mktable	1	0.002((0.0)	1.659	8415.3	4897.4	99.51	255.8	0.0000	0.0000	0.0000
chkabort	20	0.001((0.0)	0.056	16.7	0.0	0.00	0.0	0.0007	0.0003	0.0000
mkmsh	1	0.001((0.0)	0.831	58.1	0.9	10.41	35.7	0.0003	0.0001	0.0000
finalise	1	0.001((0.0)	0.807	59.3	0.0	0.06	1.1	0.0002	0.0002	0.0000
inputarg	1	0.000((0.0)	0.317	155.4	0.0	0.06	8.0	0.0000	0.0000	0.0000
alloccp	21	0.000((0.0)	0.015	27.2	0.0	0.00	0.0	0.0001	0.0001	0.0000
substop	220	0.000((0.0)	0.001	30.9	0.0	0.00	0.0	0.0000	0.0000	0.0000
substart	219	0.000((0.0)	0.001	40.5	0.0	0.00	0.0	0.0000	0.0000	0.0000
chkbc	1	0.000((0.0)	0.114	29.7	0.1	0.36	3.0	0.0001	0.0000	0.0000
meshalloc	1	0.000((0.0)	0.084	49.9	0.1	6.89	72.5	0.0000	0.0000	0.0000
rmabort	1	0.000((0.0)	0.053	14.1	0.0	0.00	0.0	0.0000	0.0000	0.0000
mkmshsz	1	0.000((0.0)	0.052	30.7	0.3	0.38	1.0	0.0000	0.0000	0.0000
uppercase	17	0.000((0.0)	0.002	63.3	0.0	0.00	0.0	0.0000	0.0000	0.0000
defaults	1	0.000((0.0)	0.028	93.4	0.0	0.46	6.0	0.0000	0.0000	0.0000
const	1	0.000((0.0)	0.012	24.6	4.7	0.00	0.0	0.0000	0.0000	0.0000
mp_time	2	0.000((0.0)	0.004	15.7	0.3	0.00	0.0	0.0000	0.0000	0.0000
defneigh	1	0.000((0.0)	0.005	70.1	0.0	0.00	0.0	0.0000	0.0000	0.0000
updbc	1	0.000((0.0)	0.004	33.7	0.0	0.00	0.0	0.0000	0.0000	0.0000
neighproc	1	0.000(0.0)	0.003	31.1	0.0	0.00	0.0	0.0000	0.0000	0.0000
total	 871	783.762(10)	0.0)	899.841	3545.0	1944.5	99.02	135.0	3.3439	2.2542	32.3237

It can be seen that the PSE subroutine takes 97% of all execution time and almost all main memory (as indicated by BANK CONF), making it clearly the most performance relevant part of the code. It is called 20 times (according to the 20 time steps), takes around 38 seconds per call and has a sustained performance of 2 GFLOPS, thus 25% of the peak performance. It is interesting to note that the subroutine mktable gets 4.9 GFLOPS (61% of peak). This is due to the fact that it vectorizes almost perfectly (99.51% vectorized with an average vector length of 255.8). The task of this subroutine is to build a look-up table for the exponential function to speed-up the numerous evaluations of the PSE kernel function. Also the PSE core routine pse vectorizes 99.1%, but the average vector length is only 134.9 while 256.0 would be the maximum. On average, this routine thus uses only 52.7% of the vector registers, which explains the difference in the operation count. It would be possible to increase the vector length of the core routine by building particle interaction lists prior to doing the actual interactions. This way, all interactions of one particle would be combined into a huge vector and one could possibly improve the over-all performance by a factor of two. Also the communication and memory allocation subroutines updatep, bc, diag and partalloc vectorized almost completely.

Another interesting detail is to notice that the routines init and do_io produce a significant amount of instruction cache misses. This is due to the fact that they do I/O on changing data. However they are only called once (init at the beginning, do_io at the end of the simulation) so it does not really harm.

Chapter 10

Comparison to experiments

The simulations conducted in chapter 9 will be validated against experiments in this chapter. A FRAP experiment as well as a stack of micrograph sections at $\Delta z = 0.1 \,\mu$ m has been made from the same cell by Anna Mezzacasa according to the following protocol:

Cell culture:

VERO cells were cultured essentially as described in (Pelkmans et. al., Inst. of Biochemistry, ETHZ), using a 100 mm tissue-culture grade plastic dish at 37°C with 5% CO₂ in DMEM cell culture media containing 10% FCS, and then passaged to 18 mm glass cover slips. At 70% confluence, a reporter gene containing the ER targeting signal sequence fused to GFP and the ER retention sequence (KDEL) was transiently transfected using Superfect (Sigma) for expression in mammalian cells. 12 – 16 hours post transfection, live cells were exchanged into CO₂-independent buffer for imaging.

Confocal microscopy:

Live cells were maintained at 37°C during the FRAP experiment using a Zeiss stage warmer. Confocal images were acquired with an inverted Zeiss LSM 510 (X63 PlanApo Nikon objective, N.A. 1.4) using the 488 nm argon laser line and detected using a BP505-530 emission filter (488/543 dichroic) (pinhole setting at 1 Airy unit, image pixel size 80 nm). Images were acquired as 8-bit tiff files (512×512 pixel frame), analyzed using Openlab 3.0.4 and processed in Photoshop 5.0 (Adobe).

z-series and Fluorescence Recovery After Photobleaching (FRAP):

Prior to FRAP, 55 0.1 $\mu \rm m$ optical z-sections were collected and processed for 3D reconstruction. Experimental details of FRAP have been described in the protocols in [Sbalzarini (2001)]. To a wild type cell expressing the fusion protein: ss-GFP-KDEL, image scans of the entire cell were acquired at low laser power (30% power, 0.3% transmission). A defined region of interest (ROI) (5×5 $\mu \rm m^2)$ was photobleached at full laser power (100% power, 100% transmission, 50 ms sampling interval); recovery of fluorescence was monitored by scanning the whole cell at low laser power (30% power, 0.3% transmission).

As indicated in the protocol, the cell was transfected to express a protein called ssGFP-KDEL. Hereby, "ss" stands for "signal sequence", a polypeptide needed for the ribosomes to recognize the transcript and start synthesizing it into the ER lumen. "GFP" means the familiar green fluorescent protein and "KDEL" is a targeting sequence for ER resident proteins, which causes the protein to be retained in the ER once it is synthesized (or being transferred back to the ER if it escapes). The ssGFP-KDEL protein is soluble meaning that it freely diffuses within the ER lumen (as opposed to membrane-bound proteins). This makes the results of the simulations comparable as the PSE algorithm simulates diffusion *inside* the reconstructed ER surface rather than diffusion *on* that surface. Figure 10.1 shows a sample micrograph of a VERO¹ cell with the ER shining in green. After filtering the data with a Gaussian filter of kernel size 46.576, the reconstruction is done according to chapter 2 with voxels of size 66.54×27.0 . Figure 10.2 shows a shaded surface view of the resulting triangulated 3D surface.

¹Cell from green monkey intestine epithelium





Figure 10.1: Micrograph of the ER sample 8s (courtesy of Anna Mezza-casa)

Figure 10.2: 3D reconstruction of ER surface

The bounding box of the triangulated ER geometry is:

$[6146.58, 27349.1] \times [4502.35, 31204.9] \times [-40.2303, 40.2303]$

In the lateral directions x and y, one unit of length corresponds to $1.2 \,\mathrm{nm}^2$, in the z direction the unit is $67.1 \,\mathrm{nm}^3$. The lateral resolution of the microscope is $220 \,\mathrm{nm}$, which means that all the pictures have been taken at a 2.75-fold oversampling. The depth resolution of the microscope is $500 \,\mathrm{nm}$. Since the slices are taken at $100 \,\mathrm{nm}$ distance, they actually overlap, leading to a slightly blurred representation in the third direction. Due to these resolution limitations and the fact that the unit length is two orders of magnitude smaller than for all the runs in section 9.1, the global geometry tolerance TOL is increased from 10^{-10} to 10^{-8} .

To get a meaningful validation, two runs (and experiments) for different bleached boxes are performed using this geometry. Table 10.1 lists the coordinates of the bleached boxes used in the experiments and the subsequent simulations. The first run, termed 8s features a small bleached box at the periphery of the cell, the second run (8.2) a larger one closer to the nucleus.

Run	a	С	b	d
8s	23500	25000	26000	27500
8.2	16500	13000	20500	17000

Table 10.1: Comparison of bleached box geometries

²One pixel of the micrograph image represents a 80×80 nm square, the image has 320×407 pixel and the bounding box of the triangulation is 21202×26703 .

 $^{^3 {\}rm The}~z{\rm -distance}$ between each of the 55 micrograph slices is $0.1\,\mu{\rm m}$ and the z-length of the bounding box of the geometry is 80.46.

For the subsequent FRAP curve, the fluorescence intensity in the originally bleached box was measured at a sampling interval of 50 ms. The quantification has been done by Anna Mezzacasa after properly normalizing the images and subtracting any background noise. The resulting measured data are listed below and the FRAP curves are shown in figure 10.3.

time	FRAP 8.2	FRAP 8s
0.00	0.0000000	0.000000000
0.05	0.642359994	0.768091498
0.10	0.761357527	0.85031668
0.15	0.825385128	0.875191667
0.20	0.85971692	0.925180215
0.25	0.891158161	0.946877274
0.30	0.914321773	0.946250857
0.35	0.925693525	0.946139788
0.40	0.924324516	0.978810841
0.45	0.937138416	0.96555209
0.50	0.954323345	0.985550568
0.55	0.946796145	0.999120908
0.60	0.937829731	0.98808123
0.65	0.958437983	1.029331826
0.70	0.952619479	1.033946096
0.75	0.962860212	
0.80	0.962289884	
0.85	0.961190977	
0.90	0.983247206	
0.95	0.96625929	
1.00	0.973192383	
1.05	0.976437893	
1.10	0.965600989	



Figure 10.3: Measured FRAP curves for experiments 8s (red) and 8.2 (blue) (courtesy of Anna Mezzacasa)

One problem with the protein used for the experiments is its fast diffusion due to its small size of ~ 30 kDa and the free space diffusion. This means that the FRAP curve will recover very quickly, asking for a large diffusion constant in the simulations. According to the stability criterion given in section 9.1 this asks for a very small time step which in turn requires a large number of time steps to be simulated. To keep the simulations computationally feasible, the number of particles therefore is drastically reduced compared to all previous runs. The inter-particle spacings are chosen to be $h_{1,2,3} = [100.0, 100.0, 20.0]$ which results in a total number of 46486 particles. This is roughly 1/10 of the numbers used so far. According to the timing results given in section 9.3 this means that one time steps will complete in about 1.5 s on prometeo. The required number of 450000 time steps will thus approximately need 187.5 hours on 4 processors (amounting to 750 CPU-hours !).

Table 10.2 lists the resulting numbers of particles inside the bleached area (N_B) , the number of grid points inside the bleached box (N_g) as well as the asymptotic steady-state FRAP value (c_{∞}) and the volume-filling coefficient (ϕ_0) . Definitions and explanations of all of these properties have been given in section 9.1. It can be seem from the volume-filling coefficient that the second run takes place in an area where the ER is much denser since it is closer the the nucleus.

Sample	N	N_B	N_g	c_{∞}	ϕ_0
8s	46486	112	2500	0.9976	0.0448
8.2	46486	3556	6400	0.9235	0.5556

Table 10.2: Asymptotic FRAP values and volume-filling coefficients

According to the algorithms presented in section 7.2.5, the particles are initialized and placed inside the reconstructed ER geometry. Figure 10.4 shows the corresponding plots for both examples. As before, blue points signify particles of initial concentration 1, red points such of concentration 0.



Figure 10.4: Initial particle distributions for comparative simulations. Blue points correspond to particles of concentration 1, red points to those of concentration 0.

The diffusion constant is arbitrarily set to D = 750.0 for all runs. This requires the time step size to be less than 0.267 for the simulation to be stable. A time step of $\delta t = 0.1$ is chosen for all runs. The simulations are performed using the techniques introduced in chapter 7 in order to obtain simulated FRAP curves. Using MATLAB, the normalized FRAP curves are then fitted to the experimental data such as to minimize the quadratic error. This is done by allowing time scaling only which is –

as shown in section 9.1 – equivalent to varying the diffusion constant. Figure 10.5 shows the results for both examples along with the resulting time stretching factors s_t . Both factors are about $s_t = 2 \cdot 10^{-5}$, indicating that both examples have been run using the same diffusion constant. This is in fact true and can be considered a consistency check for this validation.



Figure 10.5: Simulated FRAP curve (red) compared to experimental measurement data (blue circles). The time stretching factors are listed below the corresponding figure.

The fitting error for sample 8s is 0.021 and thus much larger than the one for the 8.2 sample of $6.389 \cdot 10^{-4}$. This has different reasons, among them:

- Fluorescence recovery in the first example is very fast, resulting in an almost complete recovery after only three measurement intervals. This leads to large measurement uncertainties since the speed of the process under observation challenges the measurement equipment.
- Due to the large inter-particle spacing that had to be chosen in order to finish the simulation runs on time, the resolution is quite poor having only 112 particles inside the bleached box (compared to 3556 for the 8.2 example).
- The simulation is done in an area where the ER is not very dense. This leads to a high sensitivity to geometry errors due to picture noise or reconstruction uncertainties.
- The measured FRAP curve for the first example recovers to values > 1, which is not possible. This indicates a generally higher level of measurement uncertainty for this run.

Taking these points into account, one has to state that the second run (8.2) is more significant and reliable than the first one. The simulation should therefore be validated against the 8.2 run and the 8s run should be taken as a consistency check only. It is noticeable that the resulting time stretching factors are about the same, despite the larger fitting error for 8s. This could mean that the method presented is quite robust against measurement errors and simulation resolution. However, to make a statement about this, more samples have to be considered of course. Another explanation for the deviations in the first case could be that diffusion in reality is anisotropic and thus the isotropic simulation results fail to match it exactly. Running anisotropic simulations with different ratios of anisotropy would be needed to address this question in future work (cf. chapter 13). Using the time stretch information and the known length scale of the micrographs it is now even possible to calculate a numerical estimation for the diffusion constant. The simulation has been run with a diffusion constant of 750 (units length)²/(unit time). From the fitting of the resulting curve to the experimental data we know that (unit time)= $2 \cdot 10^{-5}$ s and from the microscope settings (see page 94) we know (unit length)=1.2 nm for the lateral directions. Inserting this yields a diffusion constant of $D = 54.0 \,\mu\text{m}^2/\text{s}$. To compare with, the reported diffusion constant of pure GFP in water is $87 \,\mu\text{m}^2/\text{s}$.

In conclusion, one can say that the simulations have been shown to perform reasonably well, even under unfavorable conditions. They generally exhibit good coincidence with experimental data and make estimations of the diffusion constant possible.

Chapter 11

Towards a novel FRAP data model

It has been shown in chapter 10 that it is possible to obtain numerical values for the diffusion constant by fitting the FRAP curve from a simulation in the very same geometry to the measured intensity values. However, it would be too time consuming and require too much knowledge and resources to do so on a regular basis in productive laboratory work. Therefore, we are looking for a mathematical model that describes the time behavior of diffusion in the ER (i.e. the FRAP curve as a function of time). The parameters of this model will then be identified using the simulated FRAP curves for which we know the value of the diffusion constant. The model should then be able to interpolate and – if based on first physical principles – also extrapolate to new experimental FRAP curves of unknown diffusion constant in order to extract the value thereof.

11.1 Review of current models

To start with, briefly recall the standard models used today as already presented in [Sbalzarini (2001)]. Note that besides the ones presented hereafter, other models for special cases such as strip bleaching experiments exist ([Ellenberg et al. (1997)]). The first standard model is the *exponential model* emerging from the solution of the diffusion equation on a square 2D plate assuming homogeneous and isotropic diffusion. The fluorescence intensity in an initially bleached square region will then recover as:

$$F(t) = F_{\infty} \left(1 - e^{-\alpha Dt} \right) \tag{11.1}$$

where F(t) is the intensity (FRAP value) at time t, F_{∞} is the asymptotic intensity for $t \to \infty$ and α is a model parameter.

Another model commonly used is a purely *empirical correlation* originally proposed by [Dayel, Hom & Verkman (1999)] and [Reits & Neefjes (2001)]. It is defined by:

$$F(t) = F_a + \frac{\left[F_a + R\left(F_0 - F_a\right)\right] \left(t/t_{1/2}\right)^{\alpha}}{1 + \left(t/t_{1/2}\right)^{\alpha}}$$
(11.2)

where F_0 is the intensity just before bleaching, F_a is the intensity just after bleaching and α and $t_{1/2}$ are model parameters. The *mobile fraction* R is defined as:

$$R = \frac{F_{\infty} - F_a}{F_0 - F_a}$$

where F_{∞} is the asymptotic fluorescence value for $t \to \infty$.

Both models described above lack a solid physical foundation and are thus questionable in their extrapolation capabilities. Due to this fundamental problem, there are a number of practical insufficiencies. The models are for example not capable to account for the fact that different ER geometries will yield different FRAP curves, even at the same value of the diffusion constant (cf. figure 9.3 in section 9.1). They will thus predict a difference in the apparent diffusion constant even though there is none. This is simply due to the fact that they do not contain any information about the geometry at hand (more accurate: about the restriction of diffusion to the geometry at hand) and will interpret all variations in the FRAP curve as caused by the diffusion constant.

Another problem with these models is that they do not include any information about the position or the size of the bleached box. As shown in section 9.2 this is however an important piece of information that indeed does have an influence on the resulting FRAP curve (see figure 9.7). Current models will thus predict a dependence of the diffusion constant on the geometry of the bleached box. This however is completely unphysical and an artifact of the model used. The diffusion constant is physically given by nature and it cannot be that it depends itself on the method used to measure it. New models should thus take this into account.

11.2 New models

Now knowing the currently used models and their insufficiencies, two new models as developed during this project are presented.

11.2.1 A power law model

In order to have a model that takes the ER's local geometry into account, the concepts of diffusion on fractal sets are used. It has been shown in section 6.2 that the dimension of the walk is a sufficient parameter to capture the influences of geometry on the time behavior of diffusion. It has also been shown in the same section that the mean square displacement of a particle undergoing Brownian motion on a fractal is given by:

$$\mathsf{E}\left(\left|X\left(t+\delta t\right)-X\left(t\right)\right|^{2}\right) \asymp \delta t^{2/d_{w}}$$

where d_w is the dimension of the walk for the specific geometry under consideration. Using this parameter in a model would thus incorporate the sought-after geometrical information and moreover it does this very efficiently using just a single parameter. For unrestricted diffusion in a Euclidean shape, d_w is always equal to 2. For fractals (or geometrically restricted diffusion), we have $d_w \neq 2$. For the first model, we will thus assume that the FRAP value is proportional to t^{α} for some parameter α . In order to meet the global model constraints that the FRAP curve starts at (0,0) and asymptotically goes to a steady-state value F_{∞} , we set:

$$F(t) = F_{\infty} \left(1 - t^{-\alpha} \right)$$

The problem now is that this would be $-\infty$ for t = 0, so we shift the origin by one, replacing t with (t+1). To account for the fact that the mean square displacement is only proportional to t^{2/d_w} but not equal, we also introduce a proportionality factor β . The final *power law model* is thus given by:

$$F(t) = F_{\infty} \left(1 - \left(\beta t + 1\right)^{-\alpha} \right)$$
(11.3)

with α and β being its parameters. Compared to the current models of section 11.1, this now incorporates some physical knowledge about the process and also takes into account the influences of the local ER geometry on diffusion. However, it still fails to include any information about the geometry of the bleached box and will thus still be insufficient.

11.2.2 A second order physical model

The task is to now include some information about the bleached box into the model. This is done by starting from a very simple and basic physical reasoning. Consider the situation in the vicinity of the bleached box as depicted in figure 11.1. Instead of looking at diffusion of fluorescent protein *into* the box, consider the complementary problem of diffusion of bleached protein *out of* the box. This process is analogous to the original one and also happens in reality, governed by the same laws and the same diffusion coefficient. Considering it has however the advantage that the asymptotics of the model will automatically come out right whereas a model for diffusion *into* the box would need a fix (just as the previous power law model needed) to prevent it from diverging to infinity.



Figure 11.1: Geometrical situation around the bleached box

Let a, b and c be the lengths of the edges of the bleached box in all three dimensions. Without loss of generality, assume that the initial concentration of bleached protein inside the box is 1. The total mass of bleached protein in the box is thus equal to *abc*. After some time elapsed, the concentration front will have traveled a mean distance of \overline{x} in both the x and y direction due to diffusion. The new volume in which the bleached proteins can distribute themselves is thus given by $(a + 2\overline{x})(b + 2\overline{x})c$. Assuming a homogeneous distribution and using the conservation of mass, the new mean concentration in the bleached box hence becomes:

$$\frac{abc}{(a+2\overline{x})(b+2\overline{x})c}$$

The second assumption to be made is that there is a finite and constant number of places for particles to be. This means that if a bleached protein moves out of the bleached box, a fluorescent one will come in for it in exchange. Under this assumption, the fluorescence intensity (or the FRAP value) is given by the complement of above concentration. At the same time we also rescale the model to an asymptotic level of F_{∞} instead of 1, so it becomes:

$$F(\overline{x}) = F_{\infty} \left(1 - \frac{ab}{(a+2\overline{x})(b+2\overline{x})} \right)$$

To also take into account the influences of the local ER geometry, we set $\overline{x}^2 = \alpha^2 t^{2\beta}$ using the same concepts and proofs of diffusion on fractal sets as for the power law model. Inserting this into above equation leaves us with the final formulation of the *second order physical model*:

$$F(t) = F_{\infty} \left(1 - \frac{ab}{ab + 2(a+b)\alpha t^{\beta} + 4\alpha^2 t^{2\beta}} \right)$$
(11.4)

with parameters α and β . This model now includes both information about the ER geometry (using the fractal concepts) and the bleached box (in *a* and *b*). Moreover, it is based on first physical principles and can thus be expected to even have some extrapolation capabilities. It is termed *second order* physical model since it contains a second order polynomial in αt in the denominator.

To get the half-time of recovery from this model, one simply sets $F(t) = \frac{1}{2}F_{\infty}$ and solves for t taking into account that t > 0. This yields:

$$t_{1/2} = \left[\frac{\sqrt{(a+b)^2\alpha^2 + 4ab\alpha^2} - (a+b)\alpha}{4\alpha^2}\right]^{1/\beta}$$
(11.5)

This can now be used to compare the results to the ones of the empirical model given by equation 11.2 or to draw certain physical conclusions (cf. also section 11.4).

11.3 Identification of the model parameters

In order for the models described in the previous sections to be of any use, their parameters have to be identified (i.e. assigned some numerical value). This is done by fitting them to the simulated FRAP curves obtained in section 9.1.

11.3.1 The gradient descent algorithm

To fit the (nonlinear) models to simulation or experimental data, a gradient descent algorithm is used. Let $F^*(t_n)$ be the given data to match at time $t_n = (n-1)\delta t$ and $F(t = t_n, \alpha, \beta)$ the corresponding value of the model under consideration. Now define the matching error as:

$$E = \sum_{n=1}^{n_{max}} E(t_n) = \frac{1}{2} \sum_{n=1}^{n_{max}} \left[F^*(t_n) - F(t = t_n, \alpha, \beta) \right]^2$$
(11.6)

The problem now consists of finding the pair of model parameters (α, β) that minimizes the error E. This is done by following the gradient of E with respect to α and β :

$$\frac{\partial E}{\partial \alpha} = \sum_{n=1}^{n_{max}} \frac{\partial E(t_n)}{\partial \alpha}$$
$$\frac{\partial E}{\partial \beta} = \sum_{n=1}^{n_{max}} \frac{\partial E(t_n)}{\partial \beta}$$

The algorithm thus proceeds as follows:

Algorithm 11.1 (Gradient descent).

Step 1: Choose starting values α_0 and β_0

Step 2: Iterate for $k = 1, 2, \ldots$

2.2:

2.1: Update α and β according to:

$$\alpha_{k} = \alpha_{k-1} - \eta_{\alpha} \sum_{n=1}^{n_{max}} \frac{\partial E(t_{n})}{\partial \alpha_{k-1}}$$

$$\beta_{k} = \beta_{k-1} - \eta_{\beta} \sum_{n=1}^{n_{max}} \frac{\partial E(t_{n})}{\partial \beta_{k-1}}$$
112-139
If $\frac{\partial E}{\partial \alpha_{k-1}} < TOL$ and $\frac{\partial E}{\partial \beta_{k-1}} < TOL$: exit
149-151

The learning rates η_{α} and η_{β} are chosen to maximize convergence but prevent instability. The larger they are, the faster the algorithm will converge but at some point is starts to become unstable or to miss the minimum because it takes too large steps. Above algorithm is implemented as a Fortran program called nlfit.f90. See appendix C.4.2.1 for its source code.

11.3.2 Model gradients

To be able to apply the gradient descent algorithm, the gradients of the error function for all models are needed for all time steps. For the derivation of some of the gradients, the following law of logarithmic differentiation has been used:

$$\frac{\partial}{\partial\beta} \left(t^{c\beta} \right) = \frac{\partial}{\partial\beta} \left(e^{c\beta \log t} \right) = c \log t e^{c\beta \log t} = c t^{c\beta} \log t$$

All the model equations including their gradients will in turn be given hereafter.

Exponential model

$$F(t) = F_{\infty} \left(1 - e^{-\alpha Dt} \right)$$
$$\frac{\partial E(t_n)}{\partial \alpha} = \frac{\partial E(t_n)}{\partial F(t)} \frac{\partial F(t)}{\partial \alpha} = \left(F(t_n) - F^*(t_n) \right) \left(F_{\infty} - F(t_n) \right) t_n$$

Empirical model

To represent the situation of the simulations, the empirical model has to be taken with $F_a = 0$ and $F_0 = 1$. Therefore, we have $R = F_{\infty}$ and the model becomes: 35-36

$$F(t) = \frac{F_{\infty} (t/t_{1/2})^{\alpha}}{1 + (t/t_{1/2})^{\alpha}}$$

$$\frac{\partial E(t_n)}{\partial \alpha} = F_{\infty} (F(t_n) - F^*(t_n)) \left[(t_n/t_{1/2})^{\alpha} \log (t_n/t_{1/2}) \left[1 + (t_n/t_{1/2})^{\alpha} \right]^{-1} - (t_n/t_{1/2})^{2\alpha} \left[1 + (t_n/t_{1/2})^{\alpha} \right]^{-2} \log (t_n/t_{1/2}) \right]$$

$$\frac{\partial E(t_n)}{\partial t_{1/2}} = F_{\infty} (F(t_n) - F^*(t_n)) \left[(t_n/t_{1/2})^{\alpha} \left[1 + (t_n/t_{1/2})^{\alpha} \right]^{-2} \frac{\alpha}{t_{1/2}^2} (t_n/t_{1/2})^{\alpha-1} - \frac{\alpha}{t_{1/2}^2} (t_n/t_{1/2})^{\alpha-1} \left[1 + (t_n/t_{1/2})^{\alpha} \right]^{-1} \right]$$

Power law model

$$F(t) = F_{\infty} \left(1 - (\beta t + 1)^{-\alpha} \right)$$
$$\frac{\partial E(t_n)}{\partial \alpha} = (F(t_n) - F^*(t_n)) \left(F_{\infty} - F(t_n) \right) \log \left(\beta t_n + 1 \right)$$
$$\frac{\partial E(t_n)}{\partial \beta} = (F(t_n) - F^*(t_n)) \left(F_{\infty} - F(t_n) \right) \frac{\alpha t_n}{\beta t_n + 1}$$

Second order physical model

$$\begin{split} F(t) &= F_{\infty} \left(1 - \frac{ab}{ab + 2 (a + b) \alpha t^{\beta} + 4\alpha^2 t^{2\beta}} \right) \\ \frac{\partial E(t_n)}{\partial \alpha} &= F_{\infty} ab \left(F(t_n) - F^*(t_n) \right) \left[ab + 2 \left(a + b \right) \alpha t_n^{\beta} + 4\alpha^2 t_n^{2\beta} \right]^{-2} \\ &\cdot \left[2 \left(a + b \right) t_n^{\beta} + 8\alpha t_n^{2\beta} \right] \\ \frac{\partial E(t_n)}{\partial \beta} &= F_{\infty} ab \left(F(t_n) - F^*(t_n) \right) \left[ab + 2 \left(a + b \right) \alpha t_n^{\beta} + 4\alpha^2 t_n^{2\beta} \right]^{-2} \\ &\cdot \left[2 \left(a + b \right) \alpha t_n^{\beta} \log t_n + 8\alpha^2 t_n^{2\beta} \log t_n \right] \end{split}$$

11.3.3 Properties of the error functions

To get an idea of what the error functions for the different model fits looks like, the parameter space is sampled on a regular cartesian grid of 500×500 grid points for the models having two parameters and on 500 linear interval points for the exponential model having just one parameter. At each grid point (α_i, β_j) , the RMS matching error is calculated as:

$$E_{\rm RMS} = \sqrt{\frac{1}{n_{max}} \sum_{n=1}^{n_{max}} \left[F^*(t_n) - F(t = t_n, \alpha_i, \beta_j)\right]^2}$$
(11.7)

and its logarithm is plotted. The ER sample clx has been taken for all visualizations. Figure 11.2 shows the resulting functional shape for the exponential model, figures



Figure 11.2: Logarithmic RMS matching error for the exponential model



Figure 11.4: Logarithmic RMS matching error for the power law model



Figure 11.3: Logarithmic RMS matching error for the empirical model



Figure 11.5: Logarithmic RMS matching error for the second order physical model

11.3 to 11.5 show the contour plots for the empirical, power law and second order physical model, respectively.

It can be seen from the plots that within the area of interest there is only one minimum. This justifies the application of a gradient descent algorithm as there is no danger for it to get stuck in local minima.

Figure 11.6 shows the model matching error for the current empirical model (blue dashed line) and the second order physical model (red) as a function of the iteration step number of the gradient descent algorithm. To be able to see anything at all, logarithmic scaling is used on both axes.

Another interesting detail is that the error function for the power law model resembles the Rosenbrock banana function¹ in two characteristic properties:

- 1. It is very ill-conditioned with a narrow valley that has a soft slope on one side of the minimum and a very steep one on the other side (notice the logarithmic scaling).
- 2. The valley makes a banana-like turn.

These properties make the function very difficult to handle for optimization algorithms as they might overshoot, miss the minimum or fail to follow the narrow

 $^{^1}A$ popular multimodal test function for minimization algorithms. In its 2-parameter form: $f(x)=100(\beta-\alpha^2)^2+(1-\alpha)^2$



Figure 11.6: Model matching error for empirical model (blue dashed) and second order physical model (red). Notice the double logarithmic scaling.

turning valley. Therefore, Rosenbrock's function is very popular as a test function for optimization techniques. The function shown in figure 11.4 exhibits the same properties. However, it emerges from the real-world problem of fitting fractal diffusion models to FRAP curves whereas the Rosenbrock function is something completely artificial.

Looking at the same function one also notices that gradient descent is one of the best methods to employ due to its robustness. Faster algorithms such as conjugate gradients would not be better as the main problem is the non-linearity of the function and not its dimensionality (CG is optimal for linear problems of high dimension). Figure 11.3 however shows a different picture. This function is very flat and regular with elongated ellipses as contour lines, such that a gradient descent will have a very poor convergence rate and need a lot of iterations. Looking at the results in appendix B in fact shows this.

11.3.4 Fitting results for all ER samples

Using all the knowledge about the models and the functional shapes, the second order physical model is fitted to all experimental FRAP curves obtained in chapter 9. The power law model will not be pursued any further as it fails to account for the influences of the bleached box geometry. All fitting runs are started from the same initial point in parameter space, namely:

$$\alpha_0 = 4.0 \qquad \qquad \beta_0 = 0.5$$

in order to be able to extract comparative information from the iteration counts. The learning rates are fixed to $\eta_{\alpha} = \eta_{\beta} = 10^{-3}$ for all runs and the tolerance for the termination criterion is set to $TOL = 10^{-6}$, meaning that both derivatives of the error function have to be less than 10^{-6} for the solution to be considered converged. The simulated FRAP curves all consist of $n_{max} = 2001$ data points for $t \in [0...20]$.

Figure 11.7 shows the results. The simulated FRAP curves are shown in dashed blue, the best model fit in red and the minimum value of the RMS fitting error as

defined by equation 11.7 is included in the captions. The numerical values for the best model parameters as well as the number of iterations it took for the solution to converge are listed in appendix B.





Figure 11.7: Second order physical model fit (red) to simulated FRAP curves (dashed blue)

The plots show that the new model almost perfectly fits all of the ER samples and even the runs erp5722, erp5723 and erp5724 with different bleached box geometries are well captured. Recall from table 8.2 in section 8.4 that the RMS error of the PSE simulations as compared to the analytic solution on the box test case is at the order of $5.7 \cdot 10^{-3}$ for the resulting FRAP curve. Since the fitting errors for the new model are at the same order of magnitude (in fact always slightly below), they are not significant. It is in principle even possible that the model perfectly describes the physics and that above RMS errors stem from the PSE simulations only. It could however also be possible that some contribution to the error comes from the PSE and another one from the model. In the worst case they are both grossly wrong but their error contributions cancel so to make the model fits look good. However, since the PSE is well validated, it seems safe to assume that its error is small and thus the model not only fits the simulation results but also the physical reality pretty well.

Chapter 12 will be concerned with comparative studies of all the old and new models described in this chapter for the box test case, for a sample ER structure and for experimental data.

11.4 The link between model parameters and the diffusion constant

Now having a new model that fits the simulated FRAP data, the question arises of how its parameters α and β reflect the diffusion constant and the dimension of the walk. Only this link would finally allow to draw information about real-world quantities from the data model fits. We are thus looking for the following functions for the second order physical model:

$$D = f_1(\alpha, \beta)$$
$$d_w = f_2(\alpha, \beta)$$

The first function is the one we are actually interested in so as to be able to use the new model for quantitative evaluations of real FRAP experiments in the laboratory. The second one is just a "side effect" but nonetheless interesting. As mentioned in section 6.3, there exists no algorithm to measure the spectral dimension or the dimension of the walk for an arbitrary given geometry. Finding f_2 however would enable us to measure both the dimension of the walk and – since we can estimate Hausdorff's dimension using the box counting algorithm of section 5.2 – the spectral dimension by simulating a FRAP experiment inside the geometry in question and fitting the new model to the result.

The following constraints for the functions f_1 and f_2 are known from various computer experiments:

- an increase in the diffusion constant leads to an increase in α and vice versa: $D \uparrow \iff \alpha \uparrow$
- an increase in the dimension of the walk leads to a decrease in β and vice versa: $d_w \uparrow \iff \beta \downarrow$
- β does not depend on D but only on the geometry d_w since two runs in the same geometry with different diffusion constants yield the same values for β:
 d_w = f₂(β)
- α depends on both D and the geometry: $D = f_1(\alpha, \beta)$
- α depends quadratically on D: $\alpha_1^2/\alpha_2^2 = D_1/D_2$

Furthermore, it turned out that f_1 is not just a simple polynomial in α and β , nor is it a polynomial in αt^{β} or any real rational function of such polynomials.

For the second order physical model to be dimensionally consistent, it has to yield something that has the same measurement unit as intensity. Looking at equation 11.4 one notices that this means that the term in parentheses has to be dimensionless. The number 1 trivially meets this requirement. The numerator of the following fraction is of dimension $(\text{length})^2$, so the denominator must be of the same dimension. This requires α to be of dimension $(\text{length}) \times (\text{time})^{-\beta}$. β has to be dimensionless as it is only found as an exponent to time. The simplest functions that fulfill these requirements and meet all constraints stated above are:

$$D = f_1(\alpha, \beta) = \frac{(\alpha T^{\beta})^2}{\tau}$$
$$d_w = f_2(\beta) = \frac{\kappa}{\beta}$$

for some constants $\tau > 0$ [time], T > 0 [time] and $\kappa > 0$ [dimensionless]. Unfortunately, this set-up fails to recover the correct diffusion constants for the simulated FRAP curves.

Due to time resource limitations it has not been possible to find the correct functions f_1 and f_2 during this work. However, enough simulations in different ER geometries have been conducted and it is possible to produce a large data set of (α, β) and corresponding D values by proper time scaling of those simulated FRAP curves. As mentioned in section 9.1 varying the diffusion constant simply corresponds to rescaling the time and no new PSE simulation is necessary. This way, a few hundred FRAP curves with different but known D values could be created for the various ER samples considered. The model can then be fitted to each of these curves to get the corresponding parameter values. This would finally yield enough data to use techniques like neural networks, data assimilation or machine learning methods to find the sought-after functional relationships.

Meanwhile, the new model can already be used to get recovery half-times according to equation 11.5. Statements about relative differences and ratios of diffusion constants in the same geometry are possible directly from the half-times as they are linearly related to each other. However, this does not work across different geometries since the geometry also has an influence on the functional relationship between the recovery half-time and the diffusion constant. This becomes obvious when calculating the half-times for all the model fits of section 11.3.4 giving $11.78^{+12.40}_{-6.08}$. However since all of them are simulation runs with the same diffusion constant, they should be equal if geometry had no influence. Existing models such as the one by [Dayel, Hom & Verkman (1999)] however completely neglect this fact using the same calibration measurement for all experiments². This still allows the diffusion constants to be compared among each other for the same cell (since only ratios are considered) but it makes them completely incomparable among different cells that have differently shaped ERs or even among different groups that use different calibration points.

 $^{^2} One$ such calibration point is that a recovery half-time of 100 ms corresponds to a diffusion constant of $5-10\,\mu m^2/s$

Chapter 12

Comparison to currently used models

The new models developed in section 11.2 will be compared to the current models as stated in section 11.1. All four models are to be compared to each other both for the box test case as defined in section 8.1 and a simulated ER sample from section 9.1. In addition, the empirical and the second order physical model are compared to each other on the experimental data of chapter 10. All model fits are made using the gradient descent algorithm as described in section 11.3.1.

12.1 Comparison on the box test case

Table 12.1 gives the initial parameters, learning rates and tolerance settings used for fitting the various models to the $n_{max} = 601$ data points of the analytical solution of the box test case. Table 12.2 summarizes the optimal model parameters found, the minimum value of the quadratic fitting error as defined by equation 11.6 and the number of iterations it took for the gradient descent algorithm to converge below the prescribed tolerance.

Model	Initial parameters	Learning rates	Tolerance
Exponential	$\alpha_0 = 0.4$	$\eta_{lpha} = 10^{-5}$	10^{-6}
Empirical	$\alpha_0 = 0.4, t_{1/2,0} = 2.0$	$\eta_{\alpha} = \eta_{t_{1/2}} = 10^{-3}$	10^{-6}
Power law	$\alpha_0 = 0.5, \beta_0 = 0.9$	$\eta_{\alpha} = \eta_{\beta} = 10^{-3}$	10^{-6}
2^{nd} order physical	$\alpha_0 = 2.0, \beta_0 = 0.5$	$\eta_{\alpha}=\eta_{\beta}=10^{-3}$	10^{-6}

Table 12.1: Parameter settings for box test case model fits

Model	Optimal parameters	E_{min}	N_{Iter}
Exponential	$\alpha_{opt} = 0.23811$	1.225543	8422
Empirical	$\alpha_{opt} = 0.94695, t_{1/2,opt} = 1.58038$	0.140867	14158
Power law	$\alpha_{opt} = 1.11405, \beta_{opt} = 0.43515$	0.085755	15115
2^{nd} order physical	$\alpha_{opt} = 0.12707, \beta_{opt} = 0.78848$	0.030658	1704

Table 12.2: Optimal parameters for box test case model fits

Figure 12.1 shows the resulting plots of the fitted models (red curves) to the analytical solution for the box test case (blue curves). The corresponding minimum value of the quadratic fitting error is listed in each figure's caption. It can be seen that even for the simple (non-fractal, homogeneous and isotropic) test case, the current models seem to be insufficient. For the exponential model this is due to the fact that it is based on 2D plate diffusion whereas the box test case is a 3D problem. The effects of dimensionality as investigated in [Sbalzarini (2001)] are thus reflected in the resulting fitting error. In fact (after proper temporal scaling), the exponential model's plot below looks very similar to figure 5.2 (page 29) in [Sbalzarini (2001)] which compares a 2D and a 3D random walk simulation. This is evidence that the influence of dimensionality largely accounts for the error in the exponential model fit. The empirical model on the other hand is completely based on measurement data for the ER and lacks physical background. It is therefore not surprising that its performance on the box test case (which is different from ER diffusion) is not convincing as it fails to extrapolate to new geometries.

Both new models perform better than the current ones since they are based on physical principles. In fact the RMS fitting error for the second order physical model is again at the order of 10^{-3} as for all the ER samples considered in section 11.3.4. This indicates that it manages to extrapolate over a wide range of geometries without grossly deteriorating in performance. It fits about one order of magnitude better than the empirical model and two orders of magnitude better than the exponential model. For the box test case, also the power law model seems quite usable. However, we will see later that it gets slightly worse for the ER samples.



Figure 12.1: Best fit of different models (red) to the analytic solution for the box test case (blue dashed)

12.2 Comparison on ER samples

Although many of the models' features can already be observed on the box test case, their performance on real simulated FRAP data from ER samples is what actually counts. Therefore, the models are also fitted to the simulation results obtained in section 9.1. Exemplarily, the results for the erp5724 run are given hereafter since this run turned out to pose the hardest problem to the models because it has a different bleached box size than all the other ones (cf. section 9.2). Table 12.3 lists the initial parameters, learning rates and tolerance settings used for fitting the various models to the $n_{max} = 2001$ data points of the simulated FRAP curve. Table 12.4 summarizes the optimal model parameters found, the minimum value of the quadratic fitting error as defined by equation 11.6 and the number of iterations it took for the gradient descent algorithm to converge below the prescribed tolerance.

Model	Initial parameters	Learning rates	Tolerance
Exponential	$\alpha_0 = 0.1$	$\eta_{lpha} = 10^{-5}$	10^{-6}
Empirical	$\alpha_0 = 0.7, t_{1/2,0} = 6.5$	$\eta_{\alpha} = \eta_{t_{1/2}} = 10^{-3}$	10^{-6}
Power law	$\alpha_0 = 0.5, \ \beta_0 = 1.0$	$\eta_{\alpha} = \eta_{\beta} = 10^{-4}$	10^{-6}
2^{nd} order physical	$\alpha_0 = 4.0, \beta_0 = 0.5$	$\eta_{\alpha} = \eta_{\beta} = 10^{-3}$	10^{-6}

Table 12.3: Parameter settings for erp5724 model fits

Model	Optimal parameters	E_{min}	N_{Iter}
Exponential	$\alpha_{opt} = 0.04342$	6.433889	24
Empirical	$\alpha_{opt} = 0.66775, t_{1/2,opt} = 16.28989$	1.36757	991209
Power law	$\alpha_{opt} = 0.18428, \beta_{opt} = 1.38605$	0.166035	344597
2^{nd} order physical	$\alpha_{opt} = 3.61430, \beta_{opt} = 0.45771$	0.012437	18318

Table 12.4: Optimal parameters for erp5724 model fits

Figure 12.2 again shows the resulting plots of the fitted models (red curves) to the simulated FRAP data (blue curves). The minimum values of the quadratic fitting error are listed along with the figures' captions. The picture is even more striking as it was for the box test case. The exponential model now seems to completely fail to capture the characteristics of the curve. Comparing this picture to the one of the box test case shows that the influence of geometry accounts for an error of 5.208 whereas the influence of dimensionality caused an error of 1.226. The geometrical structures (precise: the restriction of diffusion to them) thus have a large influence on the resulting FRAP curve even at the same value of the diffusion constant. The empirical model is also worse than it was for the box test case. One should however note, that it usually performs better on ER samples (cf. table B.2 in appendix B). This sample however features a bleached box that deviates in size and position from the ones of the other samples. The large error of the empirical model for this run indicates that it is incapable of dealing with different bleached box geometries, as already stated in section 11.1. Another interesting detail is the fact that the empirical model takes about 50 times as many iterations to converge as the second order physical one. This is in agreement with the findings of section 11.3.3.



Figure 12.2: Best fit of different models (red) to simulated erp5724 FRAP curve (blue dashed)

Again, both new models perform significantly better¹ than the empirical model, even for this worst case ER sample. The RMS error of the second order physical model fit is again below the RMS error of the PSE simulations indicating a possibly perfect match.

To allow further comparisons between the best currently used model (i.e. the empirical model) and the best new model (the second order physical model), the fitting results for all ER samples are given in appendix B. Apart from the exception erp573.2, the new model constantly fits one to two orders of magnitude better than the empirical model.

Figure 12.3 shows the model matching errors for the empirical model (blue dashed curve) and the second order physical model (red curve) as a function of the iteration step number k of the gradient descent algorithm 11.1. Both axes use logarithmic scaling to get a clear plot.

 $^{^1 \}mathrm{one}$ order of magnitude for the power law model, two orders of magnitude for the second order physical model



Figure 12.3: Model matching error for sample erp5724 and empirical model (blue dashed) as well as second order physical model (red). Notice the double logarithmic scaling.

12.3 Comparison on experimental data

Finally, the two best models, i.e. the empirical one and the new second order physical model, are compared with real experimental data (courtesy of Anna Mezzacasa). Therefore, the same two experiments as presented in chapter 10 are used. Example 8s has 15 data points and 8.2 has 23 data points. Both experiments have been made with a bleached spot of size $5 \,\mu m \times 5 \,\mu m$. The initial parameters and optimization settings are listed in table 12.5. Table 12.6 gives the optimal model parameters, the minimum value of the quadratic fitting error and the number of iterations needed for the 8s sample and table 12.7 for the 8.2 sample.

Model	Initial parameters	Learning rates	Tolerance
Empirical	$\alpha_0 = 1.0, t_{1/2,0} = 0.05$	$\eta_{\alpha} = 10^{-4}, \ \eta_{t_{1/2}} = 5 \cdot 10^{-5}$	10^{-6}
2^{nd} order physical	$\alpha_0 = 5.0, \beta_0 = 0.8$	$\eta_{\alpha} = 0.1, \ \eta_{\beta} = 5 \cdot 10^{-4}$	10^{-6}

Table 12.5: Parameter settings for experimental data model fits

Model	Optimal parameters	E_{min}	N_{Iter}
Empirical	$\alpha_{opt} = 11.19487, t_{1/2,opt} = 0.01775$	$4.49922 \cdot 10^{-3}$	1118631
2^{nd} order physical	$\alpha_{opt} = 25.93876, \ \beta_{opt} = 0.77531$	$3.90525 \cdot 10^{-3}$	2794396

Table 12.6: Optimal parameters for 8s model fits

Figure 12.4 shows the corresponding plots. The measured data points are shown as blue crosses, the optimally fitted model is drawn as a red line. Again, the second order physical model performs better than the empirical one. However, the

Model	Optimal parameters	E_{min}	N_{Iter}
Empirical	$\alpha_{opt} = 11.02793, t_{1/2,opt} = 0.02912$	$1.34563 \cdot 10^{-3}$	523349
2^{nd} order physical	$\dot{\alpha_{opt}} = 17.49833, \beta_{opt} = 0.78028$	$7.44625 \cdot 10^{-4}$	677758

Table 12.7: Optimal parameters for 8.2 model fits

difference is smaller than for the simulated data sets. This was to be expected as the empirical model is based on such experimental data and has been trained to fit it (including its measurement errors and noise).



Figure 12.4: Best fit of different models (red) to experimental FRAP data (blue crosses)

Again, the model matching error is plotted against the iteration step number of the gradient descent algorithm using a double logarithmic scaling. The results for both experimental data sets are shown in figure 12.5 where the blue dashed curve stands for the empirical model and the red curve for the second order physical one.

According to equation 11.5, the recovery half-times for the second order physical model can also be calculated. Table 12.8 lists the results and the relative deviation from the empirical model. The models are in agreement insofar as their mutual difference is less than the statistical scatter (as both experiments have been made using the same ER, the diffusion constant and therefore the half-times for both samples should be equal).

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Figure 12.5: Model matching error for empirical model (blue dashed) as well as second order physical model (red) for both experiments. Notice the double logarithmic scaling.

Sample	$t_{1/2}$ empirical	$t_{1/2}$ physical	relative deviation
8s	$17.75\mathrm{ms}$	$15.70\mathrm{ms}$	11.5%
8.2	$29.12\mathrm{ms}$	$26.69\mathrm{ms}$	8.3%
average	$23.4\pm5.7\mathrm{ms}$	$21.2\pm5.5\mathrm{ms}$	9.4%

Table 12.8: Recovery half-times and their deviation

As stated in section 11.4, this would now allow one to deduce the diffusion constant if the influence of geometry was known. Wrongly applying the calibration point of $[Dayel, Hom \& Verkman (1999)]^2$ to the geometry of our samples yields $D = 21.8 - 43.5 \,\mu\text{m}^2/\text{s}$ for the empirical model and $D = 23.6 - 47.2 \,\mu\text{m}^2/\text{s}$ for the second order physical model. Comparing these numbers to the result $D = 54 \,\mu \text{m}^2/\text{s}$ obtained from direct numerical simulation in chapter 10 shows that they differ by a factor of 1.2 to 2.5 for the empirical model and 1.1 to 2.3 for the second order physical model. This difference is caused by neglecting geometrical influences. The fact that the deviation for the second order physical model is slightly smaller is due to its better fitting behavior. Using the empirical model, various papers (e.g. [Dayel, Hom & Verkman (1999)]) report diffusion coefficients between $25\,\mu\text{m}^2/\text{s}$ and $30\,\mu\text{m}^2/\text{s}$ for GFP in the cytoplasm, between $5\,\mu\text{m}^2/\text{s}$ and $10\,\mu\text{m}^2/\text{s}$ in the ER lumen and between $20 \,\mu \text{m}^2/\text{s}$ and $30 \,\mu \text{m}^2/\text{s}$ in the mitochondrial matrix (see [Lippincott-Schwartz, Snapp & Kenworthy (2001)] for a summary). The difference between these numbers and the one obtained with the old empirical model for the present experiments is thus caused by differences in the cells, proteins or experimental set-ups used and has nothing to do with the models involved.

 $^{^2 {\}rm Stating}$ that $t_{1/2} = 100\,{\rm ms}$ corresponds to a diffusion constant of $5-10\,\mu{\rm m}^2/{\rm s}$

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Chapter 13

Conclusions and future work

This project successfully extended the work of [Sbalzarini (2001)] to study not only the influence of dimensionality but also the one of geometry on diffusion processes in the endoplasmic reticulum. The use of fractal concepts and latest computer simulation techniques made it possible to move on to three-dimensional realistic ER shapes. Changing the simulation technique from random walk to the method of particle strength exchange allowed accurate treatment of diffusion in arbitrary domains. In order to do so, 3D reconstructions of real ER samples were made as described in chapter 2. The triangulation defining their surface was then checked for various validity criteria in chapter 3.

As stated in the future work list of [Sbalzarini (2001)], a sound mathematical treatment of diffusion on fractal sets was needed in order to proceed with using concepts from fractal geometry to find novel FRAP data models for diffusion in the endoplasmic reticulum. Towards this end, chapter 4 introduced and summarized some of the theoretical foundations about measures and dimensions for later use. In this context, different fractal dimensions such as the box counting dimension or Hausdorff's dimension were properly introduced and defined, extending section 10.1 of [Sbalzarini (2001)]. Chapter 6 then made use of this theoretical background to investigate diffusion on fractal sets. Starting from Brownian motion on the Sierpinski gasket, a proof has been presented that Hausdorff's dimension is not a sufficient parameter to capture the time behavior of diffusion on fractals as was suspected in [Sbalzarini (2001)]. Considering transition densities and eigenvalues of the Laplacian, the *dimension of the walk* has been found to actually be the sought-after sufficient geometry parameter. This finding has also been extended to infinitely ramified fractals and anisotropic diffusion.

Chapter 5 was concerned with measuring the fractal dimension of the ER surface as requested in [Sbalzarini (2001)]. Therefore, an algorithm to transform the triangulated representation to a voxel set was presented and validated. A 3D generalization of the box counting algorithm has then been developed and used to determine the box counting dimension of the ER surface to be 2.8. This value is within the theoretically predicted range of [Sbalzarini (2001)].

Chapter 7 introduced the numerical simulation techniques used in this work. An existing code for the anisotropic particle strength exchange method was successfully extended to handle bounded computational domains of arbitrary shapes with either Neumann or Dirichlet boundary conditions. To this end, a fast geometry preprocessor has been developed as well as proper post-processing and visualization techniques. A new set of algorithms is now available that is much more sophisticated and trustworthy than the 2D finite difference tools applied

by [Siggia, Lippincott-Schwartz & Bekiranov (2000)], [Ellenberg et al. (1997)] and [Dayel, Hom & Verkman (1999)].

To validate the simulation algorithms, a test case has been introduced in chapter 8 and its analytic solution derived and analyzed. A simple finite difference code served as a cross-validation. Random walk and PSE simulations have been shown to yield results close to the exact solution, justifying their use for productive simulations. Comparative error considerations and timings have been made to analyze the codes' performance.

Chapter 9 presented diffusion simulations in a number of ER samples in order to obtain simulated FRAP curves. Moreover, the influence of the position and size of the bleached box was investigated and shown to be relevant. The results obtained for the ER geometry differ significantly from the one of the cubic box considered in [Sbalzarini (2001)]. Comparative timings and parallel speed-up estimations concluded this chapter.

In chapter 10, a simulation has been presented that was run in the exact same geometry as used for a real laboratory experiment. Very good coincidence was observed for sufficiently resolved simulations and still good coincidence for underresolved runs. Moreover, it has been shown that it is possible to determine diffusion constants by fitting simulation data to experimental measurements. The diffusion constant obtained this way differs from the one predicted by the current empirical model of [Dayel, Hom & Verkman (1999)] by a factor of 1.2 to 2.5. Neglect of the geometrical influences hereby accounts for 1.1 to 2.3, the rest is caused by using a model that does not optimally fit the data.

In chapter 11, new FRAP data models have been proposed after briefly reviewing some existing ones. The new models are based on physical principles. Using the theory of diffusion on fractal sets as outlined in chapter 6, the new models also take into account the local ER geometry which has been shown in chapter 9 to significantly influence the FRAP results. One model also includes basic information about the bleached box geometry which has been shown in section 9.2 to affect the FRAP results. After noting some properties of the model matching error functions, the new model's parameters have been determined for all simulated FRAP curves using a gradient descent method. For the new model, almost perfect correspondence was observed for all ER samples with an RMS error which is below the one of the PSE simulation itself. The chapter was concluded by some considerations on how to interpret the models' parameters in terms of real-world physical properties.

The new models have been compared to the current ones in chapter 12 on the box test case, on ER samples and on experimental data. All comparisons show the same picture: the new physical model is better than the current models while fully accounting for geometry and bleached box influences. Comparing the fitting behavior of the 2D exponential model for the box test case and the ER sample allowed to make the observation that the influences of geometrical structures are some 4 to 5 times more important to the time behavior of diffusion (and thus the FRAP curves) than the influences of dimensionality as investigated in [Sbalzarini (2001)].

In the limited time of this diploma thesis, it was of course not possible to address all the questions in desirable detail. Namely the following points will have to be subjects to future work:

• A functional relationship between the parameters of the new model and the diffusion constant has to be found in order to complete the considerations of section 11.4. Due to the required extrapolation capabilities, it is desirable to find a physical relationship rather than using neural networks or machine learning methods.

- Once such a function is known, a small and stable program for automatic quantitative evaluation of experimental FRAP data has to be developed. Such an application would allow robust and coherent determination of numerical diffusion coefficients in daily laboratory use, making results comparable among different groups.
- More accurate (i.e. higher resolution) simulations should be run in order to compare them to experiments and investigate whether the deviations observed in one case in chapter 10 stem from simulation or measurement errors.
- More accurate initial conditions using an initial fluorophore concentration profile as proposed in [Axelrod et al. (1976)] should be included in the simulations and their influence on the resulting FRAP curve should be investigated.
- More and better FRAP data models have to be found and evaluated.
- Finally, diffusion of membrane-bound proteins should be addressed. An extension of the method of particle strength exchange to surface-bound diffusion will however first have to be theoretically derived. This consists of finding a discretization of the Beltrami operator¹ on particles as well as a set of (local) mappings to theoretically describe the ER's surface. The work of [Spekreijse, Hagmeijer & Boerstoel (1996)] could serve as a starting point.

Ultimately, it is desirable to go beyond simulating existing experiments in order to start investigating things not amenable to them and create novel knowledge by using computer simulations only. Examples of such applications include:

- Investigations on whether the medium filling the ER lumen is anisotropic (what ratio of anisotropy ?) or not. This can be done by conducting anisotropic simulations (recall that all simulation codes already include the functionality needed) and fitting their results to experimental data. If those results fit better than the isotropic ones, this could be evidence that diffusion in reality is anisotropic.
- Investigation on whether the ratio and spatial structure of anisotropy do have an influence on the resulting FRAP curves at all (cf. chapter 10). This consists of running different anisotropic simulations and comparing their FRAP curves.
- Use of the new data model to analyze experimental FRAP curves in order to obtain diffusion coefficients that allow biological conclusions about protein folding and sorting.
- Applications of the fractal diffusion theory to completely different fields such as transport phenomena in porous media, complex electric systems or heat conduction.

Altogether, a large step in understanding diffusion processes in the endoplasmic reticulum as well as geometrical influences on diffusion in general has been made when comparing to the state in [Sbalzarini (2001)] and all assigned goals of the project have been achieved (cf. foreword). Nevertheless, some open points remain and there is still work to do until the new models are ready for productive, unattended laboratory use. A sound basis however exists by now and the hope remains that the final goal is not too far any more.

¹The Beltrami operator describes diffusion which is restricted to a curved surface in space just as the Laplace operator describes free diffusion in space.

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Bibliography

- [Alberts et al. (1997)] ALBERTS, B., BRAY, D., JOHNSON, A., LEWIS, J., RAFF, M., ROBERTS, K. & WALTER, P., 1997, Essential Cell Biology, Garland publishing, Inc., New York.
- [Axelrod et al. (1976)] AXELROD, D, KOPPEL, D. E., SCHLESSINGER, J., ELSON, E. & WEBB, W. W., 1976, Mobility Measurement by Analysis of Fluorescence Photobleaching Recovery Kinetics, *Biophysical Journal*, 16, 1055-1069.
- [Baldock & Graham (2000)] BALDOCK, R. & GRAHAM, J. (EDS.), 2000, Image Processing and Analysis - A Practical Approach, Oxford University Press, New-York.
- [Barlow & Bass (1992)] BARLOW, M. T. & BASS, R. F., 1992, Transition densities for Brownian motion on the Sierpinski carpet, *Probability Theory and related fields*, **91**, Springer, 307-330.
- [Barlow, Hattori, Hattori & Watanabe (1997)] BARLOW, M. T., HATTORI, K., HATTORI, T. & WATANABE, H., 1997, Weak Homogenization of Anisotropic Diffusion on Pre-Sierpinski Carpets, *Communications in Mathematical Physics*, 188, Springer, 1-27.
- [Barlow & Perkins (1988)] BARLOW, M. T. & PERKINS, E. A., 1988, Brownian Motion on the Sierpinski Gasket, Probability Theory and related fields, 79, Springer, 543-623.
- [Beaudoin, Huberson & Rivoalen (2001)] BEAUDOIN, A., HUBERSON, S. & RIVOALEN, E., Simulation of anisotropic diffusion by means of a diffusion velocity method, *submitted to J. Comp. Phys.*.
- [Cheezum, Walker & Guilford (2001)], CHEEZUM, M. K., WALKER, W. F. & GUILFORD, W. H., 2001, Quantitative Comparison of Algorithms for Tracking Single Fluorescent Particles, *Biophysical Journal*, 81, 2378-2388.
- [Cottet & Koumoutsakos (2000)] COTTET, G.-H. & KOUMOUTSAKOS, P. D., 2000, Vortex Methods Theory and Practice, Cambridge University Press.
- [Dayel, Hom & Verkman (1999)] DAYEL, M. J., HOM, E. F. Y. & VERKMAN, A. S., 1999, Diffusion of Green Fluorescent Protein in the Aqueous-Phase Lumen of the Endoplasmic Reticulum, *Biophysical Journal*, **76**, 2843-2851.
- [Degond & Mas-Gallic (1989a)] DEGOND, P. & MAS-GALLIC, S., 1989, The weighted particle method for convection-diffusion equations, Part I: The case of isotropic viscosity, *Math. Comp.*, 53, 485.
- [Degond & Mas-Gallic (1989b)] DEGOND, P. & MAS-GALLIC, S., 1989, The weighted particle method for convection-diffusion equations, Part 2: The anisotropic case, *Math. Comp.*, **53**, 509.

- [Dietrich et al. (2002)] DIETRICH, C., YANG, B., FUJIWARA, T., KUSUMI, A. & JACOBSON, K., 2002, Relationship of Lipid Rafts to Transient Confinement Zones Detected by Single Particle Tracking, *Biophysical Journal*, 82, 274-284.
- [Ellenberg et al. (1997)] ELLENBERG, J, SIGGIA, E. D., MOREIRA, J. E., SMITH, C. L., PRESLEY, J. F., WORMAN, H. J. & LIPPINCOTT-SCHWARTZ, J., 1997, Nuclear Membrane Dynamics and Reassembly in Living Cells: Targeting of an Inner Nuclear Membrane Protein in Interphase and Mitosis, *Journal of Cell Biology*, Vol. 138, Nr. 6, 1193-1206.
- [Falconer (1985)] FALCONER, K. J., 1990, The Geometry of Fractal Sets, Cambridge University Press.
- [Falconer (1990)] FALCONER, K. J., 1990, Fractal Geometry Mathematical Foundations and Applications, John Wiley & Sons, Chichester, UK.
- [Falconer (1997)] FALCONER, K. J., 1997, Techniques in Fractal Geometry, John Wiley & Sons, Chichester, UK.
- [Hattori, Hattori & Watanabe (1994)] HATTORI, K., HATTORI, T. & WATANABE, H., 1994, Asymptotically one-dimensional diffusions on the Sierpinski gasket and the abc-gaskets, *Probability Theory and related fields*, **100**, Springer, 85-116.
- [Hausdorff (1919)] HAUSDORFF, F., 1919, Dimension und äusseres Mass, Mathematische Annalen, 79, 157-179.
- [Hockney & Eastwood (1988)] HOCKNEY, R. W. & EASTWOOD, J. W., 1988, Computer simulation using particles, Institute of Physics Publishing, Bristol/Philadelphia.
- [Hou, Gilmore, Mindlin & Solari (1990)] HOU, X.-J., GILMORE, R., MINDLIN, G.
 B. & SOLARI, H. G., 1990, An efficient algorithm for fast O(N*log N) box counting, *Physics Letters A*, Vol. 151, No. 1 2, 43-46.
- [Kamm (2002)] KAMM, R. D., 2002, Cellular Fluid Mechanics, Annual Revisions in Fluid Mechanics, 34, 211-232.
- [Kusumi, Sako & Yamamoto (1993)] KUSUMI, A., SAKO, Y. & YAMAMOTO, M., 1993, Confined Lateral Diffusion of Membrane Receptors as Studied by Single Particle Tracking (Nanovid Microscopy). Effects of Calcium-Induced Differentiation in Cultured Epithelial Cells, *Biophysical Journal*, 65, 2021-2040.
- [Liebovitch & Toth (1989)] LIEBOVITCH, L. S. & TOTH, T., 1989, A fast algorithm to determine fractal dimensions by box-counting, *Physics Letters A*, Vol. 141, No. 8 9, 386-390.
- [Lippincott-Schwartz, Snapp & Kenworthy (2001)] LIPPINCOTT-SCHWARTZ, J., SNAPP, E. & KENWORTHY, A., 2001, Studying Protein Dynamics in living Cells, Nature Rev. Cell Biology, 2, 444.
- [Mandelbrot (1982)] MANDELBROT, B. B., 1982, The fractal geometry of nature, W. H. Freeman and Co., San Francisco.
- [McCorquodale, Colella & Johansen (2001)] MCCORQUODALE, P., COLELLA, P. & JOHANSEN, H., 2001, A Cartesian Grid Embedded Boundary Method for the Heat Equation on Irregular Domains, *Journal of Computational Physics*, 173, 620-635.
- [Mitra, Murthy, Kundu & Bhattacharya (2001)] MITRA, S. K., MURTHY, C. A., KUNDU, M. K. & BHATTACHARYA, B. B., 2001, Fractal Image Compression Using Iterated Function System with Probabilities, *Proc. Intl. Conference on Information Technology (ITCC'01)*, IEEE Computer Society, 191-195.
- [Monks (2001)] MONKS, K., 2001, Definitions and Theorems for Chaos and Fractals, *lecture monograph*, Course Math-320, University of Scranton.
- [Numerical Recipes in Fortran 90 (1996)] PRESS, W. H., TEUKOLSKY, S. A., VETTERLING, W. T. & FLANNERY, B. P., 1996, Numerical Recipes in Fortran 90 - The Art of Parallel Scientific Computing, 2nd edition, Cambridge University Press.
- [Olveczky & Verkman (1998)] OLVECZKY, B. P. & VERKMAN, A. S., 1998, Monte Carlo Analysis of Obstructed Diffusion in Three Dimensions: Application to Molecular Diffusion in Organelles, *Biophysical Journal*, 74, 2722-2730.
- [Reits & Neefjes (2001)] REITS, E. A. J. & NEEFJES, J. J., 2001, From fixed to FRAP: measuring protein mobility and activity in living cells, *Nature Cell Biology*, Vol. 3, June 2001, E145-E147.
- [Sbalzarini (2001)] SBALZARINI, I. F., 2001, On Protein Diffusion in the Endoplasmic Reticulum - a computational approach using particle simulations, *ICoS/ETHZ semester thesis*, July 2001.
- [SGI (1992a)] SILICON GRAPHICS, INC., 1992, How to write an IRIS Inventor File Translator, Release 1.0.
- [SGI (1992b)] SILICON GRAPHICS, INC., 1992, IRIS Inventor Nodes Quick Reference, Release 1.0.
- [Shannon (1948)] SHANNON, C. E., 1948, A Mathematical Theory of Communication, The Bell System Technical Journal, 27, 379-656.
- [Siggia, Lippincott-Schwartz & Bekiranov (2000)] SIGGIA, E. D., LIPPINCOTT-SCHWARTZ, J. & BEKIRANOV, S., 2000, Diffusion in Inhomogeneous Media: Theory and Simulations Applied to Whole Cell Photobleach Recovery, *Bio-physical Journal*, **79**, 1761-1770.
- [Smith, Marks, Lang, Sheriff & Neal (1989)] SMITH, T. G. JR., MARKS, W. B., LANG, G. D., SHERIFF, W. H. JR. & NEAL, E. A., 1989, A Fractal Analysis of Cell Images, *Journal of Neuroscience Methods*, 27, 173-180.
- [Spekreijse, Hagmeijer & Boerstoel (1996)] SPEKREIJSE, S. P., HAGMEIJER, R. & BOERSTOEL, J. W., 1996, Adaptive grid generation by usgin the Laplace-Beltrami operator on a monitor surface, *National Aerospace Laboratory Technical Report*, TP 96297, Amsterdam, The Netherlands.
- [Stoll, Stern & Stucki (1996)] STOLL, E. P., STERN, C. & STUCKI, P., 1996, Fractals in isotropic systems generated with attracting spheres, *Physica A*, 230, 11-18.
- [Zimmermann, Koumoutsakos & Kinzelbach (2001)] ZIMMERMANN, S., KOUMOU-TSAKOS, P. & KINZELBACH, W., 2001, Simulation of Pollutant Transport Using a Particle Method, Journal of Computational Physics, 173, 322-347.

BIBLIOGRAPHY

Appendix A

Contents of the companion CD-ROMs

All data files that have been produced in digital form during this diploma project are contained on the two attached CD-ROMs.

The CDs contain this report (source texts and final PS and PDF versions as well as all pictures), all Fortran source code files, shell scripts and perl scripts, some movies and the final PowerPoint presentation of the work. Moreover, all original micrographs are included as well as the original input and output files of all the runs presented in this report and all experimental data. Finally, the CDs also contain some 3^{rd} party papers that are related to this work.

All directories on both CDs contain **README** files with further information on the files and their usage.

CD #2 only contains a single directory (runs) which holds compressed archive files (.tar.gz) of all the simulation runs. They can be uncompressed using the command

\$ tar zxvf <run>.tar.gz

This will create a directory called **<run>** that contains all the input and output files of the simulation run according to the information in the **README** file.

The file system of CD #1 is structured as follows:

DA02_Sbalzarini_1/	
--------------------	--

$3 \mathrm{div}/$	3D reconstructions of all ER samples
$\operatorname{codes}/$	Source code files of all programs
analytic/	Calculates the analytic solution (section 8.2)
bdiff3d/	Random walk simulation code (section 7.1)
$\operatorname{run}/$	Input files and utilities
$\operatorname{src}/$	Source code
$\operatorname{concat}/$	Concatenates PSE3D output files
crayfishpak/	Fast finite difference solver libraries
$\mathrm{FD}/$	Finite difference solver (section 8.3)
$\mathrm{nlfit}/$	Gradient descent algorithm (section $11.3.1$)
$\operatorname{runs}/$	Results of all model fits made
m src/	Source code
preproc/	Geometry preprocessor (section $7.2.5$)
$\operatorname{run}/$	Input files and visualization utilities
$\operatorname{src}/$	Source code
PSE3D/	PSE simulation code (section 7.2)
$\operatorname{run}/$	Input files
$\operatorname{src}/$	Source code
experiments/	Experimental data and evaluation tools used in chapter 10
m micrographs/	Stained ER micrograph stacks
8s/	Stack of 8s and 8.2 experiments (chapter 10)
movies/	Movies of experiments and exact solution
utilities/	Scripts and helper programs to make movies
papers/	Some relevant 3^{rd} party papers
postproc/	OpenDX converter and visualization tools presented in section 7.3
presentation/	Final project presentation (PowerPoint)
$\operatorname{report}/$	This report
source/	${\rm I\!AT}_{\!E\!} \! X$ source files of this report
$\operatorname{code}/$	Source code files of appendix C
pics/	All EPS and PS figures of this report
erp572/	Sample micrograph stack erp572 (chapter 2) $$
masters/	Original xfig drawings and pixel images

Appendix B

Fitted model parameters

B.1 Second order physical model fit

The following table lists the optimal model parameters α_{opt} and β_{opt} for the second order physical model given by equation 11.4 in section 11.2.2 found for the different simulation runs of chapter 9. In order not to loose any information, they are given to one more significant digit than the accuracy they were computed with, well aware of the fact that they cannot be trusted for more than 3 digits. The column E_{min} lists the minimum value of the quadratic (not RMS !) fitting error E as defined by equation 11.6 and N_{Iter} is the number of iterations it took for the gradient descent algorithm to converge up to a tolerance of 10^{-6} . Since all runs started from the same initial point in parameter space, this number is a measure for the distance between the initial and the optimal point. Details about the runs as well as comparative plots are contained in section 11.3.4.

Sample	α_{opt}	eta_{opt}	E_{min}	N_{Iter}
bip2	1.9440048	0.63324605	$3.629066 \cdot 10^{-2}$	7924
clx	3.3294665	0.58295736	$1.825963 \cdot 10^{-2}$	11773
erp57	2.4120744	0.63037812	$6.61639 \cdot 10^{-3}$	8881
erp572	2.8619264	0.48297255	$1.12797 \cdot 10^{-2}$	10586
$erp573_1$	3.2154385	0.62770381	$3.07065 \cdot 10^{-2}$	11444
$erp573_2$	3.4254690	0.46527276	$1.66678 \cdot 10^{-2}$	11758
erp573_3	2.9901725	0.63353217	$2.574665 \cdot 10^{-2}$	11356
$erp574_1$	2.7352222	0.51042596	$2.789649 \cdot 10^{-2}$	10467
$erp574_2$	3.8617008	0.56668043	$3.05717\cdot 10^{-3}$	12512
erp5722	2.4727129	0.57562438	$2.120234 \cdot 10^{-2}$	9303
erp5723	1.6467670	0.66137837	$5.019553 \cdot 10^{-2}$	6922
erp5724	3.6143039	0.45770528	$1.243710 \cdot 10^{-2}$	18318

Table B.1: Fitting results for the second order physical model

B.2 Empirical model fit

Table B.2 lists the optimal model parameters α_{opt} and $t_{1/2,opt}$ for the empirical model given by equation 11.2 in section 11.1 found for the different simulation runs of chapter 9. In order not to loose any information, they are given to one more significant digit than the accuracy they were computed with, well aware of the fact that they cannot be trusted for more than 3 digits. The column E_{min} lists the minimum value of the quadratic (not RMS !) fitting error E as defined by equation 11.6 and N_{Iter} is the number of iterations it took for the gradient descent algorithm to converge up to a tolerance of 10^{-6} . Since all runs started from the same initial point in parameter space, this number is a measure for the distance between the initial and the optimal point.

Sample	$lpha_{opt}$	$t_{1/2,opt}$	E_{min}	N_{Iter}
bip2	0.74410405	12.882408	0.22977	401454
$_{\rm clx}$	0.66395129	6.6161354	0.0852237	43693
erp57	0.71880026	9.6646334	0.0514976	142477
erp572	0.56902663	13.142333	0.127082	735753
$erp573_1$	0.71491157	6.0489575	0.12589	38403
erp573_2	0.51711024	11.242306	0.0376993	695268
erp573_3	0.72134601	6.7077755	0.106629	46024
$erp574_1$	0.59855700	12.441742	0.163978	524343
$erp574_2$	0.65499708	5.4388216	0.0249841	37575
erp5722	0.66276097	11.301493	0.121126	285086
erp5723	0.80072016	14.398334	0.371151	546584
erp5724	0.66775435	16.289888	1.36757	991209

Table B.2: Fitting results for the empirical model

Appendix C The simulation codes

For future reference and the purpose of complete transparency and reproducibility of the algorithms developed and used in this work, the source code listings of some programs are given in this appendix. Each section contains information about a different program starting with a general description of the code and its usage, syntax and explanation of all input and output files as well as an overview of the code's structure and calling tree (if meaningful). In all syntax examples, mandatory values are enclosed by triangular brackets <.>, optional ones by square brackets [.]. The dollar sign \$ symbolizes a command line prompt.

These explanations are then followed by a reprint of the source code listing. Each subroutine is presented in a different subsection preceded by a descriptive header stating its purpose and calling syntax. All programs have been written in the Fortran 90 programming language. For easier reference and readability, statement numbers have been included in the listings to which the algorithms described in this report contain references. They are for documentation purposes only and not present in the original "compilable" files. The files as printed hereafter are as of February 21, 2002. Due to the ongoing nature of the project, they can and will be altered and further developed in the future.

Please note that not all programs utilized in this project are given in this appendix. The following codes are omitted for various reasons (but are contained on CD #1 nevertheless):

- The new version of the random walk code because it only contains minor changes compared to the version that has been printed in [Sbalzarini (2001)] and all the changes are fully documented in section 7.1.
- The PSE simulation code because it is too large and has mainly been developed by other people in other projects. Moreover, none of its algorithms are presented in this report in detail making any references to it obsolete.
- The code that calculates FRAP value predictions based on current models because it is too trivial.
- All the small MATLAB files for visualization and least squares fitting because they are straightforward.
- The tools used to concatenate and average data files because they are implemented in an obvious manner.

C.1 Calculating the analytic solution

C.1.1 General description and usage notes

This section contains the program code used to numerically calculate the analytic solution for the box test case as derived in section 8.2. The code only consists of a single source file and has the following input parameters that are directly hard-coded into its header (statements 4 through 15):

Parameter	Meaning
Len (L)	Edge length of the cubic computational domain
a(a)	Lower x-coordinate of the bleached box
b (b)	Upper x-coordinate of the bleached box
c (<i>c</i>)	Lower y-coordinate of the bleached box
d(d)	Upper y-coordinate of the bleached box
dt (δt)	Time step size
tfinal (t_{final})	Final time of solution
tO (t_0)	Starting time
$\mathtt{Diff}(D)$	Diffusion constant
cO (c_0)	Initial concentration outside the bleached box
N(N)	Number of grid points in each direction
M(M)	Number of terms in the series expansion

Table C.1: Input parameters for analytical.f90

The output consists of a set of files called uXXX.out where XXX stands for the corresponding time step number. These files contain the analytical solution at the time steps indicated. Each line contains the solution at one grid point using the following syntax:

<x-pos> <y-pos> <z-pos> <solution u>

They are used to visualize the exact solution at certain time steps such as in figure 8.3. Two other output files are called frap.out and mass.out and they contain the FRAP values and the total mass in the system versus time using the syntax:

<time> <value>

on every line. They are used to create various FRAP curve plots as well as to check the accuracy of the simulation using the principle of conservation of mass. Examples are given in figures 8.4 and 8.5, respectively. Additionally, the program prints a message like

Completed time step 20 of 601

to the system's standard output every time it completes a time step. This is for progress monitoring purposes only.

C.1.2 Source code listing

The following source code lists the complete program used to calculate the analytic solution given in section 8.2 as an infinite sum over eigenfunctions. It already contains the high performance Fortran directive needed to parallelize it on a vector processor.

C.1.2.1 Analytic solution (analyt.f90)

```
1 program analytical
```

```
2 IMPLICIT NONE
```

!analytical Calculates the analytic solution for the box test case. ANALYTICAL calculates the analytic solution for diffusion in a cube of side length L with zero flux boundary conditions. The initial condition is 1 everywhere except in the cylinder [a,b]x[c,d]x[0,L]where it is zero initially. The solution is for homogenious, isotropic diffusion. In addition, the resulting FRAP curve is calculated and written to the file frap.out. THIS IS OPTIMIZED FOR PGHPF HIGH PERFORMANCE FORTRAN (PARALLELISM) ETH-ZUERICH ! DIPLOMA THESIS WS01/02 ICOS _____ PROTEIN DIFFUSION INSIDE THE ENDOPLASMIC RETICULUM I _____ ! Declaration of parameters _____ ! precision (4=single, 8=double) 3 INTEGER, PARAMETER :: MK = KIND(1.0D0) ! edge length of the cube 4 REAL(MK), PARAMETER :: Len = 4.0! lower x coordinate of bleached box 5 REAL(MK), PARAMETER :: a = 2.0 ! upper x coordinate of bleached box REAL(MK), PARAMETER :: b = 3.0 6 ! lower y coordinate of bleached box REAL(MK), PARAMETER 7 :: c = 2.0 ! upper y coordinate of bleached box REAL(MK), PARAMETER :: d = 3.0 8 ! time step size 9 REAL(MK), PARAMETER :: dt = 0.05! final time to calculate solution up to REAL(MK), PARAMETER 10 :: tfinal = 30.0 ! initial time to start from REAL(MK), PARAMETER :: t0 = 0.011 ! Diffusion coefficient REAL(MK), PARAMETER 12 :: Diff = 0.03! initial concentration outside bleached box REAL(MK), PARAMETER :: c0 = 1.013 ! number of grid points in each direction to get solution at INTEGER, PARAMETER :: N = 41 14 ! number of terms to calculate in the solutions series expansion 15 INTEGER, PARAMETER :: M = 3001----------! Declaration of local variables !---! loop counters 16 INTEGER :: i, j ,k, l, istep ! the four integral parts of the solution 17 REAL(MK) :: I1, I2, I3, I4 ! intermediate terms 18 REAL(MK) :: t1, t2, t3, h, im, xm ! often used sqrt terms 19 REAL(MK) :: w2u1, w21, w1u1 ! grid spacing 20 REAL(MK) :: dx ! current time 21 REAL(MK) :: t ! current position 22 REAL(MK) :: x, y

```
! solution
     REAL(MK), DIMENSION(:,:), ALLOCATABLE
23
                                                :: u
      ! coordinates of points at which solution is avaluated
24
     REAL(MK), DIMENSION(:), ALLOCATABLE
                                                :: px
      ! error trap variable
25
     INTEGER
                                                :: istat
      ! values of x and y eigenfunctions
26
     REAL(MK)
                                                :: ukx, uly
      ! eigenvalue
27
     REAL(MK)
                                                :: 11
      ! FRAP data of all time steps
     REAL(MK), DIMENSION(:), ALLOCATABLE
28
                                                :: frap
      ! total mass of all times (to check conservation)
     REAL(MK), DIMENSION(:), ALLOCATABLE
29
                                                :: mass
      ! name of current output file
30
     CHARACTER(LEN=80)
                                                :: filename
      ! minimum and maximum time step
     INTEGER
31
                                                :: Tmin, Tmax
      ! index boundaries of bleached box
32
     INTEGER
                                                :: iblx, ibux, ibly, ibuy
      ! total number of grid points inside bleached box
33
     INTEGER
                                                :: nb
      ! the inverse length of the domain
34
     REAL(MK)
                                                :: Leninv
      ! guess what ...
35
     REAL(MK)
                                                :: PI
      ! Set up solution grid and allocate memory
      1----
36
     PI = 4.0_MK*datan(1.0_MK)
     dx = Len/(N-1.0)
37
                                     ! grid spacing
     Tmax = ceiling(tfinal/dt)
38
                                   ! number of time steps
     Tmin = ceiling(t0/dt)
39
     iblx = ceiling(a/dx)+1
40
                                     ! boundary indices of bleached box
     ibux = floor(b/dx)+1
41
42
     ibly = ceiling(c/dx)+1
     ibuy = floor(d/dx)+1
43
     nb = (ibux-iblx+1)*(ibuy-ibly+1) ! number of pts in bleached box
44
45
     istat = 0
     ALLOCATE(u(N,N), px(N), frap(Tmax-Tmin+1), mass(Tmax-Tmin+1))
46
47
     if (istat .NE. 0) then
48
        WRITE(*,'(2A)') 'Unable to allocate memory for solution.'
        goto 9999
49
50
     end if
     px(1:N) = real((/(i,i=0,N-1,1)/))*dx ! coordinates of points
51
      1-----
      ! Calculate solution for all time steps and grid points
      1 - - -
      ! some expressions that occur often
     w2ul = dsqrt(2.0_MK/Len)
52
53
     w21 = dsqrt(2.0_MK*Len)
54
     w1ul = dsqrt(1.0_MK/Len)
55
     Leninv = 1.0_MK/Len
      !HPF$ INDEPENDENT
56
     do istep=Tmin,Tmax
57
        t = real(istep)*dt
        if(t .EQ. 0.0) then
58
           ! exact initial condition to avoid Gibbs oscillations
59
           u(1:N,1:N) = 1.0_MK
60
           u(iblx:ibux,ibly:ibuy) = 0.0_MK
61
        else
                _____
            ! Calculate solution for all points by Eigenfunction series expansion
62
           do j=0,N-1
63
             y = px(j+1)
I1 = c*Leninv
64
65
              I4 = (Len-d)*Leninv
66
              do l=M,1,-1
                                    ! sum in reverse order to avoid extinction
                 11 = dexp(-Diff*t*(real(1)*PI*Leninv)**2)
67
                 h = (2.0_MK/(real(1)*PI))*dcos(real(1)*PI*y*Leninv)
68
```

```
69
                 I1 = I1 + (h*dsin(real(l)*PI*c*Leninv)*ll)
70
                 I4 = I4 - (h*dsin(real(1)*PI*d*Leninv)*ll)
               end do
71
72
               do i=0,N-1
73
                 x = px(i+1)
74
                 I2 = 0.0_MK
75
                 I3 = 0.0_{MK}
76
                 do k=M,0,-1
                                     ! sum in reverse order to avoid extinction
77
                    if(k .EQ. 0) then
78
                       ukx = w1ul
79
                       t1 = ukx*a
80
                       t2 = ukx*(Len-b)
81
                    else
82
                       ukx = w2ul*dcos(real(k)*PI*x*Leninv)
83
                       h = w21/(real(k)*PI)
84
                       t1 = h*dsin(real(k)*PI*a*Leninv)
                       t2 = -h*dsin(real(k)*PI*b*Leninv)
85
86
                     end if
                    do l=M,0,-1
87
                                      ! sum in reverse order to avoid extinction
                       11 = (real(k)*PI*Leninv)**2
88
                       11 = 11 + (real(1)*PI*Leninv)**2
89
90
                       ll = dexp(-Diff*t*ll)
                       if(1 .EQ. 0) then
91
92
                          uly = w1ul
93
                          t3 = uly*(d-c)
94
                       else
95
                          uly = w2ul*dcos(real(1)*PI*y*Leninv)
96
                          t3 = dsin(real(1)*PI*d*Leninv)
                          t3 = t3-dsin(real(1)*PI*c*Leninv)
97
98
                          t3 = (w21/(real(1)*PI))*t3
99
                       end if
100
                       im = ukx*ulv*t3*11
                       I2 = I2 + im*t1
101
                       I3 = I3 + im * t2
102
                    end do
103
104
                 end do
                 u(i+1,j+1) = c0*(I1+I2+I3+I4)
105
106
              end do
107
           end do
         end if
108
         ! get current FRAP value and total mass
         1-----
109
         frap(istep-Tmin+1) = 0.0_MK
         ! loop over all cells in the bleached box
110
         do j=iblx,ibux-1
111
            do k=ibly,ibuy-1
112
              xm = u(j,k)+u(j+1,k)+u(j,k+1)+u(j+1,k+1)
113
              frap(istep-Tmin+1) = frap(istep-Tmin+1)+xm
114
           end do
115
         end do
         ! multiply with cell volume to get total mass in bleached box
116
         frap(istep-Tmin+1) = dx*dx*0.25_MK*frap(istep-Tmin+1)
         ! divide by bleached box volume to get concentration
117
         frap(istep-Tmin+1) = frap(istep-Tmin+1)/((b-a)*(d-c))
118
         mass(istep-Tmin+1) = 0.0_MK
         ! loop over all cells
119
         do j=1,N-1
120
            do k=1,N-1
            ! average the concentrations of all 4 corner vertices of the cell
            ! and multiply with the cells volume to get the mass
121
              xm = u(j,k)+u(j+1,k)+u(j,k+1)+u(j+1,k+1)
              mass(istep-Tmin+1) = mass(istep-Tmin+1)+xm
122
123
            end do
124
         end do
         mass(istep-Tmin+1) = dx*dx*0.25_MK*mass(istep-Tmin+1)
125
         ! Write solution at current time to file
126
         WRITE(filename, '(A, I3.3, A)') 'u', istep, '.out'
         OPEN(30, FILE=filename, STATUS='REPLACE', ACTION='WRITE')
127
128
         do j=1,N
129
            do i=1,N
130
              WRITE(30,*) px(i), px(j), u(i,j)
            end do
131
```

```
132
            WRITE(30,'(A)')
133
         end do
134
         CLOSE(30)
         WRITE(*,'(A,I5,A,I5)') 'Completed time step ',istep+1,' of ',Tmax+1
135
136
      end do
      ! Write FRAP data to file
      OPEN(30, FILE='frap.out', STATUS='REPLACE', ACTION='WRITE')
137
138
      do i=Tmin+1.Tmax+1
139
         WRITE(30,*) (i-1)*dt, frap(i-Tmin)
140
      end do
      CLOSE(30)
141
      OPEN(30, FILE='mass.out', STATUS='REPLACE', ACTION='WRITE')
142
      do i=Tmin+1,Tmax+1
143
         WRITE(30,*) (i-1)*dt, mass(i-Tmin)
144
145
      end do
146
      CLOSE(30)
      ! Deallocate and terminate
147
      if(ALLOCATED(u)) DEALLOCATE(u)
      if(ALLOCATED(px)) DEALLOCATE(px)
148
149
      if(ALLOCATED(frap)) DEALLOCATE(frap)
150
      if(ALLOCATED(mass)) DEALLOCATE(mass)
      9999 CONTINUE
151
```

```
152 END program analytical
```

C.2 A finite difference code

C.2.1 General description and usage notes

This finite difference code has been used in section 8.3 to validate the analytic solution for the box test case. The code only consists of a single source file and has the input parameters that are directly hard-coded in statements 30 through 49. A list of them is given in table C.2.

The boundary condition types that can be chosen for LBDCND, MBDCND and NBDCND are: 0 (for periodic boundary conditions), 1 (for Dirichlet) or 3 (for Neumann). The boundary values are assumed to be zero but they could be explicitly set in statements 88 to 93. The initial condition is chosen to be given by equation 8.1 if INCOND is set to .FALSE.. If it is set to .TRUE., the initial condition is read from a file called init.in that contains nv(1)*nv(2) lines giving the initial values on the grid points as:

```
<x-pos> <y-pos> <value u0>
```

where the inner loop is over nv(2), the outer over nv(1). The initial condition is then extruded in z direction (i.e. it is assumed to be the same on all z-planes). This enables the exact solution to serve as an initial condition as was done for figure 8.8.

The output consists of the three files init.out, fd.frap and fd.mass. init.out contains the positions of all the grid points that have an initial value of zero using the following syntax:

```
for i=1,nv(1)
  for j=1,nv(2)
    for k=1,nv(3)
        if(u(i,j,k)=0)
```

```
 <x-pos> <y-pos> <z-pos>
    end if
    end for
  end for
end for
```

It mainly serves for control purposes. The files fd.frap and fd.mass contain the FRAP values and the total mass in the system versus time using the syntax:

```
<time> <value>
```

on every line. They are used to create the various FRAP curve plots as well as to check the accuracy of the simulation using the principle of conservation of mass. Examples are given in figures 8.6 and 8.7, respectively. In addition to above output files, the code also prints some messages to the system's standard output. After successful initialization, it prints:

```
Solver initialized.
```

Then, it prints a message like:

Completed time step 20 of 6000

every time a time step is completed. At the very end, the following output is given:

Solver finished.

Parameter	Meaning
D(D)	Diffusion constant
bblx(a)	Lower x-coordinate of the bleached box
bbux (b)	Upper x-coordinate of the bleached box
bbly (c)	Lower y-coordinate of the bleached box
bbuy (d)	Upper y-coordinate of the bleached box
dt (δt)	Time step size
tfinal (t_{final})	Final time of solution
mmin(1)	Lower x-coordinate of cubic computational domain
mmin(2)	Lower y-coordinate of cubic computational domain
mmin(3)	Lower z-coordinate of cubic computational domain
mmax(1)	Upper x-coordinate of cubic computational domain
mmax(2)	Upper y-coordinate of cubic computational domain
mmax(3)	Upper z-coordinate of cubic computational domain
nv(1)	Number of grid points in x direction
nv(2)	Number of grid points in y direction
nv(3)	Number of grid points in z direction
LBDCND	Boundary condition type in x direction
MBDCND	Boundary condition type in y direction
NBDCND	Boundary condition type in z direction
INCOND	Read initial condition from file or not

Table C.2: Input parameters for fd.f90

C.2.2 Source code listing

The source code of the program used to calculate the finite difference solution for the box test case in section 8.3 is given hereafter. It is based on the fast commercial solver library CRAYFISHPAK by Green Mountain software, Madeira Beach, FL, which is needed for the program to compile and run.

C.2.2.1 Finite difference solver (fd.f90)

```
1
     program fd
2
     IMPLICIT NONE
      !FD Calculates a high-resolution finite difference solution for the box
         \ensuremath{\texttt{FD}}\xspace uses CRAYFISHPAK solvers to get a finite difference solution
         for the box test case in order to validate the analytic solution.
         THIS NEEDS THE CRAYFISHPAK LIBRARIES
         (commercial and thus not included in this reprint)
         See also
         todo:
      1------
                                                             ------
      ! DIPLOMA THESIS WS01/02 ICOS
                                                                   ETH-ZUERICH
                   PROTEIN DIFFUSION INSIDE THE ENDOPLASMIC RETICULUM
           ! set precision
3
     INTEGER, PARAMETER
                                           :: MK = KIND(1.0D0)
                                  _____
      ! Declaration of local variables
      ! loop conters
4
     INTEGER
                                           :: i, j ,k, l
     ! lower and upper boundaries of physical domain
     REAL(MK), DIMENSION(3)
5
                                           :: mmin, mmax
     ! number of grid points in all 3 directions
6
     INTEGER, DIMENSION(3)
                                           :: nv
      ! boundary condition types (0: periodic, 1: dirichlet, 3: neumann)
     INTEGER
                                           :: LBDCND
7
     INTEGER
                                           :: MBDCND
8
     INTEGER
                                           :: NBDCND
9
     ! boundary data arrays
10
     REAL(MK), DIMENSION(:,:), ALLOCATABLE :: BDXS, BDXF, BDYS, BDYF, BDZS, BDZF
     ! dimensions
     INTEGER
                                           :: LDIMF, MDIMF
11
     ! lambda in Helmholz equation
     REAL(MK)
12
                                           :: lambda
      ! error trap
13
     INTEGER
                                           :: info
     ! solution array
     REAL(MK), DIMENSION(:,:,:), ALLOCATABLE :: u
14
      ! work arrays for CRAYFISHPAK
     REAL(MK), DIMENSION(:), ALLOCATABLE
                                           :: wrk1, wrk2
15
      ! time step size
16
     REAL(MK)
                                           :: dt
      ! final time of simulation
17
     REAL(MK)
                                           :: tfinal
     ! number of time steps
18
     INTEGER
                                           :: TMAX
     ! diffusion coefficient
19
     REAL(MK)
                                           :: D
     ! dimension of work array needed by CRAYFISHPAK
20
     INTEGER
                                           :: dim
      ! Perturbation returned from solver
21
     REAL(MK)
                                           :: PERTRB
     ! grid spacing in x, y and z
```

C.2. A FINITE DIFFERENCE CODE

```
22
     REAL(MK)
                                            :: dx, dy, dz
      ! lower/upper x/y coordinates of bleached box
23
     REAL(MK)
                                            :: bblx, bbux, bbly, bbuy
      ! lower and upper indices of bleached box in solution grid
24
     INTEGER
                                            :: il, ih, jl, jh
     ! number of points inside bleached box
25
     INTEGER
                                            :: nb
      ! array to hold frap data
26
     REAL(MK), DIMENSION(:), ALLOCATABLE
                                            :: frap
     ! array to hold total mass at each time step
27
     REAL(MK), DIMENSION(:), ALLOCATABLE
                                           :: mass
      ! read initial condition from file?
                                            :: INCOND
28
     LOGICAL
     ! moving sum for total mass
29
     REAL(MK)
                                            :: xm
      ! Problem initialization (THIS IS USER INPUT)
                                                   ! diffusion coefficient
30
     D = 0.03_MK
     ! bleached box
     bblx = 2.0_MK
31
     bbux = 3.0 MK
32
     bbly = 2.0 MK
33
     bbuy = 3.0 MK
34
     ! time step and final time
dt = 0.005_MK
35
     tfinal = 30.0_MK
36
     ! physical domain size
37
     mmin(1) = 0.0 MK
     mmin(2) = 0.0_MK
38
     mmin(3) = 0.0 MK
39
     mmax(1) = 4.0_MK
mmax(2) = 4.0_MK
40
41
     mmax(3) = 4.0 MK
42
     ! number of grid vertices in all three directions
43
     nv(1) = 81
44
     nv(2) = 81
     nv(3) = 81
45
      ! boundary condition: Neumann (3) on all sides
     ! (1 would be Dirichlet and 0 periodic)
46
     LBDCND = 3
     MBDCND = 3
47
     NBDCND = 3
48
      ! read initial condition from file ?
49
     INCOND = .false.
      1------
      ! Initialize derived properties (NO USER INPUT BELOW THIS LINE)
      !----
     ! grid spacings
50
     dx = (mmax(1) - mmin(1)) / real(nv(1) - 1)
51
     dy = (mmax(2)-mmin(2))/real(nv(2)-1)
52
     dz = (mmax(3)-mmin(3))/real(nv(3)-1)
      ! number of time steps
53
     TMAX = ceiling(tfinal/dt)
      ! lambda in Helmholtz equation
     lambda = -1.0_MK/(D*dt)
54
      ! some dimensions (ex CRAYFISHPAK documentation)
     LDIMF = nv(1) + 4
55
     MDIMF = nv(2) + 4
56
57
     dim = (nv(1)+1)*(nv(2))*(nv(3)) + \&
              7*(nv(1)-1) + 3*(nv(2)-1) + 4*(nv(3)-1) + 76
      ! Allocate the memory needed
```

```
58
      ALLOCATE(u(LDIMF,MDIMF,nv(3)), STAT=info)
59
      if(info .NE. 0) then
60
         WRITE(*,'(A)') 'Error allocating memory for solution array u.'
         goto 9999
61
62
      endif
63
      ALLOCATE(wrk1(dim), wrk2(dim), STAT=info)
64
      if(info .NE. 0) then
65
         WRITE(*,'(A)') 'Error allocating memory for work arrays.'
66
         goto 9999
67
      end if
68
      ALLOCATE(BDXS(nv(2)+4,nv(3)+4), BDXF(nv(2)+4,nv(3)+4), STAT=info)
69
      if(info .NE. 0) then
70
         WRITE(*,'(A)') 'Error allocating memory for x boundary arrays.'
71
         goto 9999
72
      end if
73
      ALLOCATE(BDYS(nv(1)+4,nv(3)+4), BDYF(nv(1)+4,nv(3)+4), STAT=info)
74
      if(info .NE. 0) then
        WRITE(*,'(A)') 'Error allocating memory for y boundary arrays.'
75
76
         goto 9999
77
      end if
78
      ALLOCATE(BDZS(nv(1)+4,nv(2)+4), BDZF(nv(1)+4,nv(2)+4), STAT=info)
79
      if(info .NE. 0) then
         WRITE(*,'(A)') 'Error allocating memory for z boundary arrays.'
80
81
         goto 9999
82
      end if
83
      ALLOCATE(frap(TMAX+1), mass(TMAX+1), STAT=info)
      if(info .NE. 0) then
    WRITE(*,'(A)') 'Error allocating memory for frap and mass arrays.'
84
85
         goto 9999
86
87
      end if
      ! Set boundary and initial conditions
      1---
      ! boundary condition values (s: lower boundary, f: upper boundary)
      ! here meaning zero gradient at the boundaries
88
      BDXS = 0.0 MK
      BDXF = 0.0_MK
BDYS = 0.0_MK
89
90
91
      BDYF = 0.0_MK
92
      BDZS = 0.0_MK
      BDZF = 0.0_MK
93
      ! set initial condition
94
      u = 1.0_MK
95
      il = ceiling(bblx/dx)+1
                                  ! boundaries of bleached box
96
      ih = floor(bbux/dx)+1
97
      jl = ceiling(bbly/dy)+1
98
      jh = floor(bbuy/dy)+1
99
      nb = (ih-il+1)*(jh-jl+1)*nv(3) ! number of particles in bleached box
100
      if(INCOND) then
101
         OPEN(40, FILE='init.in', STATUS='OLD', ACTION='READ')
         ! read initial condition from file
102
        do i=1,nv(1)
103
           do j=1,nv(2)
104
               READ(40,*) xm, xm, u(i,j,1)
105
            end do
         end do
106
107
         CLOSE(40)
         ! extrude to 3D
108
         do i=1,nv(1)
109
           do j=1,nv(2)
110
               do k=2,nv(3)
                 u(i,j,k) = u(i,j,1)
111
112
               end do
113
            end do
114
         end do
115
      else
        u(il:ih,jl:jh,:) = 0.0_MK
116
117
      end if
      ! write initial condition to file for check
118
      OPEN(40,FILE='init.out',STATUS='REPLACE',ACTION='WRITE')
     do i=1,LDIMF
119
```

```
120
        do j=1,MDIMF
121
           do k=1,nv(3)
122
              if(u(i,j,k) .EQ. 0) WRITE(40,*) (i-1)*dx,(j-1)*dy,(k-1)*dx
123
           enddo
124
        enddo
125
      enddo
126
     CLOSE(40)
      1------
      ! Initialize solver (CALL TO CRAYFISHPAK LIBRARY)
                                                          _____
     CALL H3GCIS(mmin(1),mmax(1),nv(1)-1,LBDCND, &
127
                 mmin(2),mmax(2),nv(2)-1,MBDCND,
                                                  &
                 mmin(3), mmax(3), nv(3)-1, NBDCND,
                                                  &
                 lambda,LDIMF,MDIMF,info,wrk1)
     if(info .NE. 0) then
128
        WRITE(*,'(A)') 'ERROR: Initialization of solver failed!'
129
        WRITE(*, '(A, I5)') 'H3GCIS returns: ', info
130
131
        goto 9999
132
     else
133
        WRITE(*,'(A)') 'Solver initialized.'
134
     end if
      ! Main time step loop
      1
135
     do i=1,TMAX
        frap(i) = 0.0_MK
136
        ! loop over all cells in the bleached box
137
        do j=il,ih-1
138
           do k=jl,jh-1
              do l=1,nv(3)-1
139
140
                 xm = u(j,k,l)+u(j+1,k,l)+u(j,k+1,l)+u(j,k,l+1) &
                      +u(j+1,k+1,l)+u(j,k+1,l+1)+u(j+1,k,l+1)+u(j+1,k+1,l+1)
141
                 frap(i) = frap(i) + xm
              end do
142
           end do
143
        end do
144
         ! multiply with cell volume to get total mass in bleached box
145
        frap(i) = dx*dy*dz*0.125_MK*frap(i)
         ! divide by bleached box volume to get concentration
        frap(i) = frap(i)/((bbux-bblx)*(bbuy-bbly)*(mmax(3)-mmin(3)))
146
        mass(i) = 0.0_MK
147
        ! loop over all cells
148
        do j=1,nv(1)-1
149
           do k=1,nv(2)-1
150
              do l=1.nv(3)-1
              ! average the concentrations of all 8 corner vertices of the cell
151
                 xm = u(j,k,l)+u(j+1,k,l)+u(j,k+1,l)+u(j,k,l+1) &
                      +u(j+1,k+1,l)+u(j,k+1,l+1)+u(j+1,k,l+1)+u(j+1,k+1,l+1)
152
                 mass(i) = mass(i) + xm
153
              end do
154
           end do
155
        end do
         ! multiply with the cells volume to get the mass
156
        mass(i) = dx*dy*dz*0.125_MK*mass(i)
         ! Call solver (CALL TO CRAYFISHPAK LIBRARY)
         ! reset perturbation
        PERTRB = 0.0_MK
157
        ! new right hand side
158
        u = lambda*u
        ! solve
159
        CALL H3GCSS(BDXS,BDXF,BDYS,BDYF,BDZS,BDZF,LDIMF,MDIMF,u(1,1,1), &
                    PERTRB, wrk1, wrk2)
        if(PERTRB .NE. 0.0_MK) then
160
           WRITE(*,'(A,I3,A,E15.5)') 'Time step ',i,': Solver perturbation = ',PERTRB
161
162
        end if
        WRITE(*,'(A,I5,A,I5)') 'Completed time step ',i,' of ',TMAX
163
164
     end do
      ! get last value too
```

```
165
     frap(TMAX+1) = 0.0_MK
166
      do j=il,ih-1
167
         do k=jl,jh-1
168
            do l=1,nv(3)-1
169
               xm = u(j,k,l)+u(j+1,k,l)+u(j,k+1,l)+u(j,k,l+1) &
                    +u(j+1,k+1,1)+u(j,k+1,1+1)+u(j+1,k,1+1)+u(j+1,k+1,1+1)
170
               frap(TMAX+1) = frap(TMAX+1) + xm
171
            end do
172
         end do
173
      end do
       ! multiply with cell volume to get total mass in bleached box
174
      frap(TMAX+1) = dx*dy*dz*0.125_MK*frap(TMAX+1)
      ! divide by bleached box volume to get concentration
175
      frap(TMAX+1) = frap(TMAX+1)/((bbux-bblx)*(bbuy-bbly)*(mmax(3)-mmin(3)))
176
      mass(TMAX+1) = 0.0_MK
177
      do j=1,nv(1)-1
178
         do k=1,nv(2)-1
179
            do l=1,nv(3)-1
               xm = u(j,k,l)+u(j+1,k,l)+u(j,k+1,l)+u(j,k,l+1) &
180
                    +u(j+1,k+1,l)+u(j,k+1,l+1)+u(j+1,k,l+1)+u(j+1,k+1,l+1)
181
               mass(TMAX+1) = mass(TMAX+1) + xm
182
            end do
183
         end do
184
      end do
185
     mass(TMAX+1) = dx*dy*dz*0.125_MK*mass(TMAX+1)
      ! write frap data to file
      OPEN(40, FILE='fd.frap', STATUS='REPLACE', ACTION='WRITE')
186
      do i=1,TMAX+1
187
188
         WRITE(40,*) (i-1)*dt, frap(i)
189
      end do
      CLOSE(40)
190
      OPEN(40, FILE='fd.mass', STATUS='REPLACE', ACTION='WRITE')
191
192
      do i=1,TMAX+1
        WRITE(40,*) (i-1)*dt, mass(i)
193
194
      end do
      CLOSE(40)
195
      ! Free memory and terminate
      1----
     9999 CONTINUE
196
     if(ALLOCATED(mass)) DEALLOCATE(mass)
197
198
      if(ALLOCATED(frap)) DEALLOCATE(frap)
199
     if(ALLOCATED(BDZF)) DEALLOCATE(BDZF)
200
     if(ALLOCATED(BDZS)) DEALLOCATE(BDZS)
201
     if(ALLOCATED(BDYF)) DEALLOCATE(BDYF)
202
     if(ALLOCATED(BDYS)) DEALLOCATE(BDYS)
      if(ALLOCATED(BDXF)) DEALLOCATE(BDXF)
203
204
      if(ALLOCATED(BDXS)) DEALLOCATE(BDXS)
205
      if(ALLOCATED(wrk2)) DEALLOCATE(wrk2)
206
      if(ALLOCATED(wrk1)) DEALLOCATE(wrk1)
207
      if(ALLOCATED(u)) DEALLOCATE(u)
208
      WRITE(*,'(A)') 'Solver finished.'
209
      END program fd
```

C.3 Post-processing

C.3.1 General description and usage notes

As mentioned in section 7.3, the output files of the PSE simulation code have to be converted to an OpenDX grid file format in order to visualize the concentration fields. The code that implements the corresponding algorithm is given hereafter. It consists of a single source file and does not need any input parameters. The files one wishes to convert are passed to it as command line arguments upon invocation:

\$ res2dx <file1.res> [file2.res [file3.res ...]]

Any number of files can be processed at the same call, shell wildcards are allowed. The parameters NBX, NBY and NBZ on lines 4 to 6 give the numbers of cells in each direction to be used for the linked list algorithm. They can be tuned for better performance. The input files are PSE3D result files containing all the particles and their strengths at a certain time step of the simulation. Each line in such a file contains the information about one particle using the syntax:

```
<x-pos> <y-pos> <z-pos> <strength>
```

For each input file processed, the code outputs the two files **<infile>.grid** and **<infile>.general**. The former contains just one number per line giving the concentration value at that grid position. Hereby it first loops over the z-direction, then over the y-direction and last over the x-direction:

```
for i=1,Nx
  for j=1,Ny
    for k=1,Nz
        <value(i,j,k)>
        end for
    end for
end for
```

The latter contains the OpenDX header according to section 7.3. This is the file that is actually opened in OpenDX. No output is written to the standard output.

C.3.2 Source code listing

The source code of the converter program is given below. For a description of its algorithms including the fast linked list sorting, see section 7.3 of this report.

C.3.2.1 OpenDX grid file converter (res2dx.f90)

```
program res2dx
1
      IMPLICIT NONE
2
      !RES2DX Converts PSE3D result files to OpenDX grid input format.
         res2dx reads a file containing positions and strengths of particles
         and converts them to a regular cartesion grid file for OpenDX. It
         writes both the data file *.grid and the *.general file containing
         all the headers for OpenDX. The list of files to be processed is
         passed to res2dx as command line arguments. This program uses a
         fast o(N) box sorting with linked particle lists
         See also PSE3D
         todo:
                 _____
        DIPLOMA THESIS WS01/02 ICOS
                                                                       ETH-ZUERICH
                    PROTEIN DIFFUSION INSIDE THE ENDOPLASMIC RETICULUM
                         ========= ivo f. sbalzarini ======
      ! Declaration of external functions
       get number of command line arguments
3
      INTEGER, EXTERNAL
                                              :: iargc
```

	!			
	! Declaration of local variables			
	!			
	I number of boxes in x y z for chaining	mesh		
4	INTEGER. PARAMETER	$\therefore BX = 10$		
5	INTEGER, PARAMETER	:: NBY = 10		
6	Integer, PARAMETER	:: NBZ = 10		
	! accuarcy			
7	INTEGER, PARAMETER	:: MK = KIND(1.0D0)		
0	! loop counters			
0	INTEGER	:: 1,], K, 1, III		
9	REAL(MK)	:: dx. dv. dz		
	! number of grid points in x,y,z			
10	INTEGER	:: Nx, Ny, Nz		
	! bounding box			
11	REAL(MK)	:: xl,xh,yl,yh,zl,zh		
12	INTEGER	·· Np Nw		
12	! position	·· · · · · · · · · · · · · · · · · · ·		
13	REAL(MK), DIMENSION(3)	:: x, xold, xdiff		
	! strength			
14	REAL(MK)	:: s		
	! input file, output file, header file			
15	CHARACTER(LEN=80)	:: infile, outfile, headfile		
16	CHARACTER(LEN=80)	:: prgname		
10	! number of command line arguments (inpu	it files)		
17	INTEGER	:: nargc		
	<pre>! position and mass data</pre>			
18	REAL(MK), DIMENSION(:,:), ALLOCATABLE	:: xp		
19	REAL(MK), DIMENSION(:), ALLOCATABLE	:: mass		
20	! threshold for hearest grid point REAL(MK)	•• +		
20	! is this grid point found in the data?			
21	LOGICAL	:: found		
	! error status			
22	INTEGER	:: istat		
00	! head-of-chain for each box	1100		
23	INTEGER, DIMENSION(:,:,:), ALLUCATABLE	:: HUC		
24	INTEGER, DIMENSION(:), ALLOCATABLE	:: 11		
	! box sizes			
25	REAL(MK)	:: dbx, dby, dbz		
	! current box index			
26	INTEGER	:: idx, idy, idz		
	1			
	! Process command line arguments			
	!			
	! read program name from command line			
27	CALL getarg(0,prgname)			
	I got number of input files			
28	nargc = iargc()			
29	if(nargc .LT. 1) then			
30	print*,'Missing input file name(s). U	Jsage: ',TRIM(prgname),' &		
	<pre>inputfile [inputfile]'</pre>			
31	goto 9999			
32	end if			
	1			
	! Loop over all input files and translat	te them		
	!			
33	do i=1,nargc			
2/	! get input file name			
54	! derive other file names			
35	WRITE(outfile,'(2A)') TRIM(infile).'.	grid'		
36	WRITE(headfile,'(2A)') TRIM(infile),'	- .general'		
37	<pre>WRITE(*,'(2A)') 'processing file: ',T</pre>	TRIM(infile)		
	I Scon file and determine mid and	·····		
	: Scan ille and determine grid geomet	,		
	-			
38	WRITE(*,'(A)') 'Parsing data to deter	mine grid geometry'		

C.3. POST-PROCESSING

39	UPEN(40, FILE=infile, STATUS='ULD', ACTIUN='READ')
	! first pass: scan file and count number of particles to be
	! read in. also determine bounding box and grid spacings
40	Np = 0
41	xh = -HUGE(xh)
42	xl = HUGE(xl)
43	vh = -HUGE(vh)
44	$v_1 = HUGE(v_1)$
45	$z_{\rm b} = -\text{HIGE}(z_{\rm b})$
46	$z_1 = \operatorname{Higg}(z_1)$
47	$d_{\mathbf{x}} = \operatorname{Higg}(d_{\mathbf{x}})$
10	dx = High(dx)
10	dy = High(dx)
4 <i>3</i> 50	
50	PEAD(AO)(AE16 8)(END-110) = (1.3) = 0
52	$if(N_{\rm E} = 0)$ vold - x
52	$N_{\rm P} = N_{\rm P} \pm 1$
55	np - np + 1
54	xdiii = abs(x-xoid)
55	II(X(I), GI, XI) XI = X(I)
56	II(X(1) . LI . XI) XI = X(1)
57	if(x(2), GT, yh) yh = x(2)
58	if(x(2) . LT. y1) y1 = x(2)
59	1f(x(3) . GT. zh) zh = x(3)
60	if(x(3) .LT. z1) z1 = x(3)
61	if(xdiff(1) .LT. dx .AND. xdiff(1) .GT. 0.0_MK) dx = xdiff(1)
62	$if(xdiff(2) .LT. dy .AND. xdiff(2) .GT. 0.0_MK) dy = xdiff(2)$
63	$if(xdiff(3) .LT. dz .AND. xdiff(3) .GT. 0.0_MK) dz = xdiff(3)$
64	xold = x
65	end do
66	110 CLOSE(40)
	! determine number of grid points
67	Nx = int((xh-xl)/dx)+1
68	Ny = int((yh-yl)/dy)+1
69	Nz = int((zh-zl)/dz)+1
	! print terminal output
70	WRITE(*, (3(A, F8.4))))) 'xl = ',xl, '/vl = ',vl, '/zl = ',zl
71	WRITE $(*, (3(A, F8, 4)))$ $dx = (dx, (dx, (dx - 1)))$ $dz = (dx, (dx - 1))$
72	WRITE $(*, (3(A, IB)))$ $N_x = (N_x, (N_y, N_y))$ $N_z = (N_y, (N_y, N_y))$
73	
	WB B K ' S A B A J ' J ' Y = ' Y ' / V = ' V ' / Z = ' Z
15	WRITE(*, (3(A,F8.4)))) Xn = ',Xn,' / yn = ',yn,' / 2n = ',2n
10	WRITE(*, $(S(A, ro.4))$) 'xn = ',xn,' ' yn = ',yn,' / zn = ',zn
74	<pre>wmile(*, '(3(A,ro.4))') 'xn = ',xn,' ' yn = ',yn,' / zn = ',zn ! write OpenDX *.general header file OPEN(40_FILE=boodfile_STATUS='PED(ACE'_ACTION='UDITE'))</pre>
74	<pre>WRITE(*, '(S(A,FO.4))') 'xn = ',xn,' ' yn = ',yn,' / zn = ',zn ! write OpenDX *.general header file OPEN(40, FILE=headfile, STATUS='REPLACE', ACTION='WRITE') UPITE(40, (O)) 'file = ', TETM(outfile)</pre>
74 75 76	<pre>wRITE(*, '(3(A,F0.4))') 'xn = ',xn,' ' yn = ',yn,' / zn = ',zn ! write OpenDX *.general header file OPEN(40, FILE=headfile, STATUS='REPLACE', ACTION='WRITE') WRITE(40, '(2A)') 'file = ', TRIM(outfile) UPITE(40, '(A I A A I A I A)') 'crid = ', Nr '</pre>
74 75 76	<pre>WRITE(*, '(3(A,F0.4))') 'xn = ',xn,' ' yn = ',yn,' / zn = ',zn ! write OpenDX *.general header file OPEN(40, FILE=headfile, STATUS='REPLACE', ACTION='WRITE') WRITE(40,'(2A)') 'file = ', TRIM(outfile) WRITE(40,'(A,I4,A,I4,A,I4)') 'grid = ',Nx,' x ',Ny,' x ',Nz UDITE(40,'(A),') 'fine=t = continue</pre>
74 75 76 77	<pre>WRITE(*, '(3(A,F0.4))') 'xn = ',xn,' ' yn = ',yn,' / 2n = ',2n ! write OpenDX *.general header file OPEN(40, FILE=headfile, STATUS='REPLACE', ACTION='WRITE') WRITE(40,'(2A)') 'file = ', TRIM(outfile) WRITE(40,'(A,I4,A,I4,A,I4)') 'grid = ',Nx,' x ',Ny,' x ',Nz WRITE(40,'(A)') 'format = ascii' WRITE(40,'(A)') 'format = ascii'</pre>
74 75 76 77 78	<pre>WRITE(*, '(3(A,F5.4))') 'xn = ',xn,' ' yn = ',yn,' / 2n = ',2n ! write OpenDX *.general header file OPEN(40, FILE=headfile, STATUS='REPLACE', ACTION='WRITE') WRITE(40,'(2A)') 'file = ', TRIM(outfile) WRITE(40,'(A,I4,A,I4,A,I4)') 'grid = ',Nx,' x ',Ny,' x ',Nz WRITE(40,'(A)') 'format = ascii' WRITE(40,'(A)') 'interleaving = record' WRITE(40,'(A)') 'interleaving = record'</pre>
74 75 76 77 78 79	<pre>WRITE(*, '(3(A,F0.4))') 'xn = ',xn,' ' yn = ',yn,' ' zn = ',zn ! write OpenDX *.general header file OPEN(40, FILE=headfile, STATUS='REPLACE', ACTION='WRITE') WRITE(40, '(2A)') 'file = ', TRIM(outfile) WRITE(40, '(A,I4,A,I4,A,I4)') 'grid = ',Nx,' x ',Ny,' x ',Nz WRITE(40,'(A)') 'finat = ascii' WRITE(40,'(A)') 'interleaving = record' WRITE(40,'(A)') 'majority = row' WRITE(40,'(A)') 'majority = row'</pre>
74 75 76 77 78 79 80	<pre>WRITE(*, '(3(A,F0.4))') 'xn = ',xn,' ' yn = ',yn,' ' zn = ',zn ! write OpenDX *.general header file OPEN(40, FILE=headfile, STATUS='REPLACE', ACTION='WRITE') WRITE(40, '(2A)') 'file = ', TRIM(outfile) WRITE(40, '(A,I4,A,I4)') 'grid = ',Nx,' x ',Ny,' x ',Nz WRITE(40, '(A)') 'format = ascii' WRITE(40, '(A)') 'interleaving = record' WRITE(40, '(A)') 'majority = row' WRITE(40, '(A)') 'field = u' WRITE(40, '(A)') 'field = u'</pre>
74 75 76 77 78 79 80 81	<pre>WRITE(*, '(3(A,F0.4))') 'xn = ',xn,' ' yn = ',yn,' ' zn = ',zn ! write OpenDX *.general header file OPEN(40, FILE=headfile, STATUS='REPLACE', ACTION='WRITE') WRITE(40, '(2A)') 'file = ', TRIM(outfile) WRITE(40, '(A,I4,A,I4,A,I4)') 'grid = ',Nx,' x ',Ny,' x ',Nz WRITE(40, '(A)') 'format = ascii' WRITE(40, '(A)') 'interleaving = record' WRITE(40, '(A)') 'majority = row' WRITE(40, '(A)') 'field = u' WRITE(40, '(A)') 'structure = scalar'</pre>
74 75 76 77 78 79 80 81 82	<pre>WRITE(*, '(3(A,F5.4))') 'xn = ',xn,' / yn = ',yn,' / 2n = ',2n ! write OpenDX *.general header file OPEN(40, FILE=headfile, STATUS='REPLACE', ACTION='WRITE') WRITE(40,'(A)') 'file = ', TRIM(outfile) WRITE(40,'(A,I4,A,I4,)') 'grid = ',Nx,' x ',Ny,' x ',Nz WRITE(40,'(A)') 'format = ascii' WRITE(40,'(A)') 'interleaving = record' WRITE(40,'(A)') 'interleaving = record' WRITE(40,'(A)') 'majority = row' WRITE(40,'(A)') 'field = u' WRITE(40,'(A)') 'structure = scalar' WRITE(40,'(A)') 'type = float'</pre>
74 75 76 77 78 79 80 81 82 83	<pre>WRITE(*, '(3(A,F0.4))') 'xn = ',xn,' / yn = ',yn,' / zn = ',zn ! write OpenDX *.general header file OPEN(40, FILE=headfile, STATUS='REPLACE', ACTION='WRITE') WRITE(40, '(2A)') 'file = ', TRIM(outfile) WRITE(40, '(A,I4,A,I4,A,I4)') 'grid = ',Nx,' x ',Ny,' x ',Nz WRITE(40, '(A)') 'findra = ascii' WRITE(40, '(A)') 'interleaving = record' WRITE(40, '(A)') 'majority = row' WRITE(40, '(A)') 'field = u' WRITE(40, '(A)') 'field = u' WRITE(40, '(A)') 'structure = scalar' WRITE(40, '(A)') 'type = float' WRITE(40, '(A)') 'dpendency = positions'</pre>
74 75 76 77 78 79 80 81 82 83 84	<pre>WRITE(*, '(3(A,F0.4))') 'xn = ',xn,' 'yn = ',yn,' ' zn = ',zn ! write OpenDX *.general header file OPEN(40, FILE=headfile, STATUS='REPLACE', ACTION='WRITE') WRITE(40, '(2A)') 'file = ', TRIM(outfile) WRITE(40, '(A,I4,A,I4,A,I4)') 'grid = ',Nx,' x ',Ny,' x ',Nz WRITE(40, '(A)') 'format = ascii' WRITE(40, '(A)') 'interleaving = record' WRITE(40, '(A)') 'interleaving = record' WRITE(40, '(A)') 'filed = u' WRITE(40, '(A)') 'filed = u' WRITE(40, '(A)') 'structure = scalar' WRITE(40, '(A)') 'type = float' WRITE(40, '(A)') 'dpendency = positions' WRITE(40, '(A,5(F8.3,A),F8.3)') 'positions = regular,regular,regular, &</pre>
74 75 76 77 78 79 80 81 82 83 84	<pre>WRITE(*, '(3(A,F0.4))') 'xn = ',xn,' 'yn = ',yn,' ' zn = ',zn ! write OpenDX *.general header file OPEN(40, FILE=headfile, STATUS='REPLACE', ACTION='WRITE') WRITE(40, '(2A)') 'file = ', TRIM(outfile) WRITE(40, '(A,I4,A,I4), ') 'grid = ',Nx,' x ',Ny,' x ',Nz WRITE(40, '(A)') 'format = ascii' WRITE(40, '(A)') 'interleaving = record' WRITE(40, '(A)') 'majority = row' WRITE(40, '(A)') 'field = u' WRITE(40, '(A)') 'structure = scalar' WRITE(40, '(A)') 'structure = scalar' WRITE(40, '(A)') 'dependency = positions' WRITE(40, '(A)') 'dependency = positions = regular, regular, &</pre>
74 75 76 77 78 79 80 81 82 83 84 85	<pre>WRITE(*, '(3(A,F5.4))') 'xn = ',xn,' / yn = ',yn,' / zn = ',zn ! write OpenDX *.general header file OPEN(40, FILE=headfile, STATUS='REPLACE', ACTION='WRITE') WRITE(40, '(A)') 'file = ', TRIM(outfile) WRITE(40, '(A)') 'format = ascii' WRITE(40, '(A)') 'interleaving = record' WRITE(40, '(A)') 'interleaving = record' WRITE(40, '(A)') 'majority = row' WRITE(40, '(A)') 'field = u' WRITE(40, '(A)') 'jterleaving = scalar' WRITE(40, '(A)') 'type = float' WRITE(40, '(A)') 'dependency = positions' WRITE(40, '(A,S(F8.3,A),F8.3)') 'positions = regular,regular,regular, &</pre>
74 75 76 77 78 79 80 81 82 83 84 85 86	<pre>WAILE(*, '(3(A,F6.4))') 'xn = ',xn,' / yn = ',yn,' / zn = ',zn ! write OpenDX *.general header file OPEN(40, FILE=headfile, STATUS='REPLACE', ACTION='WRITE') WRITE(40, '(2A)') 'file = ', TRIM(outfile) WRITE(40, '(A,I4,A,I4,A,I4)') 'grid = ',Nx,' x ',Ny,' x ',Nz WRITE(40, '(A)') 'format = ascii' WRITE(40, '(A)') 'interleaving = record' WRITE(40, '(A)') 'interleaving = record' WRITE(40, '(A)') 'field = u' WRITE(40, '(A)') 'field = u' WRITE(40, '(A)') 'field = u' WRITE(40, '(A)') 'type = float' WRITE(40, '(A)') 'dependency = positions' WRITE(40, '(A,5(F8.3,A),F8.3)') 'positions = regular,regular,regular, &</pre>
74 75 76 77 78 79 80 81 82 83 84 85 86 87	<pre>WATE(*, '(3(A,F6.4))') 'xn = ',xn,' / yn = ',yn,' / zn = ',zn ! write OpenDX *.general header file OPEN(40, FILE=headfile, STATUS='REPLACE', ACTION='WRITE') WRITE(40, '(2A)') 'file = ', TRIM(outfile) WRITE(40, '(A,I4,A,I4,A,I4)') 'grid = ',Nx,' x ',Ny,' x ',Nz WRITE(40,'(A)') 'format = ascii' WRITE(40,'(A)') 'filed = u WRITE(40,'(A)') 'filed = u' WRITE(40,'(A</pre>
74 75 76 77 78 80 81 82 83 84 85 86 87	<pre>WRITE(*, '(3(A,F6.4))') 'xn = ',xn,' / yn = ',yn,' / 2n = ',2n ! write OpenDX *.general header file OPEN(40, FILE=headfile, STATUS='REPLACE', ACTION='WRITE') WRITE(40, '(2A)') 'file = ', TRIM(outfile) WRITE(40, '(A,I4,A,I4,A,I4)') 'grid = ',Nx,' x ',Ny,' x ',Nz WRITE(40, '(A)') 'format = ascii' WRITE(40, '(A)') 'interleaving = record' WRITE(40, '(A)') 'majority = row' WRITE(40, '(A)') 'field = u' WRITE(40, '(A)') 'field = u' WRITE(40, '(A)') 'structure = scalar' WRITE(40, '(A)') 'dpe = float' WRITE(40, '(A)') 'dpeendency = positions' WRITE(40, '(A,5(F8.3,A),F8.3)') 'positions = regular,regular,regular, &</pre>
74 75 76 77 78 79 80 81 82 83 84 85 86 87	<pre>WAILE(*, '(3(A,F5.4))') 'xn = ',xn,' / yn = ',yn,' / 2n = ',2n ' write OpenDX *.general header file OPEN(40, FILE=headfile, STATUS='REPLACE', ACTION='WRITE') WRITE(40, '(A)') 'file = ', TRIM(outfile) WRITE(40, '(A)') 'format = ascii' WRITE(40, '(A)') 'interleaving = record' WRITE(40, '(A)') 'interleaving = record' WRITE(40, '(A)') 'majority = row' WRITE(40, '(A)') 'field = u' WRITE(40, '(A)') 'type = float' WRITE(40, '(A)') 'type = float' WRITE(40, '(A)') 'dependency = positions' WRITE(40, '(A)') 'dependency = positions = regular,regular,regular, &</pre>
74 75 76 77 78 79 80 81 82 83 84 85 86 87	<pre>WAILE(*, '(3(A,F5.4))') 'xn = ',xn,' / yn = ',yn,' / 2n = ',2n ! write OpenDX *.general header file OPEN(40, FILE=headfile, STATUS='REPLACE', ACTION='WRITE') WRITE(40, '(2A)') 'file = ', TRIM(outfile) WRITE(40, '(A,I4,A,I4,Y)') 'grid = ',Nx,' x ',Ny,' x ',Nz WRITE(40,'(A)') 'format = ascii' WRITE(40,'(A)') 'interleaving = record' WRITE(40,'(A)') 'majority = row' WRITE(40,'(A)') 'field = u' WRITE(40,'(A)') 'field = u' WRITE(40,'(A)') 'type = float' WRITE(40,'(A)') 'dependency = positions' WRITE(40,'(A)') 'dependency = positions = regular,regular,regular, &</pre>
74 75 76 77 78 79 80 81 82 83 84 85 86 87	<pre>WAILE(*, '(3(A,F6.4))') 'xn = ',xn,' / yn = ',yn,' / 2n = ',2n ! write OpenDX *.general header file OPEN(40, FILE=headfile, STATUS='REPLACE', ACTION='WRITE') WRITE(40, '(2A)') 'file = ', TRIM(outfile) WRITE(40, '(A,I4,A,I4,A,I4)') 'grid = ',Nx,' x ',Ny,' x ',Nz WRITE(40,'(A)') 'interleaving = record' WRITE(40,'(A)') 'interleaving = record' WRITE(40,'(A)') 'field = u' WRITE(40,'(A)') 'field = u' WRITE(40,'(A)') 'field = u' WRITE(40,'(A)') 'grid = float' WRITE(40,'(A)') 'dependency = positions' WRITE(40,'(A,5(F8.3,A),F8.3)') 'positions = regular,regular,regular, & ',xl,',',dx,',',yl,',',dy,',',zl,',',dz WRITE(40,'(A)') 'end' CLOSE(40) !</pre>
74 75 76 77 78 80 81 82 83 84 85 86 87	<pre>WATE(*, '(3(A,F6.4))') 'xn = ',xn,' / yn = ',yn,' / 2n = ',2n ! write OpenDX *.general header file OPEN(40, FILE=headfile, STATUS='REPLACE', ACTION='WRITE') WRITE(40, '(2A)') 'file = ', TRIM(outfile) WRITE(40, '(A,I4,A,I4,A,I4)') 'grid = ',Nx,' x ',Ny,' x ',Nz WRITE(40,'(A)') 'format = ascii' WRITE(40,'(A)') 'interleaving = record' WRITE(40,'(A)') 'interleaving = record' WRITE(40,'(A)') 'field = u' WRITE(40,'(A)') 'field = u' WRITE(40,'(A)') 'structure = scalar' WRITE(40,'(A)') 'dependency = positions' WRITE(40,'(A)') 'dependency = positions' WRITE(40,'(A,5(F8.3,A),F8.3)') 'positions = regular,regular,regular, & ',xl,',',dx,',',yl,',',dy,',',zl,',',dz WRITE(40,'(A)') 'end' CLOSE(40) !</pre>
74 75 76 77 80 81 82 83 84 85 86 87	<pre>WAILE(*, '(3(A,F5.4))') 'xn = ',xn,' / yn = ',yn,' / 2n = ',2n ! write OpenDX *.general header file OPEN(40, FILE=headfile, STATUS='REPLACE', ACTION='WRITE') WRITE(40, '(A)') 'file = ', TRIM(outfile) WRITE(40, '(A)') 'file = ', TRIM(outfile) WRITE(40, '(A)') 'interleaving = record' WRITE(40, '(A)') 'interleaving = record' WRITE(40, '(A)') 'majority = row' WRITE(40, '(A)') 'filed = u' WRITE(40, '(A)') 'filed = u' WRITE(40, '(A)') 'type = float' WRITE(40, '(A)') 'dependency = positions' WRITE(40, '(A)') 'dependency = positions = regular,regular,regular, & ',xl,',',dx,',yl,',',dy,',',zl,',',dz WRITE(40, '(A)') 'end' CLOSE(40) ! ! Allocate memory and read all data !</pre>
74 75 76 77 78 79 80 81 82 83 84 85 86 87 88	<pre>WAILE(*, '(3(A,F5.4))') 'xn = ',xn,' / yn = ',yn,' / 2n = ',2n ! write OpenDX *.general header file OPEN(40, FILE=headfile, STATUS='REPLACE', ACTION='WRITE') WRITE(40, '(A)') 'file = ', TRIM(outfile) WRITE(40, '(A)') 'file = ascii' WRITE(40, '(A)') 'interleaving = record' WRITE(40, '(A)') 'majority = row' WRITE(40, '(A)') 'majority = row' WRITE(40, '(A)') 'field = u' WRITE(40, '(A)') 'structure = scalar' WRITE(40, '(A)') 'dependency = positions' WRITE(40, '(A)') 'dependency = positions = regular,regular,regular, &</pre>
74 75 76 77 78 79 80 81 82 83 84 85 86 87 88 88 89	<pre>WAILE(*, '(3(A,F5.4))') 'xn = ',xn,' / yn = ',yn,' / 2n = ',2n ! write OpenDX *.general header file OPEN(40, FILE=headfile, STATUS='REPLACE', ACTION='WRITE') WRITE(40, '(2A)') 'file = ', TRIM(outfile) WRITE(40, '(A,I4,A,I4,Y)') 'grid = ',Nx,' x ',Ny,' x ',Nz WRITE(40,'(A)') 'format = ascii' WRITE(40,'(A)') 'interleaving = record' WRITE(40,'(A)') 'field = u' WRITE(40,'(A)') 'field = u' WRITE(40,'(A)') 'field = u' WRITE(40,'(A)') 'field = u' WRITE(40,'(A)') 'gructure = scalar' WRITE(40,'(A)') 'dependency = positions' WRITE(40,'(A,5(F8.3,A),F8.3)') 'positions = regular,regular,regular, & ',xl,',',dx,',',yl,',',dy,',',zl,',',dz WRITE(40,'(A)') 'end' CLOSE(40) !</pre>
74 75 76 77 78 80 81 82 83 84 85 86 87 88 89 90	<pre>WAILE(*, '(3(A,F5.4))') 'xn = ',xn,' / yn = ',yn,' / 2n = ',2n ' write OpenDX *.general header file OPEN(40, FILE=headfile, STATUS='REPLACE', ACTION='WRITE') WRITE(40, '(A)') 'file = ', TRIM(outfile) WRITE(40, '(A,I4,A,I4,)') 'grid = ',Nx,' x ',Ny,' x ',Nz WRITE(40,'(A)') 'interleaving = record' WRITE(40,'(A)') 'majority = row' WRITE(40,'(A)') 'field = u' WRITE(40,'(A)') 'type = float' WRITE(40,'(A)') 'type = float' WRITE(40,'(A)') 'dependency = positions' WRITE(40,'(A)') 'dependency = positions = regular,regular,regular, & ',xl,',',dx,',',yl,',',dy,',',zl,',',dz WRITE(40,'(A)') 'end' CLOSE(40) !</pre>
74 75 76 77 80 81 82 83 84 85 86 87 88 89 90	<pre>WAILE(*, '(3(A,F5.4))') 'Xn = ',Xn,' / yn = ',yn,' / 2n = ',2n ' write OpenDX *.general header file OPEN(40, FILE=headfile, STATUS='REPLACE', ACTION='WRITE') WRITE(40, '(A)') 'file = ', TRIM(outfile) WRITE(40, '(A)') 'file = ', TRIM(outfile) WRITE(40, '(A)') 'interleaving = record' WRITE(40, '(A)') 'majority = row' WRITE(40, '(A)') 'majority = row' WRITE(40, '(A)') 'field = u' WRITE(40, '(A)') 'field = u' WRITE(40, '(A)') 'dependency = positions' WRITE(40, '(A)') 'dependency = positions' WRITE(40, '(A)') 'dependency = positions = regular,regular,regular, & ',xl,',',dx,',',yl,',',dy,',',zl,',',dz WRITE(40, '(A)') 'end' CLOSE(40) '</pre>
74 75 76 77 78 79 80 81 82 83 84 85 86 87 88 89 90 91 92	<pre>WAILE(*, '(3(A,F5.4))') 'Xn = ',Xn,' / yn = ',yn,' / 2n = ',2n ! write OpenDX *.general header file OPEN(40, FILE=headfile, STATUS='REPLACE', ACTION='WRITE') WRITE(40, '(A)') 'file = ', TRIM(outfile) WRITE(40, '(A)') 'file = ascii' WRITE(40, '(A)') 'majority = record' WRITE(40, '(A)') 'majority = record' WRITE(40, '(A)') 'majority = row' WRITE(40, '(A)') 'field = u' WRITE(40, '(A)') 'structure = scalar' WRITE(40, '(A)') 'dependency = positions' WRITE(40, '(A)') 'dependency = positions = regular,regular,regular, & ',xl,',',dx,','yl,',',dy,','zl,',',dz WRITE(40, '(A)') 'end' CLOSE(40) !</pre>
74 75 76 77 78 79 80 81 82 83 84 85 86 87 85 86 87 89 90 91 92	<pre>WAILE(*, '(3(A,F5.4))') 'xh = ',xh,' / yh = ',yh,' / 2h = ',2h ! write OpenDX *.general header file OPEN(40, FILE=headfile, STATUS='REPLACE', ACTION='WRITE') WRITE(40, '(2A)') 'file = ', TRIM(outfile) WRITE(40, '(A,I4,A,I4,A,I4)') 'grid = ',Nx,' x ',Ny,' x ',Nz WRITE(40, '(A)') 'finterleaving = record' WRITE(40, '(A)') 'majority = row' WRITE(40, '(A)') 'field = u' WRITE(40, '(A)') 'field = u' WRITE(40, '(A)') 'type = float' WRITE(40, '(A)') 'dependency = positions' WRITE(40, '(A,5(F8.3,A),F8.3)') 'positions = regular,regular,regular, & ',xl, ',dx,',',yl,',',dy,',',zl,',',dz WRITE(40, '(A)') 'end' CLOSE(40) !</pre>
74 75 76 77 78 80 81 82 83 84 85 86 87 88 89 90 91 92	<pre>WATE(*, '(3(A,F5.4))') 'Xn = ',Xn,' / yn = ',yn,' / 2n = ',2n ' write OpenDX *.general header file OPEN(40, FILE=headfile, STATUS='REPLACE', ACTION='WRITE') WRITE(40, '(A)') 'file = ', TRIM(outfile) WRITE(40, '(A)') 'file = ', TRIM(outfile) WRITE(40, '(A)') 'interleaving = record' WRITE(40, '(A)') 'majority = row' WRITE(40, '(A)') 'majority = row' WRITE(40, '(A)') 'field = u' WRITE(40, '(A)') 'type = float' WRITE(40, '(A)') 'type = float' WRITE(40, '(A)') 'dependency = positions' WRITE(40, '(A)') 'dependency = positions = regular, regular, regular, &</pre>
74 75 76 77 78 80 81 82 83 84 85 86 87 88 89 90 91 92 92	<pre>WAILE(*, '(3(A,F5.4))') 'Xn = ',Xn,' / yn = ',yn,' / 2n = ',2n ! write OpenDX *.general header file OPEN(40, FILE=headfile, STATUS='REPLACE', ACTION='WRITE') WRITE(40, '(A)') 'file = ', TRIM(outfile) WRITE(40, '(A)') 'file = ascii' WRITE(40, '(A)') 'interleaving = record' WRITE(40, '(A)') 'majority = row' WRITE(40, '(A)') 'majority = row' WRITE(40, '(A)') 'field = u' WRITE(40, '(A)') 'dependency = positions' WRITE(40, '(A)') 'dependency = positions' WRITE(40, '(A)') 'dependency = positions = regular,regular,regular, & ',xl,',',dx,',',yl,',',dy,',',zl,',',dz WRITE(40, '(A)') 'end' CLOSE(40) ! ! Allocate arrays for positions and masses ALLOCATE(xp(3,Np), mass(Np), STAT=istat) if(istat .NE. 0) then WRITE(*, '(A)') 'Error allocating memory for data.' goto 9999 end if ! allocate arrays for particle sorting and chaining mesh ALLOCATE(MPC(MPX NPY NPZ) 11(Mp) STAT=istat)</pre>
74 75 76 77 78 79 80 81 82 83 84 85 86 87 88 89 90 91 92 93 94	<pre>WATE(*, '(S(A,F6.4))') 'xh = ',xh,' / yh = ',yh,' / zh = ',zh ! write OpenDX *.general header file OPEN(40, FILE=headfile, STATUS='REPLACE', ACTION='WRITE') WRITE(40,'(2A)') 'file = ', TRIM(outfile) WRITE(40,'(A)') 'finerleaving = record' WRITE(40,'(A)') 'interleaving = record' WRITE(40,'(A)') 'finefleaving = record' WRITE(40,'(A)') 'finefleaving = record' WRITE(40,'(A)') 'finefleaving = resolar' WRITE(40,'(A)') 'finefleaving = resolar' WRITE(40,'(A)') 'finefleaving = regular, 'w'''''''''''''''''''''''''''''''''''</pre>
74 75 76 77 78 80 81 82 83 84 85 86 87 88 89 90 91 92 93 94 95	<pre>WATE(*, '(S(A,F6.4))') 'xh = ',xh,' / yh = ',yh,' / zh = ',zh ! write OpenDX *.general header file OPEN(40, FILE=headfile, STATUS='REPLACE', ACTION='WRITE') WRITE(40,'(A)') 'file = ', TRIM(outfile) WRITE(40,'(A)') 'format = ascii' WRITE(40,'(A)') 'interleaving = record' WRITE(40,'(A)') 'interleaving = record' WRITE(40,'(A)') 'interleaving = record' WRITE(40,'(A)') 'interleaving = record' WRITE(40,'(A)') 'structure = scalar' WRITE(40,'(A)') 'structure = scalar' WRITE(40,'(A)') 'type = float' WRITE(40,'(A)') 'type = float' WRITE(40,'(A,5(F8.3,A),F8.3)') 'positions = regular,regular,regular, & ',xl,',',dx,',',yl,',',dy,',',zl,',',dz WRITE(40,'(A)') 'end' CLOSE(40) !</pre>
74 75 76 77 78 80 81 82 83 84 85 86 87 88 89 90 91 92 93 94 95	<pre>WATE(*, '(S(A,FG.4))') 'Xh = ',Xh,' / Yh = ',Yh,' / Zh = ',Zh ! write OpenDX *.general header file OPEN(40, FILE=headfile, STATUS='REPLACE', ACTION='WRITE') WRITE(40,'(2A)') 'file = ', TRIM(outfile) WRITE(40,'(A)') 'finerat = ascii' WRITE(40,'(A)') 'interleaving = record' WRITE(40,'(A)') 'interleaving = record' WRITE(40,'(A)') 'filed = u' WRITE(40,'(A)') 'structure = scalar' WRITE(40,'(A)') 'structure = scalar' WRITE(40,'(A)') 'type = float' WRITE(40,'(A)') 'type = float' WRITE(40,'(A)') 'dependency = positions' WRITE(40,'(A)') 'dependency = positions = regular,regular,regular, & ',xl,',',dx,',',yl,',',dy,',',zl,',',dz WRITE(40,'(A)') 'end' CLOSE(40) !</pre>
74 75 76 77 78 80 81 82 83 84 85 86 87 88 89 90 91 92 93 94 95 96 77	<pre>WRITE(*, '(S(A,F6.4))') 'M = ',M,' / yn = ',yn,' / 2n = ',2n ! write OpenDX *.general header file OPEN(40, FILE-headfile, STATUS='REPLACE', ACTION='WRITE') WRITE(40,'(2A)') 'file = ', TRIM(outfile) WRITE(40,'(A), / 14,A,I4,A,I4)') 'grid = ',Nx,' x ',Ny,' x ',Nz WRITE(40,'(A)') 'interleaving = record' WRITE(40,'(A)') 'interleaving = record' ! Allocate memory and read all data !</pre>
74 75 76 77 78 79 80 81 82 83 84 85 86 87 88 89 90 91 92 93 94 95 96 97	<pre>WRITE(*,'(S(A,F6.4))')''An = ',An,' / yn = ',yn,' / 2n = ',2n ! write OpenDX *.general header file OPEN(40, FILE-headfile, STATUS='REPLACE', ACTION='WRITE') WRITE(40,'(2A)') 'file = ', TRIM(outfile) WRITE(40,'(A),14,A,14,A,14)') 'grid = ',Nx,' x ',Ny,' x ',Nz WRITE(40,'(A)') 'majority = record' WRITE(40,'(A)') 'majority = row' WRITE(40,'(A)') 'structure = scalar' WRITE(40,'(A)') 'type = float' WRITE(40,'(A)') 'dependency = positions' WRITE(40,'(A)') 'dependency = positions = regular,regular,regular, & ',x1,',',dx,',y1,',',dy,',',21,',',dz WRITE(40,'(A)') 'end' CLOSE(40) !</pre>
74 75 76 77 78 80 81 82 83 84 85 86 87 88 89 90 91 92 93 94 95 96 97	<pre>wkllE(*, (3(A,F0.4))') 'Xh = ',Xh,' / Yh = ',Yh,' / 2h = ',Zh ! write OpenDX *.general header file OPEN(40, FILE=headfile, STATUS='REPLACE', ACTION='WRITE') WRITE(40,'(A)') 'file = ', TRIM(outfile) WRITE(40,'(A)') 'interleaving = record' WRITE(40,'(A)') 'interleaving = record' WRITE(40,'(A)') 'interleaving = record' WRITE(40,'(A)') 'itper = float' WRITE(40,'(A)') 'structure = scalar' WRITE(40,'(A)') 'type = float' WRITE(40,'(A)') 'dependency = positions' WRITE(40,'(A)') 'dependency = positions = regular,regular,regular, & ',Xl,',',Xr,',',yl,',',dy,',',zl,',',dz WRITE(40,'(A)') WRITE(40,'(A)') WRITE(40,'(A)') 'end' CLOSE(40) !</pre>
74 75 76 77 78 80 81 82 83 84 85 86 87 88 89 90 91 92 93 94 95 96 97	<pre>wkllE(*, (3(A,F0.4))') 'Xh = ',Xh,' / Yh = ',Yh,' / 2h = ',Zh ! write OpenDX *.general header file OPEN(40, FILE=headfile, STATUS='REPLACE', ACTION='WRITE') WRITE(40,'(2A)') 'file = ', TRIM(outfile) WRITE(40,'(A)') 'format = ascii' WRITE(40,'(A)') 'interleaving = record' WRITE(40,'(A)') 'majority = row' WRITE(40,'(A)') 'field = u' WRITE(40,'(A)') 'field = u' WRITE(40,'(A)') 'type = float' WRITE(40,'(A)') 'dependency = positions' WRITE(40,'(A)') 'dependency = positions = regular,regular,regular, & ',xl,',',dx,',yl,',',dy,',',zl,',',dz WRITE(40,'(A)') WRITE(40,'(A)') end' CLOSE(40) !</pre>
74 75 76 77 78 80 81 82 83 84 85 86 87 88 89 90 91 92 93 94 95 96 97 98	<pre>wkllE(*, (3(A,F0.4))') 'Xh = ',Xh,' / Yh = ',Yh,' / 2h = ',Zh ! write OpenDX *.general header file OPEN(40, FILE=headfile, STATUS='REFLACE', ACTION='WRITE') WRITE(40,'(A)') 'file=' ', TRIM(outfile) WRITE(40,'(A)') 'format = ascil' WRITE(40,'(A)') 'interleaving = record' WRITE(40,'(A)') 'interleaving = record' WRITE(40,'(A)') 'itple = n' WRITE(40,'(A)') 'structure = scalar' WRITE(40,'(A)') 'structure = scalar' ! Allocate memory and read all data !</pre>

```
! second pass: read points
100
        WRITE(*,'(A)') 'Reading data points ...'
101
        OPEN(40, FILE=infile, STATUS='OLD', ACTION='READ')
102
        do j=1,Np
103
           READ(40,'(4E16.8)') xp(1:3,j), mass(j)
104
         end do
105
        CLOSE(40)
         ! Set up chaining mesh and sort the particles into boxes
         !----
106
        WRITE(*,'(A)') 'Sorting particles in linked lists ...'
         ! determine box sizes
107
        dbx = (xh-xl)/real(NBX)
        dby = (yh-yl)/real(NBY)
108
        dbz = (zh-zl)/real(NBZ)
109
         ! loop over all particles
        do j=1,Np
110
            ! determine index of box we are in
111
           idx = ceiling((xp(1,j)-xl)/dbx)
112
           if(idx .EQ. 0) idx = 1
           if(idx .EQ. NBX+1) idx = NBX
113
           idy = ceiling((xp(2,j)-yl)/dby)
114
           if (idy .EQ. 0) idy = 1
if (idy .EQ. NBY+1) idy = NBY
115
116
           idz = ceiling((xp(3,j)-zl)/dbz)
117
           if(idz . EQ. 0) idz = 1
118
           if(idz .EQ. NBZ+1) idz = NBZ
119
           ! add particle to linked list of this box
           ll(j) = HOC(idx,idy,idz)
120
           HOC(idx,idy,idz) = j
121
122
        end do
         1-----
         ! Write regular grid output
         1-----
                   _____
         ! write output file
        WRITE(*,'(A)') 'Writing regular grid output file ...'
123
        OPEN(50, FILE=outfile, STATUS='REPLACE', ACTION='WRITE')
124
125
        Nw = 0
        ! threshold for nearest grid point
126
        t = (0.5*minval((/dx,dy,dz/)))**2
      ! t = 0.1
        ! loop over all grid points
127
        do j=1,Nx
128
           do k=1,Ny
129
              do l=1,Nz
                 ! current position
                 xold(1) = xl+(j-1)*dx
xold(2) = yl+(k-1)*dy
130
131
                 xold(3) = zl+(l-1)*dz
132
                 ! determine box index
133
                 idx = ceiling((xold(1)-xl)/dbx)
134
                 if(idx . EQ. 0) idx = 1
135
                 if(idx .EQ. NBX+1) idx = NBX
136
                 idy = ceiling((xold(2)-yl)/dby)
137
                 if(idy .EQ. 0) idy = 1
138
                 if(idy .EQ. NBY+1) idy = NBY
                 idz = ceiling((xold(3)-zl)/dbz)
139
140
                 if(idz . EQ. 0) idz = 1
                 if(idz .EQ. NBZ+1) idz = NBZ
141
                  ! search for this grid point in the data
142
                 found = .FALSE.
                 ! traverse linked list for all particles in the same box
143
                 m = HOC(idx,idy,idz)
144
                 do while (m .NE. O .AND. found .EQ. .FALSE.)
                    ! if particle is close enough to grid point, take it and
                     ! exit loop
145
                    if(sum((xp(1:3,m)-xold(1:3))**2) .LT. t) then
146
                       WRITE(50,*) mass(m)
147
                       Nw = Nw + 1
                       found = .TRUE.
148
149
                    end if
                    m = ll(m)
150
                 end do
151
```

```
! if not found, output zero
152
                  if(.NOT. found) WRITE(50,*) 0.0_MK
               end do
153
154
            end do
155
         end do
156
         CLOSE(50)
         ! Terminate and deallocate memory
         ! terminal output
         WRITE(*, '(2(A, I8))') 'Particles read: ',Np,' / Particles written: ',Nw
157
         WRITE(*,'(A)')
158
         ! deallocate memory
         DEALLOCATE(HOC, 11)
159
         DEALLOCATE(xp, mass)
160
161
      end do
      9999 CONTINUE
162
         if(ALLOCATED(HOC)) DEALLOCATE(HOC)
163
164
         if(ALLOCATED(11)) DEALLOCATE(11)
         if(ALLOCATED(xp)) DEALLOCATE(xp)
165
         if(ALLOCATED(mass)) DEALLOCATE(mass)
166
167
      END program res2dx
```

C.4 Gradient descent fitting

C.4.1 General description and usage notes

The program code for the gradient descent algorithm described in section 11.3.1 and used in section 11.3.4 and chapter 12 to fit the various FRAP data models to both simulation and experimental data is given next. The program does a nonlinear least squares fit of a model having 1 or 2 parameters called α and β . It uses the slow but robust (cf. the considerations in section 11.3.3) method of gradient descent and consists of a single source file. The following input parameters are specified in statements 5 to 16:

Parameter	Meaning
TOL (TOL)	Tolerance for termination criterion
a	Size (extension) of the bleached box in x -direction
b	Size (extension) of the bleached box in y -direction
MODEL	Model to be fitted to the data
VIZ	Create visualization output of error function ?
Nalpha	Number of visualization points in α -direction
Nbeta	Number of visualization points in β -direction
alpha0	Lower α bound for visualization
beta0	Lower β bound for visualization
alphamax	Upper α bound for visualization
betamax	Upper β bound for visualization
POSTOUT	Create PostScript figures or not

Table C.3: Input parameters for nlfit.f90

The parameter MODEL specifies which model to use. It can assume the following values: 1 (exponential model given by equation 11.1), 2 (power law model given by equation 11.3), 3 (second order physical model given by equation 11.4) or 4 (empirical model given by equation 11.2). The parameter VIZ can be either .TRUE. or .FALSE.. It specifies whether the error function should be probed on a regular

cartesian grid over the parameter space in order to visualize it prior to the gradient descent. This is for example useful in finding initial parameter values for the optimization or to investigate shape and properties of the error function.

The initial parameter values for α and β need to be specified in statements 35 and 36 in variables alpha and beta. The corresponding learning rates η_{α} and η_{β} are to be defined in statements 37 and 38 in etaa and etab.

In addition to these parameters, the correct level of the theoretical steady-state FRAP value F_{∞} for $t \to \infty$ has to be set in y0. The values for all the samples considered in this project are given around statement 34 and one only needs to comment out the right one.

The files containing the data one wishes to fit the model to are passed to the program as command line arguments upon invocation:

```
$ nlfit <file1.frap> [file2.frap [file3.frap ...]]
```

Any number of files can be processed at the same call, shell wildcards are allowed. The input files contain FRAP data points using the following syntax on each line:

```
<time> <FRAP value>
```

The output files that are generated for each such input file are **<infile>.grad**, **<infile>.eta** and **<infile>.mac**. The .grad file contains the model fitting error E as defined by equation 11.6 and the error gradients $\frac{\partial E}{\partial \alpha}$ and $\frac{\partial E}{\partial \beta}$ for each iteration step as:

```
<error E> <dEdalpha> <dEdbeta>
```

This information can be used to create convergence plots such as the ones in figures 11.6, 12.3 or 12.5. In addition, these values are also printed to the standard output. The .eta file contains the learning rates η_{α} and η_{β} of each iteration step as:

```
<etaa> <etab>
```

The .mac file finally contains the gnuplot macro to plot the data and the best model fit in a picture like figure 11.7 or 12.4. If POSTOUT is set to .TRUE., gnuplot is instructed to generate PostScript figure files, otherwise the plot is output to a graphics window.

If the parameter VIZ is set to .TRUE., an additional output file called E.out is written that contains the values of the absolute and relative RMS matching errors on a Nalpha×Nbeta grid for α between alpha0 and alphamax and β between beta0 and betamax:

```
for i=1,Nbeta
  for j=1,Nalpha
      <alpha> <beta> <E> <Erel>
      end for
end for
```

This data can then be used to create contour plots of the "error landscape" such as the ones in figures 11.3 to 11.5.

C.4.2 Source code listing

The program source code for the gradient descent algorithm and all the visualization is given below.

C.4.2.1 Curve fitting using gradient descent (nlfit.f90)

```
1 program nlfit
```

```
2 IMPLICIT NONE
```

```
!NLFIT Does a non-linear model fit using a gradient decent technique.
        nlfit reads simulated FRAP data from the file(s) specified on the
        \ensuremath{\mathsf{command}} line and estimates the model parameters for various models
        to best fit the simulation results in a least squares sense.
        See also PSE3D
        todo:
     DIPLOMA THESIS WS01/02 ICOS
                                                              ETH-ZUERICH
     L
                _____
                  PROTEIN DIFFUSION INSIDE THE ENDOPLASMIC RETICULUM
     ! Declaration of external functions
     ! get number of command line arguments
3
     INTEGER, EXTERNAL
                                         :: iargc
                 ! Declaration of parameters
     ! precision
4
     INTEGER, PARAMETER
                                        :: MK = KIND(1.0D0)
     ! convergence tolerance (for termination criterion)
                                       :: TOL = 1e-6_MK
     REAL(MK), PARAMETER
5
     ! x and y dimensions of bleached box
     REAL(MK), PARAMETER
                                        :: a = 50.0
6
     REAL(MK), PARAMETER
                                        :: b = 50.0
     ! Which model (1=exponential, 2=power law, 3=second order physical,
     ! 4=empirical as in SA,p.14)
8
     INTEGER, PARAMETER
                                        :: MODEL = 3
     ! Do area bombing to visualize function?
     LOGICAL, PARAMETER
                                         :: VIZ = .FALSE.
9
     ! Number of alpha and beta grid points for visualization
     INTEGER, PARAMETER
                                        :: Nalpha = 500
10
     INTEGER, PARAMETER
                                        :: Nbeta = 500
11
     ! lower boundary for alpha and beta for visualization
12
     REAL(MK), PARAMETER
                                        :: alpha0 = 0.1
     REAL(MK), PARAMETER
                                         :: beta0 = 0.01
13
     ! upper boundary for alpha and beta for visualization
     REAL(MK), PARAMETER
                                        :: alphamax = 5.0
14
     REAL(MK), PARAMETER
                                        :: betamax = 1.00
15
     ! Produce postscript gnuplot output?
     LOGICAL, PARAMETER
16
                                       :: POSTOUT = .FALSE.
     ! Declaration of local variables
     1 ----
     ! loop counters
17
     INTEGER
                                        :: i, j, k, l
     ! name of the program
18
     CHARACTER(LEN=80)
                                        :: prgname
     ! name of the input and output files
19
     CHARACTER(LEN=80)
                                         :: infile, outfile, gradfile, etafile
     ! number of command line arguments (input files)
20
     INTEGER
                                        :: nargc
     ! simulated FRAP data from file
21
     REAL(MK), DIMENSION(:), ALLOCATABLE
                                        :: frap
     ! number of data points
22
     INTEGER
                                         :: N
     ! current model value
23
     REAL(MK)
                                         :: у
     ! error status
```

```
24
     INTEGER
                                         :: istat
     ! model parameters to be fitted
25
     REAL(MK)
                                         :: alpha, beta, alphaold, betaold
     ! gradients of the error function
26
     REAL(MK)
                                         :: dEda, dEdb, dEdaold, dEdbold
     ! time array
27
     REAL(MK), DIMENSION(:), ALLOCATABLE
                                         :: t
     ! learning rates for alpha and beta, respectively
28
     REAL(MK)
                                         :: etaa, etab
     ! steady-state value of frap curve
29
     REAL(MK)
                                         :: y0
     ! value of the error function
30
     REAL(MK)
                                         :: E, Erel
     ! sign of the dot product of subsequent gradients
31
     INTEGER
                                         :: s
     ! intermediates
32
     REAL(MK)
                                         :: den, dyda, dydb
     ! alpha and beta step sizes for visualization
33
     REAL(MK)
                                         :: deltaa, deltab
                                   _____
                                                      _____
     ! steady state values
                       _____
     ! y0 = 0.9375
                                    ! for Box test case
     ! y0 = 481958.0_MK/493850.0_MK
                                   ! for bip2
     ! y0 = 475859.0_MK/488054.0_MK
                                   ! for clx
     ! y0 = 503182.0_MK/509578.0_MK
                                   ! for erp57
                                   ! for erp572
     y0 = 538092.0_MK/545859.0_MK
                                   ! for erp573_1
     y0 = 463516.0_MK/473466.0_MK
     y0 = 482469.0_MK/495373.0_MK
                                   ! for erp573_2
     ! y0 = 319538.0_MK/335480.0_MK
                                   ! for erp573_3
     ! y0 = 469745.0_MK/481922.0_MK
                                   ! for erp574_1
     ! y0 = 481486.0_MK/504770.0_MK
                                   ! for erp574_2
     ! y0 = 539400.0_MK/545859.0_MK
                                   ! for erp5722
     ! y0 = 538412.0_MK/545859.0_MK
                                   ! for erp5723
     ! y0 = 528973.0_MK/545859.0_MK
                                  ! for erp5724
      y0 = 477774.0_MK/480214.0_MK
34
                                  ! for 8s9
     ! y0 = 440104.0_MK/480214.0_MK
                                  ! for 8.2
     ! y0 = 0.98_MK
     I-----
     ! start values
     1------
35
     alpha = 1.0_MK
36
     beta = 0.05_MK
37
     etaa = 1e-4_MK
38
     etab = 5e-5 MK
     1-----
     ! Process command line arguments
     1---
     ! read program name from command line
     CALL getarg(0,prgname)
39
     ! get number of input files
40
     nargc = iargc()
41
     if(nargc .LT. 1) then
42
       print*,'Missing input file name(s). Usage: ',TRIM(prgname), &
          ' inputfile [inputfile ...]'
43
       goto 9999
44
     end if
     ! Delete old result file
                          _____
     OPEN(40, FILE='All_Params.out', ACTION='WRITE')
45
46
     CLOSE(40, STATUS='DELETE')
     ! Loop over all input files
     !---
47
     do i=1,nargc
                                  _____
```

! Scan file and read its contents

C.4. GRADIENT DESCENT FITTING

! get input file name 48 CALL getarg(i, infile) ! first pass: count points 49 OPEN(40, FILE=infile, STATUS='OLD', ACTION='READ') 50 N = O51 do 52 READ(40,*,END=110) y, y 53 N = N + 154 end do 55 110 CLOSE(40) ! allocate memory 56 ALLOCATE(frap(N), t(N), STAT=istat) 57 if(istat .NE. 0) then 58 WRITE(*,'(A)') 'Error allocating memory for frap data.' goto 9999 59 60 end if ! second pass: read data 61 OPEN(40, FILE=infile, STATUS='OLD', ACTION='READ') 62 do j=1,N READ(40,*) t(j), frap(j) 63 64 end do 65 CLOSE(40) ! Visualization of error function if(VIZ) then 66 67 OPEN(40, FILE='E.out', STATUS='REPLACE', ACTION='WRITE') 68 deltab = (betamax-beta0)/real(Nbeta-1,MK) deltaa = (alphamax-alpha0)/real(Nalpha-1,MK) 69 do j=1,Nbeta 70 71 beta = beta0+(j-1)*deltab 72 do k=1,Nalpha alpha = alpha0+(k-1)*deltaa E = 0.0_MK 73 74 $Erel = 0.0_MK$ 75 do l=1,N 76 ! current model value 77 if(MODEL .EQ. 2) then ! POWER MODEL y = y0*(1.0_MK-(beta*t(1)+1)**(-alpha)) 78 elseif(MODEL .EQ. 1) then ! EXPONENTIAL MODEL 79 80 $y = y0*(1.0_MK-exp(-alpha*t(1)))$ elseif(MODEL .EQ. 3) then 81 ! SECOND ORDER PHYSICAL MODEL den = (a*b+2.0_MK*(a+b)*alpha*t(l)**beta & 82 +4.0_MK*alpha**2*t(1)**(2.0_MK*beta)) 83 $y = y0*(1.0_MK-a*b/den)$ 84 elseif(MODEL .EQ. 4) then ! EMPIRICAL MODEL FROM SA, P.14 85 den = (t(1)/beta)**alpha 86 $y = y0*den/(1.0_MK+den)$ else 87 88 print*,'Unknown model parameter specified !' goto 9999 89 90 end if ! error function 91 E = E + (frap(1)-y) **292 Erel = Erel + ((frap(1)-y)/(y0-frap(1)))**2 93 end do 94 E = sqrt(E/real(N)) 95 Erel = sqrt(Erel/real(N)) 96 WRITE(40, '(4E12.4)') alpha, beta, E, Erel 97 end do end do 98 CLOSE(40) 99 100 end if ! open some output files 101 WRITE(gradfile, '(2A)') TRIM(infile), '.grad' OPEN(40, FILE=gradfile, STATUS='REPLACE', ACTION='WRITE') 102 WRITE(etafile,'(2A)') TRIM(infile),'.eta' 103 OPEN(50, FILE=etafile, STATUS='REPLACE', ACTION='WRITE') 104 _____

1-----

! Gradient decent

!-----

105	k = 0
	! iterate for alpha and beta
106	do
107	sum the gradients over all time points
107	if(MODEL GT 1) dEdbold = dEdb
109	dEda = 0.0 MK
110	dEdb = 0.0 MK
111	E = 0.0 MK
112	do j=2,ℕ
	! current model value
113	if(MODEL .EQ. 1) then
11/	! EXPUNENTIAL MUDEL
114	$y = y_0 + (1.0_{\text{rk}} - exp(-a_1p_1a + t(j)))$
116	elseif(MODEL .EQ. 2) then
	! POWER MODEL
117	y = y0*(1.0_MK-(beta*t(j)+1)**(-alpha))
118	dEda = dEda + (y-frap(j))*(y0-y)*log(beta*t(j)+1)
119	dEdb = dEdb + (y-frap(j))*(y0-y)*((alpha*t(j))/(beta*t(j)+1))
120	elseif(MODEL .EQ. 3) then
101	! SECUND URDER PHYSICAL MUDEL
121	$dem = (a*D+2.0_mA*(a+D)*aipia*i(j)**Deta & \\ +4 \ 0 \ MK*aipha**2*+(i)**(2 \ 0 \ MK*heta))$
122	v = v0*(1.0 MK-a*b/den)
123	dyda = y0*a*b*den**(-2.0_MK)*(2.0_MK*(a+b)*t(j)**beta &
	+8.0_MK*alpha*t(j)**(2.0_MK*beta))
124	dydb = y0*a*b*den**(-2.0_MK)*(2.0_MK*(a+b) &
	<pre>*alpha*t(j)**beta*log(t(j))+8.0_MK*alpha**2*t(j) &</pre>
405	**(2.0_MK*beta)*log(t(j)))
125	dEda = dEda + (y - Irap(j)) * dyda $dEdb = dEdb + (y - frap(j)) * dyda$
127	elseif(MODEL .EQ. 4) then
	! EMPIRICAL MODEL FROM SA, P.14
128	den = $(t(j)/beta)**alpha$
129	$y = y0*den/(1.0_MK+den)$
130	dyda = y0*(den*log(t(j)/beta)*(1.0_MK+den)**(-1.0_MK) - &
131	$den*(1.0_mh+den)**(-2.0_mh)*den*10g(t(j)/beta))$ $dvdh = v0*(-a)nha/(heta*heta)*(t(i)/heta)**(a)nha-1.0 MK) $
101	*(1.0 MK+den)**(-1.0 MK)+den*(1.0 MK+den)**(-2.0 MK) &
	<pre>*alpha/(beta*beta)*(t(j)/beta)**(alpha-1.0_MK))</pre>
132	dEda = dEda + (y-frap(j))*dyda
133	dEdb = dEdb + (y-frap(j))*dydb
134	print* 'Unknown model specified '
136	goto 9999
137	end if
138	E = E + (frap(j)-y)**2
139	end do
140	E = 0.5 MK * E
141	WRITE(40,*) E, dEda, dEdb
142	WRITE(50,*) etaa, etab WRITE(* *) E dEda dEda
1 10	
	! update model parameters
144	alphaold = alpha
145	if(MODEL .GT. 1) betaold = beta
146	alpha = alphaold - etaa*dEda if(MODEL (T. 1) bata = bataold = atab#dEdb
148	k = k + 1
1 10	
	! terminateion criterion
149	if(abs(dEda) .LE. TOL .AND. abs(dEdb) .LE. TOL) then
150	goto 200
151	end if
152	
154	CLOSE(40) CLOSE(50)
	!
	! Output results to terminal and result file
	!
155	200 WRITE(* '(34 T8 4)') 'Innut file' ' TRIM(infile) ' containe ' N %
100	' data points'
156	WRITE(*,'(A,I8,A)') 'Solution converged after ',k,' Iterations.'
157	WRITE(*,'(A,E16.8)') 'Best parameters found: alpha = ', alpha

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```
158
           WRITE(*,'(A,E16.8)') '
                                                               beta = '. beta
           OPEN(40, FILE='All_Params.out', POSITION='APPEND', ACTION='WRITE')
159
160
           WRITE(40,'(I4,3E16.8,I7)') i, alpha, beta, E, k
161
           CLOSE(40)
162
           WRITE(*,'(A)')
163
          DEALLOCATE(frap, t)
           ! Write gnuplot comparison macro
164
          WRITE(outfile,'(2A)') TRIM(infile), '.mac'
165
          OPEN(40, FILE=outfile, STATUS='REPLACE', ACTION='WRITE')
166
          WRITE(40, '(A)') 'set nokey'
          if(POSTOUT) then
167
              WRITE(40, '(A)') 'set term postscript eps color 14'
168
              WRITE(40, '(3A)') 'set output ''', TRIM(infile), '.eps'''
169
170
          end if
          if(MODEL .EQ. 2) then
171
              ! POWER MODEL
172
              WRITE(40, '(3A, F8.5, A, F8.5, A)') 'p ''', TRIM(infile), ''' w p, ', y0, &
                  '*(1.0-(',beta,'*x+1.0)\\'
173
              WRITE(40, '(A, F8.5, A)') '**(-', alpha, ')) w l'
          elseif(MODEL .EQ. 3) then
174
              ! SECOND ORDER PHYSICAL MODEL
             WRITE(40,'(34,F8.5,A)') 'p ''',TRIM(infile),''' w p, ',y0,'*(1.0-( \\'
WRITE(40,'(6(F8.5,A)') a,'*',b,')/(',a,'*',b,'+2.0*(',a,'+',b,')* \\'
175
176
              WRITE(40, '(3(F8.5, A)') alpha, '*x**', beta, '+4.0*', alpha, '**2*x**( \\'
177
              WRITE(40, '(F8.5, A)') 2.0*beta, '))) w 1'
178
          elseif(MODEL .EQ. 4) then
179
              ! EMPIRICAL MODEL FROM SA, P.14
             WRITE(40,'(3A,F8.5,A)') 'p ''',TRIM(infile),''' w p, ',y0,'*((x/ \\'
WRITE(40,'(4(F8.5,A))') beta,')**',alpha,')/(1.0+(x/',beta,')** &
180
181
          ',alpha,') w l'
elseif(MODEL .EQ. 1) then
182
              ! EXPONENTIAL MODEL
              WRITE(40,'(3A,F8.5,A)') 'p ''',TRIM(infile),''' w p, ',y0,'*(1.0- \\'
183
              WRITE(40, '(A,F8.5,A)') 'exp(',-alpha,'*x)) w l'
184
185
          else
             print*,'Unknown model specified.'
186
187
              goto 9999
188
          end if
          WRITE(40, '(A)') 'pause -1'
189
190
          CLOSE(40)
191
       end do
192
       9999 CONTINUE
193
      END program nlfit
```

C.5 Geometry preprocessor

C.5.1 Code structure and calling tree

The geometry preprocessor code implements the algorithms described in chapters 3 and 5 as well as section 7.2.5. It is parallelized using MPI and uses C preprocessor directives to encapsulate both the serial version and the parallel version of the code in the same source files. The parallel version is built if the code is preprocessed with <u>_MPI_</u> defined:

\$ cpp -D__MPI__ -P init_part.f90 __init_part.f90

to build the serial version, no defines are needed:

```
$ cpp -P init_part.f90 __init_part.f90
```

__init_part.f90 then is the program to be compiled. The code consists of the following modules at this stage of development:

Name	Function
init_part.f90	Main program
globals.f90	Global definitions and variables
Defaults.f90	Set default parameter values
readinput.f90	Reads all input files
ReadParams.f90	Read parameter input file
readiv.f90	Read OpenInventor 3D files
chktriang.f90	Check triangulation for validity
intersect.f90	Intersect a line with a triangle
SortT.f90	Sort triangles to bin lists and cell lists
AllocateLL.f90	Bin and cell list memory management
<pre>point_in_domain.f90</pre>	Check if a point is inside the computational domain
Voxelize.f90	Convert triangulation to voxel representation
BCdim.f90	Measure the box counting dimension
UpperCase.f90	Convert a string to upper case letters
Util.f90	Dynamic list resizing

Table C.4: Modules of the preprocessor code

The code's structure given by its calling tree is:



where "Util.reallocate" means a call to function reallocate that is contained in module Util and "(BCdim)" means that function BCdim can be called if its services are wanted but will generally be commented out (see lines between statement 289 and 290).

C.5.2 Input files and parameters

The code reads two input files: preproc.in, which contains all parameters and settings, and a geometry input file containing the triangulated surface description. The parameter input file preproc.in uses the same syntax as the one for the PSE code. It contains several lines like:

<keyword> = <value>

Additionally, it can contain any number of blank lines and comment lines beginning with a hash symbol (#). Those lines are ignored by the program. Table C.5 lists all supported keywords, their purpose and default values. If a certain keyword is not explicitly specified in the parameter file, its default value will be assumed. The keyword specifications can occur in any order in the parameter file

Keyword	Meaning	Default value
surfacefile	Name of geometry input file to read	box.surf
restartfile	Name of particle output file to write	restart.dat
dx	Grid spacings in x, y, z	0.1,0.1,0.1
bleachbox	$x_{min}, y_{min}, x_{max}$ and y_{max} of bleached box	2.0, 2.0, 3.0, 3.0
cutoff	PSE cutoff radius (for cell lists)	1.0
boundarycondition	BC type: 0=none, 1=Dirichlet, 2=Neumann	0
gnuout	Produce gnuplot output ?	.FALSE.
checktopology	Check triangulation for validity ?	.FALSE.
tolerance	Global geometry tolerance	10^{-10}
NbinsY	Number of bins in y -direction (for bin lists)	1
NbinsZ	Number of bins in z -direction (for bin lists)	1

Table C.5: Keywords for preprocessor parameter file

For keywords that take more than a single number, a comma separated list is to be given. For **bleachbox**, this list contains the *x*-coordinate of the lower left corner, the *y*-coordinate of the lower left corner, the *x*-coordinate of the upper right corner, the *y*-coordinate of the upper right corner of the bleached box. As an example, the file for the erp572 simulation run of section 9.1 is given below.

```
# Parameter file for preprocessing tool
# _____
# Name of the surface description file to read in:
surfacefile
                = ../3div/erp572.iv
# Name of PSE3D restart file to be written
restartfile
                = restart.dat
# Lattice spacing in x, y and z direction respectively:
                = 2.0, 2.0, 0.4
dx
# Bleached box geometry: lower and upper boundaries in x and y
bleachbox
                = 250.0, 125.0, 300.0, 175.0
# Cutoff radius as used in PSE (default: 10*eps)
cutoff
                 = 22.23
# Type of B.C. (0: none, 1: zero Dirichlet, 2: zero Neumann)
boundary condition = 2
```

Produce additional ASCII output for gnuplot?

gnuout = .false.
Run topology check ?
checktopology = .true.
Geometry tolerance (at least 10 times smaller than smallest geometry struct.)
tolerance = 1e-10
number of bins in y and z
NbinsY = 100
NbinsZ = 100

The geometry input file specified using the surfacefile keyword in the parameter input file can either be a triangulation .surf file as described in [Sbalzarini (2001)] or an OpenInventor 3D file in ASCII format (see [SGI (1992a)] and [SGI (1992b)]). The code recognizes the file type by the first line it contains. If it is

#ASCII triangulated surface description

the file is considered to be in .surf-format according to the specifications in appendix C.2.5 of [Sbalzarini (2001)]. If the first line is

#Inventor V2.1 ascii

the file is treated as an OpenInventor 3D file and handled by the subroutine **readiv**. The only OpenInventor graphics primitives currently supported are triangulated meshes as for example produced by Imaris (cf. chapter 2).

C.5.3 Output files

The main output file of the preprocessor is the binary restart file to be read by the PSE simulation code. Its name is defined using the **restartfile** keyword in the parameter input file (see above). The restart file contains the positions, initial strengths and attributes¹ of all the particles. Since it is a binary file, one has to make sure that byte order and precision of the target platform are matched. This means that if one wishes to produce initial restart files for a PSE simulation running on a *big endian* machine, the preprocessor has to be compiled using the -byteswapio option if it is running on a *little endian* and vice versa. Moreover, it is advisable to use Fortran KIND inquiry functions to set the accuracy parameter MK (see statements 3 and 4 in globals.f90) rather than hard-coding the bit width of a certain machine architecture.

If the parameter gnuout in the global parameter input file is set to .TRUE., a set of additional ASCII output files is generated that can be used to visualize the geometry and the particle distribution using gnuplot or some other tool. The file surface.gnu explicitly contains all the triangles to be visualized in gnuplot using sp 'surface.gnu' (see e.g. figures 5.2 and 5.7). The file mass0.gnu' contains the positions of all the particles of strength zero (i.e. the ones inside the bleached box) and mass1.gnu contains the positions of all those with strength 1. These two files can be used to visualize the initial particle distribution inside the ER geometry such as in figure 9.1. The file ghost.gnu finally contains the positions of all boundary condition image particles that are – by definition – in an r_c -neighborhood outside of the computational domain.

```
156
```

 $^{^{1}\!&}gt;0$ for particles inside the computational domain and <0 for boundary condition image particles outside the computational domain. The value 0 is reserved for MPI communication layer ghost particles in the PSE code and must not be used

If the routines Voxelize.f90 and BCdim.f90 are switched on in order to estimate the box counting dimension of the geometry at hand, three additional files are produced. The first one, BC.gnu, contains the values $s_i = \log\left(\frac{1}{\sigma_i}\right)$ and $y_i = \log \eta_i$ for all reduction steps (see section 5.2):

<si> <yi>

The second one, plotit.mac, is a gnuplot macro to visualize the measured box counting points and the best linear fit through them (recall that the slope of it corresponds to the box counting dimension of the geometry). The plots in figure 5.11 have for example been made using such automatically generated macros. The third file, voxels.gnu, contains the voxel representation of the geometry. Each voxel of value 1 is given by its position in space on a line of the file. The innermost loop is over z, the outermost over x. Thus the file's syntax is

This voxel representation can then be used to create glyph visualizations (e.g. using OpenDX) such as the ones in figures 5.3 and 5.5.

For debugging and visualization purposes, additional output to create illustrations like the one in figure 7.3 is possible from inside point_in_domain. Please see section C.5.4.11 below for further information.

Besides the various file outputs, the program also writes some information about the geometry to the system's standard output (which is a terminal console in most cases). A sample output from the erp572 simulation run of section 9.1 is given hereafter.

Initializing particle positions on regular lattice Spatial resolution (tolerance) is: 0.10E-09 Lattice spacing in x,y,z: 2.00 2.00 0.40 Bleached box has lower left corner at: 250.00 125.00 And x-,y-widths: 50.00 50.00 Reading surface description file: ../3div/erp572.iv Surface triangulation contains 35810 vertices defining 72024 triangles Triangles ... read. Bounding box of domain: xmin = 54.8070 xmax = 414.2720 ymin = 102.2590 ymax = 512.7500 zmin = -1.8534 zmax = 17.2687 Calculating centroids and outer normals of triangles ... Triangles sorted into 100x 100 bins. Number of entries in largest bin list: 367 Total number of entries in bin lists: 505051 Determining indices of particles inside domain ... 5%. 10%. 15%. 20%. 25%. 30%. 35%. 40%. 45%. 50%. 55%. 60%. 65%. 70%. 75%. 80%. 85%. 90%. 95%.100%. Initializing 545859 particles ...

```
Using Neumann boundary conditions. done.
```

Wrote restart file restart.dat

Wrote gnuplot files.

C.5.4 Source code listings

The following subsections contain the program source of all the modules listed in table C.4. Each section starts with a brief text header explaining the purpose of the routine and highlighting special features of it.

C.5.4.1 Global parameters and variables (globals.f90)

The following module contains the definitions of global parameters and variables used in more than one of the other routines.

```
1 module globals
```

```
2 IMPLICIT NONE
```

```
1----
      !GLOBALS Global variables and definitions for all routines.
         See also INIT_PART
      ! DIPLOMA THESIS WS01/02 ICOS
                                                                   ETH-ZUERICH
                   PROTEIN DIFFUSION INSIDE THE ENDOPLASMIC RETICULUM
          ! Global parameters
     ! Accuracy
     INTEGER, PARAMETER
                                          :: MK = KIND(1.0D0)
3
                                            :: I4B = SELECTED_INT_KIND(9)
4
     INTEGER, PARAMETER
      !----
     ! Global TYPE declarations
5
     TYPE ptr_to_list
        REAL(MK), DIMENSION(:), POINTER :: list
6
     END TYPE
7
      1-----
     ! Declaration of global variables
     ! Precision Tolerance for geometry handling (at least 10 times smaller
      ! than smallest occuring geometry structure)
                                            :: TOL
     REAL(MK)
8
     ! Numer of surface elements
9
     INTEGER
                                            •• M
     !\ {\tt Cut-off}\ {\tt radius}\ {\tt for}\ {\tt boundary}\ {\tt condition}\ {\tt handling}\ ({\tt mirroring})\ {\tt and}\ {\tt cell}\ {\tt list}
10
     REAL(MK)
                                            :: rc
      ! Vertices of all oriented triangles of the domain surface
11
     REAL(MK), DIMENSION(:,:,:), ALLOCATABLE :: triang
      ! centroids of all triangles
12
     REAL(MK), DIMENSION(:,:), ALLOCATABLE
                                            :: c
      ! outer normals of all triangles
     REAL(MK), DIMENSION(:,:), ALLOCATABLE
13
                                            :: normal
      ! inxed of closest triangle to every particle
14
     INTEGER, DIMENSION(:), ALLOCATABLE
                                           :: clotri
      ! distance to boundary of each point
15
     REAL(MK), DIMENSION(:), ALLOCATABLE
                                            :: dtb
```

	! Nr. of CPUs and rank of current CPU		
16	INTEGER	:: ncpu, rank	
	! Precision for MPI	-	
17	INTEGER	:: MPI_PREC	
	! Lattice spacings in x, y and z direction		
18	REAL(MK)	:: ldx, ldy, ldz	
	! produce ASCII output for gnuplot?		
19	LOGICAL	:: GNUOUT	
	! run topology check ?		
20	LOGICAL	:: CHKTOPO	
	! Bleached box. coordinates of lower left	corner and widths	
21	REAL(MK)	:: bbllx,bblly, bbwx,bbwy	
	! upper boundaries in x and y of bleached	box	
22	REAL (MK)	:: bbux, bbuy	
	! Name of PSE3D restart file to be written		
23	CHARACTER(LEN=80)	:: resfile	
	! Name of surface description file		
24	CHARACTER(LEN=80)	:: surffile	
	! Type of B.C. (0: none, 1: zero Dirichlet, 2: zero Neumann)		
25	INTEGER	:: BC	
	! math. constant		
26	REAL(MK)	:: PI	
	! triangle lists for y/z bins (point_in_domain)		
27	<pre>TYPE(ptr_to_list), DIMENSION(:,:), ALLOCA</pre>	TABLE :: bin	
	! lengths of y/z lists		
28	INTEGER, DIMENSION(:,:), ALLOCATABLE	:: nbin	
	! array of triangle lists for x/y/z cells	(boundary condition only)	
29	<pre>TYPE(ptr_to_list), DIMENSION(:,:,:), ALLOCATABLE :: cell</pre>		
	! lengths of x/y/z lists		
30	<pre>INTEGER, DIMENSION(:,:,:), ALLOCATABLE</pre>	:: ncell	
	! Bounding box		
31	REAL(MK)	:: xmin, xmax, ymin, ymax, zmin, zmax	
	! number of bins in y and z		
32	INTEGER	:: NBINY, NBINZ	
	<pre>! y and z sizes of bins</pre>		
33	REAL(MK)	:: dby, dbz	
	! number of cells in x,y,z		
34	INTEGER	:: NCELLX, NCELLY, NCELLZ	
	<pre>! x, y and z sizes of cells</pre>		
35	REAL(MK)	:: dcx, dcy, dcz	
	! voxel representation of geometry		
36	<pre>INTEGER, DIMENSION(:,:,:), ALLOCATABLE</pre>	:: voxel	
27	and module globals		

C.5.4.2 Main program (init_part.f90)

The main program of the geometry processing tool is given next. It is parallelized using the MPI message passing library. Pay special attention to the MPI testing in statements 45 to 104 which makes sure that all communication channels are up and all processors ready. Load balancing is done in statements 182 to 227 including a small load overview table that is printed to the system's standard output. All calls to MPI are encapsulated in the C preprocessor directives **#ifdef _MPI_... #endif** to make them fully transparent. The region between statements 289 and 290 is commented out and contains the calls to the routines used to translate the triangulation to a voxel representation and to determine its box counting dimension.

```
1 program init_part
```

```
2 USE globals
```

```
3 IMPLICIT NONE
```

4 #ifdef __MPI__ 5 INCLUDE 'mpif.h' 6 #endif

!-----!init_part Geometry preprocessing and particle initialization for PSE3D.
! INIT_PART reads the triangulated geometry description, determines its
! bounding box and initializes the particles on a regular lattice.
! The particle distribution is then written to a file which can be read

[!] in by PSE3D as a restart file

```
See also PSE3D
         todo:
        DIPLOMA THESIS WS01/02 ICOS
                                                                       ETH-ZUERICH
                     PROTEIN DIFFUSION INSIDE THE ENDOPLASMIC RETICULUM
      ! Declaration of external functions
      ! checks whether a point is inside the domain
      INTEGER, EXTERNAL
7
                                                :: point_in_domain
      ! Checks the topological validity of the triangulation set
      LOGICAL, EXTERNAL
8
                                                :: chk_topology
      ! Checks the orientation of the triangles (syntactical validity)
9
      LOGICAL, EXTERNAL
                                                :: chk_orientation
      ! returns the box counting dimension of the surface
     REAL(MK), EXTERNAL
10
                                               :: BCdim
      ! Declaration of local variables
      ! loop counters
      INTEGER
                                              :: i, j, k, n
11
      ! total number of particles on local processor and its global sum
12
      INTEGER
                                               :: Npart, Npsum
      ! number of ghost partcles for B.C. handling and its global sum
13
      INTEGER
                                               :: Nghost, Ngsum
      ! flag array of which particles are inside the domain
      INTEGER, DIMENSION(:,:,:), ALLOCATABLE :: flag
14
      ! position of current lattice node
      REAL(MK), DIMENSION(3)
15
                                               :: x
      ! positions of particles
      REAL(MK), DIMENSION(:,:), ALLOCATABLE
16
                                             :: xp
      ! particle flags
      INTEGER, DIMENSION(:), ALLOCATABLE
17
                                               :: ap
      ! concentrations at particle locations
18
      REAL(MK), DIMENSION(:), ALLOCATABLE
                                               :: mass
      ! I/O status flag
19
      INTEGER
                                               :: istat
      ! Number of grid points in x,y,z
20
     INTEGER
                                               :: Nx, Ny, Nz
      ! number of triangles per slave
21
      INTEGER
                                               :: Nxslave
      ! local number of triangles
22
      INTEGER
                                               :: Nxlocal
      ! lower and upper boundary of Nx on local processor
23
      INTEGER
                                              :: Nxlow, Nxhigh
      ! edge vectors and direction
24
      REAL(MK), DIMENSION(3)
                                              :: a, b
      ! intermediate variables
      REAL(MK)
25
                                              :: rx
      ! Return flag for READIV
26
      LOGICAL
                                              :: success
      ! percentage done in multiples of 5%
27
      INTEGER
                                               :: pert
      ! temporary local arrays for dtb and clotri
     REAL(MK), DIMENSION(:), ALLOCATABLE :: dtbtemp
INTEGER, DIMENSION(:), ALLOCATABLE :: clotritemp
28
29
      #ifdef __MPI__
    ! MPI init error and MPI status variable
30
31
         INTEGER
                                         :: ierror, stat(MPI_STATUS_SIZE)
         ! number of characters in processor name
32
         INTEGER
                                              :: ilen
         ! name of current processor
         CHARACTER(LEN=80)
33
                                              :: pname
34
      #endif
      ! for the voxel test
```

35 INTEGER :: voxx, voxy, voxz
```
I _____
            _____
      ! Set all global parameters to their default value
                       _____
      1-----
36
     CALL Defaults
      1------
      ! initialize and test MPI
                               _____
37
     #ifdef __MPI__
        ! Define needed precision for MPI
38
        if (MK .EQ. 4) then
39
           MPI_PREC = MPI_REAL
        elseif (MK .EQ. 8) then
40
          MPI_PREC = MPI_DOUBLE_PRECISION
41
42
        else
           WRITE (*, '(A,I4)') 'Warning: Unknown precision: ', MK
43
44
        endif
        ! Initialize MPI
45
        CALL MPI_Init(ierror)
        if(ierror .NE. 0) then
46
47
           WRITE(*,'(A)') 'Error: Initializing MPI failed. Aborting.'
48
           goto 9999
49
        end if
        CALL MPI_Comm_size(MPI_COMM_WORLD, ncpu, ierror)
50
        if(ierror .NE. 0) then
WRITE(*,'(A)') 'Error: Cannot determine number of processors. Aborting.'
51
52
53
           goto 9999
54
        end if
        CALL MPI_Comm_rank(MPI_COMM_WORLD, rank, ierror)
if(ierror .NE. 0) then
55
56
           WRITE(*,'(A)') 'Error: Cannot determine CPU rank. Aborting.'
57
           goto 9999
58
59
        end if
60
        if(rank .EQ. 0) then
61
           WRITE(*,'(A,I3,A)') 'MPI activated. Starting ',ncpu,' processors ...'
        end if
62
        CALL MPI_Get_Processor_Name(pname, ilen, ierror)
WRITE(*,'(3A,I3)') ' -> started on ',pname(1:
63
                               -> started on ',pname(1:ilen),' with rank =', rank
64
65
        CALL flush(6)
        CALL MPI_Barrier(MPI_COMM_WORLD, ierror)
66
67
        if(rank .EQ. 0) \texttt{WRITE}(\texttt{*},\texttt{'}(\texttt{A})\texttt{'}) 'All processors are up and ready.'
        ! test individual communication
68
        if(rank .EQ. 0) then
69
           do i=1,ncpu-1
              WRITE(*, '(A,I3)', ADVANCE='NO') 'Probing rank ',i,' ...
70
71
              CALL MPI_SEND(i, 1, MPI_INTEGER, i, i, &
                   MPI_COMM_WORLD, ierror)
              CALL MPI_RECV(j, 1, MPI_INTEGER, i, MPI_ANY_TAG, &
72
                   MPI_COMM_WORLD, stat, ierror)
73
              if(j .EQ. i) WRITE(*,'(A)') '[ passed ]'
74
              CALL flush(6)
75
           end do
76
        else
77
           CALL MPI_RECV(j, 1, MPI_INTEGER, 0, MPI_ANY_TAG, &
                MPI_COMM_WORLD, stat, ierror)
78
           CALL MPI_SEND(j, 1, MPI_INTEGER, 0, rank, &
                MPI_COMM_WORLD, ierror)
79
        end if
        ! test broadcast
80
        i = 0
        if(rank .EQ. 0) then
81
82
           i = 23
83
           Npart = 0
           WRITE(*,'(A)')
84
85
           WRITE(*,'(A)',ADVANCE='NO') 'Probing broadcast to all ... '
86
        end if
87
        CALL MPI_BCAST(i, 1, MPI_INTEGER, 0, MPI_COMM_WORLD, ierror)
88
        if(rank .EQ. 0) then
89
           do j=1,ncpu-1
90
              CALL MPI_RECV(k, 1, MPI_INTEGER, j, MPI_ANY_TAG, &
                   MPI_COMM_WORLD, stat, ierror)
```

```
91
               if(k .EQ. 23) then
                  WRITE(*,'(I3)', ADVANCE='NO') j
92
93
                  Npart = Npart + 1
94
               end if
95
            end do
96
            if(Npart .EQ. ncpu-1) then
97
               WRITE(*,'(A)') ' [ passed ]'
98
            else
99
               WRITE(*,'(A)') ' [ failed ]'
100
               goto 9999
101
            end if
102
         else
            CALL MPI_SEND(i, 1, MPI_INTEGER, 0, rank, &
103
                 MPI_COMM_WORLD, ierror)
104
         end if
105
      #else
106
         ncpu = 1
107
         rank = 0
108
      #endif
      ! Read input files on ROOT
      !----
109
      if(rank .EQ. 0) then
        CALL readinput(success)
110
111
        if(.NOT. success) then
112
           WRITE(*,'(A)') 'Failed reading input files. Aborting.'
           goto 9998
113
114
        end if
        ! dump geometry for GNUPLOT (splot '...' w 1)
        if(GNUOUT) then
115
           OPEN(30, FILE='surface.gnu', STATUS='REPLACE')
116
117
           do i=1.M
              do j=1,3
118
119
                 WRITE(30,*) triang(1,j,i), triang(2,j,i), triang(3,j,i)
120
              end do
              WRITE(30,'(A)')
121
122
             WRITE(30,'(A)')
           end do
123
          CLOSE(30)
124
        end if
125
126
      end if
      ! Broadcast geometry to slave processors
      !---
127
      #ifdef __MPI__
         !\ \mbox{broadcast} number of triangles, global tolerance and grid spacings
128
         CALL MPI_BCAST(M, 1, MPI_INTEGER, 0, MPI_COMM_WORLD, ierror)
129
         CALL MPI_BCAST(TOL, 1, MPI_PREC, 0, MPI_COMM_WORLD, ierror)
130
         CALL MPI_BCAST(ldx, 1, MPI_PREC, 0, MPI_COMM_WORLD, ierror)
131
         CALL MPI_BCAST(ldy, 1, MPI_PREC, 0, MPI_COMM_WORLD, ierror)
132
         CALL MPI_BCAST(ldz, 1, MPI_PREC, 0, MPI_COMM_WORLD, ierror)
133
         if(rank .NE. 0) then
            ! allocate local memory for triangles on slave
134
            istat = 0
            ! first index: x,y,z; second: point number; third: triangle number
135
            ALLOCATE(triang(3,3,M), c(3,M), normal(3,M))
136
            if (istat .NE. 0) then
137
               WRITE(*,'(2A)') 'Unable to allocate memory for triangles on ' &
                  ,pname(1:ilen)
138
               goto 9999
            endif
139
140
         end if
          wait for all slaves to allocate memory
141
         CALL MPI_Barrier(MPI_COMM_WORLD, ierror)
         ! Broadcast triangles
142
         do k=1,M
                           ! loop over all triangles
143
            do j=1,3
                           ! loop over all vertices
144
               CALL MPI_BCAST(triang(:,j,k), 3, MPI_PREC, 0,
                                                                 &
                    MPI_COMM_WORLD, ierror)
145
            end do
146
         end do
147
      #endif
```

```
1-----
      ! Determine bounding box and set up grid
                                         _____
      1-----
148
     xmin = HUGE(xmin)
149
     xmax = -HUGE(xmax)
150
     ymin = HUGE(ymin)
151
     ymax = -HUGE(ymax)
152
     zmin = HUGE(zmin)
153
     zmax = -HUGE(zmax)
154
     do i=1,M
155
        do j=1,3
           if (triang(1,j,i) .LT. xmin) xmin = triang(1,j,i)
156
157
           if (triang(1,j,i) .GT. xmax) xmax = triang(1,j,i)
158
           if (triang(2,j,i) .LT. ymin) ymin = triang(2,j,i)
159
           if (triang(2,j,i) .GT. ymax) ymax = triang(2,j,i)
           if (triang(3,j,i) .LT. zmin) zmin = triang(3,j,i)
160
           if (triang(3,j,i) .GT. zmax) zmax = triang(3,j,i)
161
        end do
162
163
     end do
164
     if(rank .EQ. 0) then
165
        WRITE(*, '(A,F12.4,A,F12.4)') 'Bounding box of domain: xmin = ',xmin,' &
          xmax = ',xmax
        WRITE(*,'(A,F12.4,A,F12.4)') '
                                                           ymin = ',ymin,' &
166
          ymax = ',ymax
        WRITE(*,'(A,F12.4,A,F12.4)') '
167
                                                           zmin = ',zmin,' &
          zmax = ',zmax
        WRITE(*,'(A)')
168
     end if
169
     ! enlarge BBox by tolerance to avoid numerical problems at borders
170
     xmin = xmin-TOL
     xmax = xmax+TOL
171
172
     ymin = ymin-TOL
173
     ymax = ymax+TOL
     zmin = zmin-TOL
174
175
     zmax = zmax+TOL
     ! determine number of particle points
176
     Nx = ceiling((xmax-xmin)/ldx)-1
177
     Ny = ceiling((ymax-ymin)/ldy)-1
     Nz = ceiling((zmax-zmin)/ldz)-1
178
     ! and adjust grid spacing if needed
179
     ldx = (xmax-xmin)/real(Nx)
     ldy = (ymax-ymin)/real(Ny)
180
    ldz = (zmax-zmin)/real(Nz)
181
     ! local chunk sizes
     Nxslave = nint(real(Nx)/real(ncpu))
182
183
     if(rank .EQ. 0) then
184
        Nxlocal = Nx-Nxslave*(ncpu-1) ! chunk size on ROOT
                         ! lower x index
185
        Nxlow = 1
186
        Nxhigh = Nxlocal
                              ! upper x index
187
     else
        Nxlocal = Nxslave
188
189
        Nxlow = (rank-1)*Nxslave+1+(Nx-Nxslave*(ncpu-1))  ! lower x index
190
        Nxhigh = rank*Nxslave+(Nx-Nxslave*(ncpu-1))
                                                      ! upper x index
191
     end if
192
     #ifdef __MPI_.
193
       if(rank .EQ. 0) then
           ! print nice load allocation table
           WRITE(*,'(A)', ADVANCE='NO') ' CPU #
                                              | ,
194
           do i=0,ncpu-1
195
             WRITE(*,'(A,I3,A)',ADVANCE='NO') ' ',i,' | '
196
197
           end do
           WRITE(*,'(A)')
198
           WRITE(*, '(A)', ADVANCE='NO') '-----|-'
199
           do i=0,ncpu-1
200
201
             WRITE(*,'(A)', ADVANCE='NO') '-----|'
202
           end do
203
           WRITE(*,'(A)')
204
           WRITE(*, '(A, I6, A)', ADVANCE='NO') ' Nxlocal | ',Nxlocal,' | '
205
           do i=1,ncpu-1
             WRITE(*,'(I6,A)',ADVANCE='NO') Nxslave,' | '
206
207
           end do
208
           WRITE(*,'(A)')
           WRITE(*,'(A,I6,A)',ADVANCE='NO') ' Nxlow | ',Nxlow,' | '
209
210
           do i=1,ncpu-1
```

```
211
              WRITE(*,'(I6,A)',ADVANCE='NO') (i-1)*Nxslave+1+ &
                   (Nx-Nxslave*(ncpu-1)),' |
212
           end do
           WRITE(*,'(A)')
213
214
           WRITE(*,'(A,I6,A)',ADVANCE='NO') ' Nxhigh | ',Nxhigh,' | '
215
           do i=1,ncpu-1
216
              WRITE(*,'(I6,A)',ADVANCE='NO') i*Nxslave+(Nx-Nxslave*(ncpu-1)),' | '
217
           end do
218
           WRITE(*,'(A)')
219
           WRITE(*,'(A)', ADVANCE='NO') '-----|-'
220
           do i=0,ncpu-1
221
             WRITE(*, '(A)', ADVANCE='NO') '-----|'
222
           end do
223
           WRITE(*,'(A)')
224
           WRITE(*,'(A)')
225
           CALL flush(6)
226
        end if
227
     #endif
      ! Calculate centroids and normals
228
     if(rank .EQ. 0) then
       WRITE(*,'(A)') 'Calculating centroids and outer normals of triangles ...'
229
230
     end if
231
     do i=1,M
        ! calculate centroids of triangles
232
        c(:,i) = (1.0/3.0)*(triang(:,1,i)+triang(:,2,i)+triang(:,3,i))
        ! calculate outer normals of triangles
233
        a = triang(:,2,i)-triang(:,1,i)
        b = triang(:,3,i)-triang(:,1,i)
234
        ! cross product => n
        normal(1,i) = a(2)*b(3) - a(3)*b(2)
235
        normal(2,i) = a(3)*b(1) - a(1)*b(3)
236
        normal(3,i) = a(1)*b(2) - a(2)*b(1)
237
        ! normalize normals
238
        rx = sqrt(sum(normal(:,i)**2))
239
        normal(:,i) = normal(:,i)/rx
240
     end do
      1-----
      ! Set up triangle lists and sort triangles
      !---
241
     CALL SortT
     if(rank .EQ. 0) then
242
243
        WRITE(*,'(A)')
        WRITE(*,'(A,I4,A,I4,A)') 'Triangles sorted into ',NBINY,'x',NBINZ,' bins.'
244
        WRITE(*,'(A,I6)') 'Number of entries in largest bin list: ',maxval(nbin)
245
        WRITE(*,'(A,I7)') 'Total number of entries in bin lists: ',sum(nbin)
246
247
        if(BC .GT. 0) then
           WRITE(*,'(A)')
WRITE(*,'(A,I4,A,I4,A,I4,A)') 'Triangles sorted into ',NCELLX,'x', &
248
249
              NCELLY,'x',NCELLZ,' cells.'
250
           WRITE(*,'(A,I6)') 'Number of entries in largest cell list: ',maxval(ncell)
           WRITE(*,'(A,I7)') 'Total number of entries in cell lists: ',sum(ncell)
251
252
        end if
253
        WRITE(*,'(A)')
254
     end if
      1-----
      ! Check topology and orientation
255
     if(rank .EQ. 0 .AND. CHKTOPO) then
        WRITE(*,'(A)')
256
257
        if(.NOT. chk_topology()) then
258
           WRITE(*,'(A)') 'Triangulation is not topologically valid. Aborting.'
259
           goto 9998
260
        end if
261
        if(.NOT. chk_orientation()) then
262
           WRITE(*,'(A)') 'Triangulation is not syntactically valid. Aborting.'
263
           goto 9998
264
        endif
        WRITE(*,'(A)')
265
     end if
266
      !-----
```

```
164
```

```
! Initialize particles on regular lattice
                                                               _____
      ! determine number of particles
267
      if(rank .EQ. 0) then
268
        WRITE(*,'(A)') 'Determining indices of particles inside domain ...'
269
      end if
      ! allocate temporary flag array
270
      if(rank .EQ. 0) then
271
         ALLOCATE(flag(Nx,Ny,Nz))
272
         if (istat .NE. 0) then
273
            WRITE(*,'(A)') 'Unable to allocate memory for flags on ROOT'
274
           goto 9998
275
         end if
276
      else
277
         ALLOCATE(flag(Nxlocal,Ny,Nz))
278
         if (istat .NE. 0) then
            WRITE(*,'(A,I3)') 'Unable to allocate memory for flags on CPU', rank
279
280
           goto 9998
281
         end if
282
      end if
283
     if(BC .GT. 0) then
         ! allocate temporary arrays for distances and triangle indices
284
         ALLOCATE(dtbtemp(Nxlocal*Ny*Nz), clotritemp(Nxlocal*Ny*Nz))
285
         if (istat .NE. 0) then
286
            WRITE(*,'(A,I3)') 'Unable to allocate memory for dtb, clotri on CPU', rank
           goto 9997
287
288
         end if
      end if
289
      ! uncomment this to get and save a voxel representation of the
      ! geometry (not needed for determining the box counting dimension
      ! as BCdim will call Voxelize on its own)
      !!$rx = 0.05_MK
                         ! voxel size
      !!$CALL Voxelize(rx,rx,rx,voxx,voxy,voxz,istat)
      !!$if(ALLOCATED(voxel)) then
      !!$ print*,'vox allocated'
           print*,'sum of voxels: ',sum(voxel)
      !!$
            print*, 'number of voxels in x,y,z: ',voxx,voxy,voxz
      11$
           OPEN(50, FILE='voxels.gnu', STATUS='REPLACE', ACTION='WRITE')
      !!$
      11$
           do i=1.voxx
      11$
              do j=1,voxy
      !!$
                  do k=1,voxz
                    if(voxel(i,j,k) .EQ. 1) WRITE(50,*) xmin+(i-0.5)*rx, &
      11$
      !!$
                       ymin+(j-0.5)*rx, zmin+(k-0.5)*rx
                  end do
      11$
      !!$
               end do
      !!$
           end do
      !!$
           CLOSE(60)
      11$
           print*,'wrote voxels to file'
      !!$
           DEALLOCATE(voxel)
      U$end if
      ! uncomment this to determine the box counting dimension of the geomerty
      !!$print*,'invoking BCdim ...'
      !!$rx = BCdim((xmax-xmin)/500.0_MK,(ymax-ymin)/500.0_MK, &
      !!$ (zmax-zmin)/300.0_MK, 0.30_MK, istat)
      !!$if(istat .NE. 0) then
      !!$ WRITE(*,'(A)') 'ERROR in BCdim'
      !!$ STOP
      !!$end if
      !!$print*,'box counting dimension is: ',rx
    flag = 0 ! 0: grid point outside,
290
                 ! 1: grid point inside,
                 ! 2: grid point in rc-neighborhood of boundary
     Npart = 0
291
292
      Nghost = 0
293
     pert = 0
294
      do i=Nxlow,Nxhigh
295
        x(1) = xmin+(real(i)-0.5_MK)*ldx ! use staggered mesh to avoid
296
         do j=1,Ny
                                            ! having points exactly on the boundary
297
           x(2) = ymin+(real(j)-0.5_MK)*ldy! as they would coincide with their
                                            ! images => PSE would be unstable!
298
            do k=1,Nz
299
               x(3) = zmin+(real(k)-0.5_MK)*ldz
300
               if(BC .GT. 0) then
                  flag(i-Nxlow+1,j,k) = point_in_domain(x,dtbtemp(Npart+1), &
301
                       clotritemp(Npart+1))
```

302 else 303 flag(i-Nxlow+1,j,k) = point_in_domain(x,rx,n) 304 end if 305 if(flag(i-Nxlow+1,j,k) .GT. 0) Npart = Npart + 1 306 if(flag(i-Nxlow+1,j,k) .EQ. 2) Nghost = Nghost + 1 307 if(rank .EQ. 0) then ! print percentage status every 5% 308 if((((i-1)*Ny*Nz+(j-1)*Nz+k)/(Nxlocal*Ny*Nz+0.0) .GT. & pert*0.05) then 309 pert = pert + 1 310 WRITE(6, '(I3, A)', ADVANCE='NO') pert*5, '%.' 311 CALL flush(6) ! force immediate output 312 end if 313 end if end do 314 315 end do end do 316 317 WRITE(*,'(A)') ! free memory of triangle lists 318 CALL AllocateLL(2,0,0,0,istat) #ifdef __MPI__
 ! wait for all processors to terminate 319 320 CALL MPI_Barrier(MPI_COMM_WORLD, ierror) ! sum Npart onto ROOT CALL MPI_Reduce(Npart, Npsum, 1, MPI_INTEGER, MPI_SUM, 0, 321 & MPI_COMM_WORLD, ierror) ! sum Nghost onto ROOT CALL MPI_Reduce(Nghost, Ngsum, 1, MPI_INTEGER, MPI_SUM, 0, & 322 MPI_COMM_WORLD, ierror) if(rank .NE. 0) then 323 ! send local part of array flag to ROOT CALL MPI_SEND(flag(:,:,:), Nxlocal*Ny*Nz, MPI_INTEGER, 0, & 324 rank, MPI_COMM_WORLD, ierror) 325 else ! receive flags from all slaves 326 do i=1.ncpu-1 CALL MPI_RECV(flag(Nxlocal+1+(i-1)*Nxslave:Nxlocal+i*Nxslave, & 327 1:Ny,1:Nz),Nxslave*Ny*Nz, MPI_INTEGER, i, i, Яr. MPI_COMM_WORLD, stat, ierror) 328 end do 329 end if 330 #else 331 Npsum = Npart Ngsum = Nghost 332 333 #endif 334 if(BC .GT. 0) then 335 if(rank .EQ. 0) then ! allocate final arrays of correct size 336 ALLOCATE(dtb(Npsum), clotri(Npsum)) 337 if (istat .NE. 0) then 338 WRITE(*,'(A,I3)') 'Unable to allocate memory for dtb, clotri on ROOT' 339 goto 9997 340 end if ! copy root part into final arrays 341 dtb(1:Npart) = dtbtemp(1:Npart) 342 clotri(1:Npart) = clotritemp(1:Npart) 343 n = Npart ! initialize position counter #ifdef __MPI__ 344 ! receive contributions of all slaves 345 do i=1,ncpu-1 ! receive number of particles from slave i 346 CALL MPI_RECV(j, 1, MPI_INTEGER, i, i, MPI_COMM_WORLD, stat, ierror) ! receive distances to boundary for its points CALL MPI_RECV(dtb(n+1:n+j),j , MPI_PREC, i, i, 347 & MPI_COMM_WORLD, stat, ierror) receive indices of closest triangles CALL MPI_RECV(clotri(n+1:n+j),j, MPI_PREC, i, i, 348 & MPI_COMM_WORLD, stat, ierror) 349 end do 350 else ! send my number of particles to ROOT CALL MPI_SEND(Npart, 1, MPI_INTEGER, 0, rank, MPI_COMM_WORLD, ierror) ! send the distances to boundary of my particles 351 CALL MPI_SEND(dtbtemp(1:Npart), Npart, MPI_PREC, 0, rank, 352 & MPI_COMM_WORLD, ierror) ! send indices of closest triangles 353 CALL MPI_SEND(clotritemp(1:Npart), Npart, MPI_PREC, 0, rank, &

```
MPI_COMM_WORLD, ierror)
354
      #endif
355
         end if
         ! deallocate temporary arrays
356
         DEALLOCATE(dtbtemp)
357
         DEALLOCATE(clotritemp)
358
      end if
359
      if(rank .EQ. 0) then
         ! be nice and talk to the user
360
         WRITE(*,'(A,I7,A)') 'Initializing ',Npsum,' particles ...'
         if(BC .EQ. 0) then
361
362
            WRITE(*,'(A)') 'Not using boundary conditions.'
363
         end if
364
         if(BC .EQ. 1) then
365
            WRITE(*,'(A,I7,A)') 'Using zero Dirichlet B.C. with ',Ngsum, &
                ' ghost particles'
366
         end if
         if(BC .EQ. 2) then
367
            WRITE(*,'(A,I7,A)') 'Using zero Neumann B.C. with ',Ngsum, &
368
                ' ghost particles'
369
         end if
         ! add ghost particles to the number of particles
         Npsum = Npsum + Ngsum
370
         ! global array to hold all particles
371
         istat = 0
         ! first index: x,y,z, second index: particle number
372
         ALLOCATE(xp(3,Npsum), mass(Npsum), ap(Npsum))
         if (istat .NE. 0) then
    WRITE(*,'(A)') 'Unable to allocate memory for particles on ROOT'
373
374
            goto 9997
375
376
         end if
         ! initializing particle positions
377
         n = 0
                             ! ID of current particle
         Npart = 0
378
                             ! total number of particles so far
379
         do i=1.Nx
            x(1) = xmin+(real(i)-0.5_MK)*ldx
380
            do j=1,Ny
    x(2) = ymin+(real(j)-0.5_MK)*ldy
381
382
383
                do k=1,Nz
                   x(3) = zmin+(real(k)-0.5_MK)*ldz
384
385
                   if(flag(i,j,k) .GT. 0) then
386
                      n = n + 1
                      Npart = Npart + 1
387
                      xp(1:3,Npart) = x(1:3)
388
                                               ! regular particle
389
                      ap(Npart) = Npart
                                               ! particle ID as attribute
                      if(flag(i,j,k) .EQ. 2) then
390
                         ! boundary particle => needs mirror image as ghost
391
                         xp(1:3,Npart+1) = x(1:3)+2.0*dtb(n)*normal(1:3,clotri(n))
                         ! negative particle index as attribute
392
                         ap(Npart+1) = -Npart
                         Npart = Npart + 1
393
394
                      end if
395
                   end if
396
                end do
397
            end do
398
         end do
399
      end if
       ! deallocate flags on all processors
      DEALLOCATE(flag)
400
401
      if(rank .EQ. 0) then
         ! initialize particle strengths
402
         do i=1,Npsum
403
            if(ap(i) .GT. 0) then
                                     ! regular particles
               if((xp(1,i) .GE. bbllx) .AND. (xp(1,i) .LE. bbllx+bbwx) .AND.
(xp(2,i) .GE. bblly) .AND. (xp(2,i) .LE. bblly+bbwy)) then
404
                                                                                   &
                   ! no concentration inside bleached box
                   mass(i) = 0.0
405
406
                else
                   ! unit concentration c0 outside
```

```
407
                 mass(i) = 1.0
408
              end if
409
            end if
410
            if(ap(i) .LT. 0) then
                                                       ! ghost particles
411
               if(BC .EQ. 1) mass(i) = -mass(-ap(i))
                                                     ! Dirichlet
412
               if(BC .EQ. 2) mass(i) = mass(-ap(i))
                                                      ! Neumann
413
            end if
414
         end do
        WRITE(*,'(A)') 'done.'
415
416
        WRITE(*,'(A)')
         ! Write new restart file for PSE3D
         1--
                                          _____
417
        OPEN(20, FILE=resfile, ACTION='WRITE', FORM='UNFORMATTED', IOSTAT=istat)
418
        CLOSE(20, STATUS='DELETE')
419
        OPEN(20, FILE=resfile, ACTION='WRITE', FORM='UNFORMATTED', IOSTAT=istat)
420
        WRITE(20)Npsum, 0.0_MK
421
        do i=1,Npsum
          WRITE(20)xp(:,i), mass(i), ap(i)
422
423
        end do
424
        CLOSE(20)
        WRITE(*,'(2A)') 'Wrote restart file ', trim(resfile)
425
426
        WRITE(*,'(A)')
         ! Write gnuplot files for testing purposes
         1---
        if(GNUOUT) then
427
           OPEN(20, FILE='mass1.gnu', STATUS='REPLACE', ACTION='WRITE')
428
           OPEN(30, FILE='mass0.gnu', STATUS='REPLACE', ACTION='WRITE')
OPEN(40, FILE='ghost.gnu', STATUS='REPLACE', ACTION='WRITE')
429
430
431
           do i=1,Npsum
              if(ap(i) .GT. 0) then
432
                 if(mass(i) .EQ. 0.0) then
WRITE(30, *) xp(:,i)
433
434
435
                 else
                   WRITE(20, *) xp(:,i)
436
                 end if
437
              end if
438
              if(ap(i) .LT. 0) then
439
440
                 WRITE(40, *) xp(:,i)
441
              end if
442
           end do
           CLOSE(20)
443
444
           CLOSE(30)
           WRITE(*,'(2A)') 'Wrote gnuplot files.'
445
446
        end if
447
     end if
             _____
      1----
      ! Deallocate memory and terminate
                                          _____
      1--
448
     9997 if(ALLOCATED(dtb)) DEALLOCATE(dtb)
449
     if(ALLOCATED(clotri)) DEALLOCATE(clotri)
450
     if(ALLOCATED(xp)) DEALLOCATE(xp)
451
     if(ALLOCATED(ap)) DEALLOCATE(ap)
452
     if(ALLOCATED(mass)) DEALLOCATE(mass)
453
     9998 if(ALLOCATED(triang)) DEALLOCATE(triang)
454
     if(ALLOCATED(c)) DEALLOCATE(c)
455
     if(ALLOCATED(normal)) DEALLOCATE(normal)
456
     9999 CONTINUE
457
     #ifdef __MPI_.
458
        CALL MPI_Finalize(ierror)
459
     #endif
460
     END program init_part
```

C.5.4.3 Set global default values (Defaults.f90)

The following subroutine assigns default values to all global parameters. This is needed to make sure that all parameters contains meaningful values even if they are not explicitly specified in the parameter input file described in section C.5.2.

```
1
     subroutine Defaults
     USE globals
2
     IMPLICIT NONE
3
     1-----
     !Defaults Assigns default values to all parameter variables.
         The settings of this file can be overridden by specifying a
         different value for a parameter in the parameter input file of
        the code.
         See also INIT_PART
         todo:
     ! DIPLOMA THESIS WS01/02 ICOS
                                                                 ETH-ZUERICH
                  PROTEIN DIFFUSION INSIDE THE ENDOPLASMIC RETICULUM
     Specify default values here
     ! Precision Tolerance for geometry handling (at least 10 times smaller
     ! than smallest occuring geometry structure)
TOL = 1e-10_MK
4
     ! Cut-off radius for boundary condition handling (mirroring)
5
                         = 1.0_MK
     rc
     ! Nr. of CPUs and rank of current CPU
6
                         = 1
     ncpu
7
                         = 0
     rank
     ! Lattice spacings in x, y and z direction
ldx = 0.1_MK
8
                          = 0.1_MK
9
     ldy
10
                         = 0.1_MK
     ldz
     ! produce ASCII output for gnuplot?
GNUOUT = .FALSE.
11
     ! run topology check ?
12
     CHKTOPO
                          = .FALSE.
     ! Bleached box. coordinates of lower left corner and widths
13
                         = 2.0 MK
     bbllx
14
                         = 2.0_MK
     bbllv
                         = 1.0_MK
15
     bbwx
                         = 1.0_MK
16
     bbwy
     ! Name of PSE3D restart file to be written
resfile = 'restart.dat'
17
     ! Name of surface description file
18
                         = 'box.surf'
     surffile
     ! Type of B.C. (0: none, 1: zero Dirichlet, 2: zero Neumann)
19
     BC
                         = 0
     ! number of bins in y and \boldsymbol{z}
20
     NBINY
                          = 1
21
     NBINZ
                          = 1
     ! math. constant
22
     if(MK .EQ. KIND(1.0D0)) then
23
       PI = 4.0_MK*datan(1.0_MK)
                                 ! double prec version of arctan
24
     else
       PI = 4.0_MK*atan(1.0_MK)
25
                                ! single prec version of arctan
26
     end if
```

27 END subroutine Defaults

C.5.4.4 Read all input files (readinput.f90)

subroutine readinput(res)

This routine is responsible for reading all the input files, i.e. the global parameter file and the triangulated geometry. Its argument **res** serves as a status flag. If it is zero upon completion of the routine, no errors occurred. Otherwise the input files contained one or more errors.

```
2
     USE globals
3
     IMPLICIT NONE
     1-----
                                                           !readinput Reads the parameter and geometry input files.
        READINPUT reads all input files for \ensuremath{\texttt{INIT\_PART}} and initializes
        basic variables on the ROOT processor. res will be .TRUE. if
        everything was successful, else .FALSE.
        See also INIT_PART
        todo:
     ! DIPLOMA THESIS WS01/02 ICOS
                                                             ETH-ZUERICH
                 PROTEIN DIFFUSION INSIDE THE ENDOPLASMIC RETICULUM
     ! Declaration of external functions
     ! funtion to read and parse the global parameter file
     INTEGER, EXTERNAL
4
                                       :: ReadParams
     ! Declaration of input/output variables
     ! .TRUE. if reading was succesful
5
     LOGICAL, INTENT(OUT)
                                        :: res
                     _____
     ! Declaration of local variables
     ! loop counter
6
     INTEGER
                                        :: i
     ! Return flag for READIV
7
     LOGICAL
                                        :: success
     ! File format description line
     CHARACTER(LEN=80)
8
                                        :: filefmt
     ! I/O status flag
     INTEGER
9
                                        :: istat
     res = .TRUE.
10
     1-----
     ! Read parameter file
     1--
11
     if(ReadParams('preproc.in') .NE. 0) then
12
       WRITE(*,'(A)') 'Error reading parameter file restart.in! aborting.'
13
       res = .FALSE.
14
       return
15
     end if
     ! if B.C. invalid: set to none
     if((BC .NE. 0) .AND. (BC .NE. 1) .AND. (BC .NE. 2)) then
16
17
       WRITE(*, '(A)') 'Invalid boundary condition specified. Disabling it.'
18
       BC = 0
19
     end if
```

```
20
      WRITE(*,'(A)')
      WRITE(*,'(A)') 'Initializing particle positions on regular lattice'
21
22
      WRITE(*,'(A,E12.2)') 'Spatial resolution (tolerance) is: ', TOL
23
      WRITE(*,'(A)')
24
      WRITE(*,'(A,3F6.2)') 'Lattice spacing in x,y,z: ',ldx,ldy,ldz
25
      WRITE(*,'(A)')
26
      WRITE(*,'(A,2F8.2)') 'Bleached box has lower left corner at: ',bbllx,bblly
27
      WRITE(*,'(A,2F8.2)') '
                                                  And x-,y-widths: ',bbwx,bbwy
28
      WRITE(*,'(A)')
      1-----
               _____
      ! Read in triangulated geometry
29
      WRITE(*,'(2A)') 'Reading surface description file: ', trim(surffile)
30
      OPEN(20, FILE=surffile, STATUS='OLD')
      ! determine file format
      READ(20,'(A)') filefmt
31
      if(filefmt .EQ. '#ASCII triangulated surface description') then
32
         READ(20, '(A)')
33
         READ(20, '(A)')
34
         READ(20,'(A)')
35
36
         READ(20, '(A)')
37
         READ(20, '(A)')
         READ(20, '(A)')
38
39
         READ(20,*) M
40
         READ(20, '(A)')
41
         READ(20, '(A)')
         READ(20,'(A)')
42
43
         READ(20, '(A)')
44
         READ(20, '(A)')
         READ(20,'(A)')
45
         WRITE(*,'(A,I6,A)') 'Surface triangulation contains ',M,' triangles'
46
47
         istat = 0
         ! first index: x,y,z; second: point number; third: triangle number
         ALLOCATE(triang(3,3,M), c(3,M), normal(3,M))
48
         if (istat .NE. 0) then
    WRITE(*,'(A)') 'Unable to allocate memory for triangles'
49
50
            res = .FALSE.
51
52
            return
         end if
53
54
         do i=1,M
           READ(20,*) triang(:,1,i), triang(:,2,i), triang(:,3,i)
55
56
         end do
57
         CLOSE(20)
         WRITE(*, '(A)') 'Triangles ... read.'
WRITE(*, '(A)')
58
59
      else if(filefmt(1:9) .EQ. '#Inventor') then
60
         ! read OpenInventor file format
61
         if(filefmt(16:20) .EQ. 'ascii') then
62
            CLOSE(20)
                                              ! readiv expects closed file
63
            CALL READIV(surffile, success)
                                              ! read Inventor file
64
            if(.NOT. success) then
65
              WRITE(*,'(A)') 'Failed reading Inventor file.'
66
               res = .FALSE.
67
               return
68
            end if
69
         end if
70
         if(filefmt(16:21) .EQ. 'binary') then
71
            WRITE(*,'(A)') 'The surface file is a binary Inventor file.'
72
            WRITE(*,'(A)') 'Binary files are not supported. Please convert it to an'
73
            WRITE(*,'(A)') 'ASCII file using ivcat on a SGI workstation.'
74
            res = .FALSE.
75
            return
76
         end if
77
      else
78
         WRITE(*,'(A)') 'Unknown geometry file format.'
79
         res = .FALSE.
80
        return
      end if
81
82
      CLOSE(20)
83
      return
84
      END subroutine readinput
```

C.5.4.5 Read parameter input file (ReadParams.f90)

function ReadParams(ctrlfile)

The routine ReadParams.f90 is called by readinput.f90 in order to read and parse the global parameter input file according to the syntax rules stated in section C.5.2. The name of the file to be read it passed to it in the argument variable ctrlfile. ReadParams is a function and it returns zero if no errors occurred, 1 otherwise.

2 USE globals 3 IMPLICIT NONE 1----_____ !ReadParams Reads the parameter file "ctrlfile" returns 0 if no errors occured, 1 otherwise. See also init_part, readinput todo: 1==== ! DIPLOMA THESIS WS01/02 ICOS ETH-ZUERICH PROTEIN DIFFUSION INSIDE THE ENDOPLASMIC RETICULUM ! Input/Output arguments ! name of the file to be read 4 CHARACTER(LEN=*), INTENT(IN) :: ctrlfile ! return flag 5 INTEGER :: ReadParams ! Declaration of local variables ! loop counters 6 INTEGER :: i,j ! position index 7 INTEGER :: idx,i1,i2 ! read buffer CHARACTER(LEN=256) 8 :: cbuf ! file unit to be used for reading the parameter file 9 INTEGER :: iUnit ! string length, I/O error flag INTEGER 10 :: ilen,ios ! length of filename INTEGER 11 :: ilenctrl ! line counter INTEGER 12 :: ibc.iline ! value and argument of current assignment statement 13 CHARACTER(LEN=256) :: cvalue,carg CHARACTER(LEN=256) 14 :: bcloc.tvalue ! flag to indicate whether the file exists at all 15 LOGICAL :: lExist ! Definition of file unit 1-iUnit = 20 16 1_____ ! Check that the parameter file exists _____ 1--17 ilenctrl = LEN_TRIM(ctrlfile) ! length of filename string 18 INQUIRE(FILE=ctrlfile(1:ilenctrl), EXIST=lExist) 19 if(.NOT. lExist) then WRITE(*,'(2A)')'No such file: ',ctrlfile(1:ilenctrl) 20

```
21
      ReadParams = 1
22
      goto 9999
23
    end if
    1------
    ! open the file
    _____
24
    OPEN(iUnit, FILE=ctrlfile(1:ilenctrl), IOSTAT=ios, ACTION='READ')
25
    if(ios .NE. 0) then
26
      WRITE(*,'(2A)')'Failed to open file: ',ctrlfile(1:ilenctrl)
27
      ReadParams = 1
28
      goto 9999
29
    end if
    1---
    ! scan file
    ! ---
30
    iline = 0
31
    do
      iline = iline + 1
32
                         ! increment line
33
      READ(iUnit,'(A)',END=100,ERR=200) cbuf
34
      ilen = LEN_TRIM(cbuf)
       !-
      ! Skip comment or empty lines
      if(ilen .GT. 0 .AND. cbuf(1:1) .NE. '#') then
35
         !-----
         ! Remove all spaces
                            -----
         !----
36
         i = 0
37
         do i=1,ilen
          if(cbuf(i:i) .NE. ' ') then
38
39
             j = j + 1
40
             cbuf(j:j) = cbuf(i:i)
41
           end if
         end do
42
43
         ilen = j
                  ! update length of string
         1-----
                           -----
         ! Find position of =
         1---
                                  _____
44
         idx = INDEX(cbuf,'=')
         1-----
         ! Exit if = is missing
                                 _____
         if(idx .LT. 0) then
45
46
           WRITE(*,'(A,I5)')'Incorrect line: ',iline
47
           ReadParams = 1
48
           goto 9999
49
         end if
         1-----
                              _____
         ! Get argument and value
         1--
50
         carg = ADJUSTL(cbuf(1:idx-1))
51
         cvalue = ADJUSTL(cbuf(idx+1:ilen))
         1-----
               _____
         ! Convert to upper case
52
         CALL UpperCase(carg,idx-1)
         I------
                             ------
         ! Parse and read input data
                                 _____
53
         if(carg .EQ. 'SURFACEFILE') then
           READ(cvalue,'(A)',IOSTAT=ios,ERR=200) surffile
54
         elseif (carg .EQ. 'RESTARTFILE') then
55
56
           READ(cvalue, '(A)', IOSTAT=ios, ERR=200) resfile
         elseif (carg .EQ. 'DX') then
57
           READ(cvalue,*,IOSTAT=ios,ERR=200) ldx,ldy,ldz
58
         elseif (carg .EQ. 'BLEACHBOX') then
59
           READ(cvalue,*,IOSTAT=ios,ERR=200) bbllx,bblly,bbux,bbuy
60
61
           bbwx = bbux-bbllx
           bbwy = bbuy-bblly
62
```

```
63
            elseif (carg .EQ. 'CUTOFF') then
64
               READ(cvalue,*,IOSTAT=ios,ERR=200) rc
65
            elseif (carg .EQ. 'BOUNDARYCONDITION') then
66
               READ(cvalue,*,IOSTAT=ios,ERR=200) BC
67
            elseif (carg .EQ. 'GNUOUT') then
68
               READ(cvalue,*,IOSTAT=ios,ERR=200) GNUOUT
69
            elseif (carg .EQ. 'CHECKTOPOLOGY') then
70
               READ(cvalue,*,IOSTAT=ios,ERR=200) CHKTOPO
71
            elseif (carg .EQ. 'TOLERANCE') then
72
               READ(cvalue,*,IOSTAT=ios,ERR=200) TOL
73
            elseif (carg .EQ. 'NBINSY') then
74
               READ(cvalue,*,IOSTAT=ios,ERR=200) NBINY
75
            elseif (carg .EQ. 'NBINSZ') then
76
               READ(cvalue,*,IOSTAT=ios,ERR=200) NBINZ
77
            end if
78
         end if
79
      end do
      I -----
      ! Something went wrong if we got here
80
      200
            CONTINUE
81
      WRITE(*,'(A,I5,2A)') 'Error reading line: ',iline,
                                                                               &
                           ' of file: ',ctrlfile(1:ilenctrl)
        &
      ilen = LEN_TRIM(cbuf)
82
      WRITE(*,'(A)') cbuf(1:ilen)
      ReadParams = 1
83
      goto 9999
84
      I ----
      ! End of file
85
      100 ReadParams = 0
      ! Close file
86
      CLOSE(iUnit)
      I-----
      ! Return
      9999 CONTINUE
87
88
      return
89
      END function ReadParams
```

C.5.4.6 Read OpenInventor 3D files (readiv.f90)

The following routine implements the elementary OpenInventor 3D file reader needed to be able to read Imaris' triangulated output (cf. chapter 2) according to the syntax standards given in [SGI (1992a)] and [SGI (1992b)]. The name of the OpenInventor file to be read is passed to this subroutine in the argument variable ivfile. The other argument success is the return value of type LOGICAL which contains .TRUE. if the file was successfully read, .FALSE. otherwise.

```
1 subroutine readiv(ivfile, success)
```

```
2 USE globals
```

```
3 IMPLICIT NONE
```

```
! DIPLOMA THESIS WS01/02 ICOS
                                                              ETH-ZUERICH
                              PROTEIN DIFFUSION INSIDE THE ENDOPLASMIC RETICULUM
               1-----
     ! Declaration of input/output variables
                                      _____
     ! Name of file to read
4
     CHARACTER(LEN=80), INTENT(IN)
                                             :: ivfile
     ! Read successful?
5
     LOGICAL, INTENT(OUT)
                                             :: success
     ! Declaration of local variables
     ! loop counters
6
     INTEGER
                                              :: i
     ! I/O error trap
     INTEGER
7
                                             :: istat
     ! read buffer
     CHARACTER(LEN=80)
8
                                             :: buf
     ! separator dummy
     INTEGER
9
                                             :: dummy
     ! Number of points
10
     INTEGER
                                             :: Npts
     ! Coordinates of vertices
     REAL(MK), DIMENSION(:,:), ALLOCATABLE
11
                                             :: pts
     ! vertex indices of current triangle
     INTEGER, DIMENSION(3)
12
                                             :: vertex
     success = .TRUE.
13
    istat = 0
buf(:) = ',
14
15
     1-----
     ! first pass through file to count number of points and triangles
     1----
     OPEN(80, FILE=ivfile, STATUS='OLD', ACTION='READ', IOSTAT=istat)
16
     if(istat .NE. 0) then
    WRITE(*,'(A)') 'Could not open file ', trim(ivfile)
17
18
       success = .FALSE.
19
20
       goto 900
21
     end if
     ! Skip file header
     do while(index(buf, 'point', .FALSE.) .EQ. 0)
22
23
       READ(80, '(A)') buf
24
     end do
     ! determine number of vertices
25
     Npts = 0
26
     do while(index(buf, '}', .FALSE.) .EQ. 0)
27
     Npts = Npts + 1
28
       READ(80, '(A)') buf
                                          ! read next line
29
     end do
     ! skip normals section
30
     do while(index(buf, 'coordIndex', .FALSE.) .EQ. 0)
       READ(80, '(A)') buf
31
32
     end do
     ! determine number of triangles
33
     M = 0
34
     do while(index(buf, '}', .FALSE.) .EQ. 0)
35
       if(index(buf, '-1', .FALSE.) .EQ. index(buf, '-1', .TRUE.)) then
36
         M = M + 1
37
       else
38
         M = M + 2
       end if
39
       READ(80, '(A)') buf
40
     end do
41
     WRITE(*,'(A,I6,A,I6,A)') 'Surface triangulation contains ',Npts, &
42
        ' vertices defining ',M,' triangles'
```

```
43
      CLOSE(80)
      1----
      ! allocate memory for points and triangles
      !---
                                                  _____
44
      istat = 0
      ! first index: x,y,z; second: point number; third: triangle number
45
      ALLOCATE(triang(3,3,M), c(3,M), normal(3,M), pts(3,Npts))
46
      if (istat .NE. 0) then
         WRITE(*,'(A)') 'Unable to allocate memory for triangles and vertives'
47
48
         success = .FALSE.
49
         goto 900
50
      end if
      1--
      ! second pass through file to read data
                                              _____
51
      OPEN(80, FILE=ivfile, STATUS='OLD', ACTION='READ', IOSTAT=istat)
      if(istat .NE. 0) then
WRITE(*,'(A)') 'Could not open file ', trim(ivfile)
52
53
54
         success = .FALSE.
         goto 900
55
      end if
56
      ! Skip file header
57
      do while(index(buf, 'point', .FALSE.) .EQ. 0)
        READ(80, '(A)') buf
58
59
      end do
60
      buf = buf(index(buf,'[')+1:)
      ! read point coordinates
61
      do i=1,Npts
                                           ! parse string for numbers
         READ(buf,*) pts(1:3,i)
62
         READ(80, '(A)') buf
63
                                                   ! read next line
64
      end do
      ! skip normals section
65
      do while(index(buf, 'coordIndex', .FALSE.) .EQ. 0)
        READ(80, '(A)') buf
66
67
      end do
      buf = buf(index(buf,'[')+1:)
68
69
      i = 1
70
      do while(i .LE. M)
         if(index(buf, '-1', .FALSE.) .EQ. index(buf, '-1', .TRUE.)) then
71
            ! only one triangle on this line
72
            READ(buf,*) vertex(1:3)
            triang(1:3,1,i) = pts(1:3,vertex(1))
triang(1:3,2,i) = pts(1:3,vertex(2))
73
74
75
            triang(1:3,3,i) = pts(1:3,vertex(3))
76
            i = i + 1
77
         else
                                                    ! two triangles on this line
78
            READ(buf,*) vertex(1:3)
            ! +1 because .iv file has C style ordering starting at zero
            triang(1:3,1,i) = pts(1:3,vertex(1)+1)
triang(1:3,2,i) = pts(1:3,vertex(2)+1)
79
80
            triang(1:3,3,i) = pts(1:3,vertex(3)+1)
81
82
            i = i + 1
83
            buf = buf(index(buf,'-1')+3:)
84
            READ(buf,*) vertex(1:3)
85
            triang(1:3,1,i) = pts(1:3,vertex(1)+1)
86
            triang(1:3,2,i) = pts(1:3,vertex(2)+1)
87
            triang(1:3,3,i) = pts(1:3,vertex(3)+1)
88
            i = i + 1
89
         end if
90
         READ(80, '(A)') buf
                                                    ! read next line
91
      end do
92
      CLOSE(80)
93
      WRITE(*,'(A)') 'Triangles ... read.'
      WRITE(*,'(A)')
94
95
      900 if(ALLOCATED(pts)) DEALLOCATE(pts)
96
      return
      END subroutine readiv
97
```

C.5.4.7Check validity of triangulation (chktriang.f90)

The following file contains the two function chk_topology and chk_orientation. They check the triangulation for its topological and syntactical validity, respectively, according to the criteria and algorithms given in chapter 3. Both functions return a value of type LOGICAL which is .TRUE. if all criteria are fulfilled and .FALSE. if one of them is not. Notice that chk_orientation needs the centroids and normals of all the triangles and cannot be called before they have been calculated in the main program (statements 231 to 240). Moreover, it expects the triangulation to be at least topologically valid and should thus always be called *after* chk_topology.

- function chk_topology() 1
- 2 USE globals

4

5

6

7

8

9

3 IMPLICIT NONE

```
_____
     !chk_topology Checks topological validity of a triangulation set
         CHK_TOPOLOGY checks the following requirements:
            A) All edges must be of length > 0
           B) No vertex must be inside another triangle
           C) Every triangle must have exactly three neighbors sharing 2 vertices
           D) No edge of any triangle must intersect any other triangle
         if all points are met, the triangulation describes a closed coherent
         surface in space. Attn.: this is o(M**2)!
         RETURN VALUE: .TRUE. if set is topologically valid, else .FALSE.
         See also INIT_PART
         todo:
        DIPLOMA THESIS WS01/02 ICOS
                                                                 ETH-ZUERICH
                   PROTEIN DIFFUSION INSIDE THE ENDOPLASMIC RETICULUM
      ! Declaration of input arguments and return value
     ! return value: .TRUE. if set is topologically valid
     LOGICAL
                                             :: chk_topology
     ! Declaration of external functions
     ! intersects a line with a triangle
     REAL(MK), EXTERNAL
                                             :: intersect
     ! Declaration of local variables
     1 ---
     ! loop counters
     INTEGER
                                             :: i ,j, k, l
     ! edge vectors
     REAL(MK), DIMENSION(3)
                                             :: a, b, g
     ! determinant of a matrix
     REAL(MK)
                                             :: det
     ! Number of common vertices of two triangles
     INTEGER
                                             :: Nvert
     ! Number of neighbors of a triangle
10
     INTEGER
                                             :: Neigh
     ! Matrix of a 2x2 system
11
     REAL(MK), DIMENSION(2,2)
                                             :: S
     ! line coefficients
12
     REAL(MK)
                                             :: alpha, beta
     ! local condition variable
```

APPENDIX C. THE SIMULATION CODES

```
13
      LOGICAL
                                                      :: lcond
14
      chk_topology = .TRUE.
      \texttt{WRITE}(\texttt{*},\texttt{'}(\texttt{A})\texttt{'}) 'Checking topology of triangulation set \ldots '
15
16
      WRITE(*,'(A,I6,A)') 'Set contains ', M, ' triangles.'
      ! a) All edges must be > 0
      !--
17
      lcond = .TRUE.
      ! Main loop over all triangles
18
      do i=1.M
19
         if(sum((triang(:,2,i)-triang(:,1,i))**2) .LT. TOL) then
20
             WRITE(*,'(A,I6,A,E12.2)') '*** WARNING: Edge between vertices 1 &
                and 2 of triangle ',i,' is < ', TOL
21
             chk_topology = .FALSE.
22
             lcond = .FALSE.
23
          end if
24
         if(sum((triang(:,3,i)-triang(:,1,i))**2) .LT. TOL) then
             WRITE(*,'(A,I6,A,E12.2)') '*** WARNING: Edge between vertices 1 &
25
                and 3 of triangle ',i,' is < ', TOL
26
             chk_topology = .FALSE.
27
            lcond = .FALSE.
28
         end if
29
         if(sum((triang(:,3,i)-triang(:,2,i))**2) .LT. TOL) then
             WRITE(*,'(A,I6,A,E12.2)') '*** WARNING: Edge between vertices 2 &
30
               and 3 of triangle ',i,' is < ', TOL
31
             chk_topology = .FALSE.
32
            lcond = .FALSE.
33
         end if
34
      end do
35
      if(lcond) then
         WRITE(*,'(A)') 'Topology[A] Good: All edges are of length > 0.'
36
37
      end if
      ! b) No vertex must be inside any other triangle
      1----
                       _____
                                                              _____
      1 \text{ cond} = .TRUE.
38
      ! Main loop over all triangles
39
      do i=1,M
40
         do j=1,M
                                       ! loop over all the other triangles
41
            if(j .NE. j) then
                                       ! loop over all vertices of triangle i
42
                do k=1,3
43
                   a = triang(:,2,j)-triang(:,1,j)
                   b = triang(:,3,j)-triang(:,1,j)
det = a(1)*b(2)-a(2)*b(1)
44
45
46
                   if(abs(det) .GE. TOL) then
47
                      det = 1.0/det
                                               ! 2x2 subsystem
48
                      S(1,1)= b(2)*det
49
                      S(1,2) = -b(1) * det
50
                      S(2,1)=-a(2)*det
51
                      S(2,2)= a(1)*det
                      g(1:2)= triang(1:2,k,i)-triang(1:2,1,j) 
! Solution
52
53
                      alpha = S(1,1)*g(1)+S(1,2)*g(2)
54
                      beta = S(2,1)*g(1)+S(2,2)*g(2)
                       ! point is in same plane if third equation is also
                        fullfilled with this solution
                      if(alpha*a(3)+beta*b(3)+triang(3,1,j)-triang(3,k,i) .LE. &
55
                            TOL) then
                          if((alpha+beta .LT. 1.0+TOL) .AND. (alpha .GE. TOL)
56
                             .AND. (beta .GE. TOL)) then ! point is insde triangle
WRITE(*,'(A,I1,A,I6,A,I6)') '*** WARNING: Vertex ',k &
57
                                ,' of triangle ',i,' lies inside triangle ',j
58
                             chk_topology = .FALSE.
59
                             lcond = .FALSE.
60
                          end if
61
      9
                      end if
62
                   end if
63
                end do
            end if
64
65
         end do
66
      end do
      if(lcond) then
67
```

68 WRITE(*, (A)) 'Topology[B] Good: No vertex is inside any other triangle.' 69 end if

```
1 ----
                    _____
      ! c) Every triangle must have exactly three neighbors sharing 2 vertices
      !--
70
      lcond = .TRUE.
      ! Main loop over all triangles
71
      do i=1,M
         Neigh = 0
72
                               ! reset neighbor counter
73
         do j=1,M
                              ! loop over all the other triangles
74
            if(j .NE. i) then
75
               Nvert = 0 ! reset number of common vertices
                             ! loop over all vertices of triangle i
! loop over all different vertices of triangle j
               do k=1,3
76
77
                  do l=1,3
78
                     if(sum((triang(:,k,i)-triang(:,1,j))**2) .LE. TOL) then
79
                        Nvert=Nvert+1
                                         ! one common vertex found
                     end if
80
                  end do
81
82
               end do
83
               if(Nvert .GT. 2) then
                  WRITE(*,'(A,I6,A,I6,A,I1,A)') '*** WARNING: Triangles ',i, &
84
                  'and ',j,'have ',Nvert,' vertices in common!
chk_topology = .FALSE.
85
                  lcond = .FALSE.
86
87
               end if
88
               if(Nvert .EQ. 2) then
                                       ! increment neighbor counter
89
                  Neigh=Neigh+1
               end if
90
91
            end if
         end do
92
         if(Neigh .NE. 3) then
93
            WRITE(*,'(A,I6,A,I1,A)') '*** WARNING: Triangle ',i,' has ' &
94
               ,Neigh,' neighbors instead of 3.'
            chk_topology = .FALSE.
95
96
            lcond = .FALSE.
97
         end if
      end do
98
      if(lcond) then
99
        WRITE(*,'(A)') 'Topology[C] Good: All triangles have 3 neighbors.'
100
      end if
101
                                                                   _____
      ! d) No edge of any triangle must intersect any other triangle
      !-----
102
     lcond = .TRUE.
      ! Main loop over all triangles
103
      do i=1,M
104
        a = triang(:,2,i)-triang(:,1,i)
105
         b = triang(:,3,i)-triang(:,1,i)
106
         g = triang(:,3,i)-triang(:,2,i)
107
         do j=1,M
            if(j .NE. i) then
108
109
               if(intersect(triang(:,1,i), a, j) .GE. 0.0) then
                  WRITE(*,'(A,I6,A,I6)') '*** WARNING: Edge between vertices 1 &
110
                     and 2 of triangle ',i,' intersects with triangle ',j
111
                  chk_topology = .FALSE.
112
                  lcond = .FALSE.
113
               end if
               if(intersect(triang(:,1,i), b, j) .GE. 0.0) then
WRITE(*,'(A,I6,A,I6)') '*** WARNING: Edge between vertices 1 &
114
115
                     and 2 of triangle ',i,' intersects with triangle ',j
116
                  chk_topology = .FALSE.
                  lcond = .FALSE.
117
118
               end if
               if(intersect(triang(:,2,i), g, j) .GE. 0.0) then
WRITE(*,'(A,I6,A,I6)') '*** WARNING: Edge between vertices 1 &
119
120
                     and 2 of triangle ',i,' intersects with triangle ',j
121
                  chk_topology = .FALSE.
                  lcond = .FALSE.
122
123
               end if
            end if
124
125
         end do
      end do
126
127
      if(lcond) then
         WRITE(*,'(A)') 'Topology[D] Good: No edge intersects any other triangle.'
128
```

```
APPENDIX C. THE SIMULATION CODES
180
129
    end if
130
    return
131 END function chk_topology
    1------
132
    function chk_orientation()
133
    USE globals
    IMPLICIT NONE
134
    !chk_orientation Checks syntactical validity of a triangulation set
       CHK_ORIENTATION checks whether all normals point out of the domain.
       Needs centroids and normals to be calculated first!
       Assumes triangulation to be topologically valid.
       RETURN VALUE: .TRUE. if set is syntactically valid, else .FALSE.
       See also INIT_PART
       todo:
    ! DIPLOMA THESIS WS01/02 ICOS
                                                      ETH-ZUERICH
               PROTEIN DIFFUSION INSIDE THE ENDOPLASMIC RETICULUM
    ! Declaration of external functions
                   _____
    1--
    ! checks whether a point is inside the domain (needs topological validity !)
135
   LOGICAL, EXTERNAL
                                     :: point_in_domain
    1-----
    ! Declaration of input arguments and return value
    1--
    ! return value: .TRUE. if set is syntactically valid
136
   LOGICAL
                                      :: chk_orientation
    I -----
                  _____
    ! Declaration of local variables
    1------
    ! loop counters
137
    INTEGER
                                      :: i
    ! point in question
138
    REAL(MK), DIMENSION(3)
                                     :: pt
139
    chk_orientation = .TRUE.
     ! main loop over all triangles
140
    do i=1,M
      ! follow the normal for a small distance (but larger than TOL
      ! for the difference to be detected by the comparison operator
141
      pt = c(:,i)+5*TOL*normal(:,i)
      if(point_in_domain(pt)) then
142
         WRITE(*,'(A,I6,A)') '*** WARNING: Triangle ',i,' is not properly oriented.'
143
         chk_orientation = .FALSE.
144
      end if
145
    end do
146
147
    if(chk_orientation) then
148
      WRITE(*,'(A)') 'Orientation Good: All normals point outward.'
149
    end if
150
    return
151
    END function chk_orientation
```

C.5.4.8 Intersect a line with a triangle (intersect.f90)

intersect implements the algorithm given in section 3.2 to intersect a line $P + \lambda v$ with a triangle given by its index t in the current triangulation set. The function takes the two vectors $P \in \mathbb{R}^3$ and $v \in \mathbb{R}^3$ as well as the integer triangle index t as input arguments and returns the absolute value of the line parameter λ (lambda) at the point of intersection or -1 if no intersection point *inside* the triangle exists.

```
function intersect(P, v, t)
1
```

```
USE globals
2
```

```
IMPLICIT NONE
3
```

```
!INTERSECT Intersects line P+lambda*v with triangle t
         Returns the absolute value of lambda.
         If no intersection point inside triangle t exists, -1 is returned.
         See also INIT_PART, CHK_TOPOLOGY, POINT_IN_DOMAIN
                  _____
      ! DIPLOMA THESIS WS01/02 ICOS
                                                                      ETH-ZUERICH
                    PROTEIN DIFFUSION INSIDE THE ENDOPLASMIC RETICULUM
              ----- ivo f. sbalzarini ------
      ! Declaration of input and output variables
      ! Starting point of line
4
     REAL(MK), DIMENSION(3), INTENT(IN) :: P
      ! direction vector
     REAL(MK), DIMENSION(3), INTENT(IN) :: v
5
     ! triangle number
     INTEGER, INTENT(IN)
6
                                        :: t
      ! return value
     REAL(MK)
                                       :: intersect
                                _____
      ! Declaration of local variables
     ! line parameters
8
     REAL(MK)
                                       :: alpha, beta, lambda
     ! edge vectors of triangle t
9
     REAL(MK), DIMENSION(3)
                                       :: a, b
      ! System matrix
     REAL(MK), DIMENSION(3,3)
10
                                       :: S
      ! determinant of system matrix and its inverse
     REAL(MK)
11
                                       :: det, detinv
     intersect = -1.0
12
     ! edge vectors
13
     a = triang(:,2,t)-triang(:,1,t)
     b = triang(:,3,t)-triang(:,1,t)
14
     ! determinant of system matrix
     det = -a(3)*b(2)*v(1)+a(2)*b(3)*v(1)+a(3)*b(1)*v(2)-
15
                                                           &
          a(1)*b(3)*v(2)-a(2)*b(1)*v(3)+a(1)*b(2)*v(3)
     if(abs(det) .GE. TOL) then
16
                                    ! intersection point exists
        ! inverse of determinant
17
        detinv=1.0/det
         ! inverse of system matrix
18
        S(1,1)=(b(2)*v(3)-b(3)*v(2))*detinv
19
        S(1,2)=(b(3)*v(1)-b(1)*v(3))*detinv
20
        S(1,3)=(b(1)*v(2)-b(2)*v(1))*detinv
21
        S(2,1)=(a(3)*v(2)-a(2)*v(3))*detinv
22
        S(2,2)=(a(1)*v(3)-a(3)*v(1))*detinv
23
        S(2,3)=(a(2)*v(1)-a(1)*v(2))*detinv
24
        S(3,1)=(a(2)*b(3)-a(3)*b(2))*detinv
25
        S(3,2)=(a(3)*b(1)-a(1)*b(3))*detinv
26
        S(3,3)=(a(1)*b(2)-a(2)*b(1))*detinv
        ! solution of linear system
```

```
27
         a(1:3)= P(1:3)-triang(:,1,t)
28
         alpha = S(1,1)*a(1)+S(1,2)*a(2)+S(1,3)*a(3)
29
         beta = S(2,1)*a(1)+S(2,2)*a(2)+S(2,3)*a(3)
30
         lambda= S(3,1)*a(1)+S(3,2)*a(2)+S(3,3)*a(3)
         ! check whether intersection point is inside triangle t
31
         if((alpha+beta .LE. 1.0) .AND. (alpha .GE. 0)
                                                         &
              .AND. (beta .GE. 0)) then
32
            intersect = abs(-lambda)
33
         end if
34
      end if
35
      return
```

36 END function intersect

C.5.4.9 Create bin lists and cell lists (SortT.f90)

This subroutine sets up the bin lists and cell lists as described in section 7.2.5 and sorts the triangles into them. The data structures needed are explained in section 7.2.6.

```
1
      subroutine SortT
2
     USE globals
3
     USE Util
4
     IMPLICIT NONE
      1-----
      SORTT Sorts triangles and builds lists.
         SortT sets up the bin lists for point_in_domain by assigning the
         triangles to bins of equal y and z coordinates as well as the cell
         lists for the nearest triangle search (if needed).
         See also INIT_PART, POINT_IN_DOMAIN, ALLOCATELL
        DIPLOMA THESIS WS01/02 ICOS
                                                                     ETH-ZUERICH
                    PROTEIN DIFFUSION INSIDE THE ENDOPLASMIC RETICULUM
                      ! Declaration of local variables
      ! loop counters
5
     INTEGER
                                             :: i, j, k, l, p, ix1, ix2
      ! minimum and maximum indices in x, y and z
     INTEGER
6
                                            :: minbinx, maxbinx, minbinv
                                             :: maxbiny, minbinz, maxbinz
     INTEGER
7
      ! x, y and z indices of current bin/cell
8
     INTEGER
                                             :: idxx, idxy, idxz
      ! lower and upper boundaries of bin
     REAL(MK), DIMENSION(4)
                                             :: bb
9
      ! error trap
     INTEGER
                                             :: istat
10
      ! edge vectors of triangle projections, vertex 1
     REAL(MK), DIMENSION(2)
                                             :: a, b, x1, y1, v1, v2
11
      ! three dimensional vectors for cell list
     REAL(MK), DIMENSION(3)
                                             :: x13, a3, b3, n3, y13
12
      ! determinants
13
     REAL(MK)
                                             :: det, det2
      ! line parameters
14
     REAL(MK)
                                             :: alpha, beta, lambda
      ! system matrix
     REAL(MK), DIMENSION(3,3)
15
                                             :: mat
      ! lower and upper boundaries of cell and its centroid
16
     REAL(MK), DIMENSION(3)
                                             :: bbl, bbu, bbc
      ! flag whether a certain triangle already belongs to a bin or not
17
     LOGICAL
                                             :: NotIn
      ! all three edge vectors of a triangle
18
     REAL(MK), DIMENSION(2,3)
                                             :: edge
```

! all edge vectors of a bin 19 REAL(MK), DIMENSION(2,2) :: side _____ 1_____ ! Initialize !----20 istat = 0! initialize lists and arrays 21 CALL AllocateLL(0,0,0,0,istat) 22 if(istat .NE. 0) then 23 WRITE(*,'(A)') 'Error in AllocateLL: returning from SortT.' 24 return 25 end if 26 dby = (ymax-ymin)/real(NBINY) 27 dbz = (zmax-zmin)/real(NBINZ) ! Assign triangles to linear bins for point_in_domain 28 do i=1,M ! loop over all triangles minbiny = NBINY+1 29 30 maxbiny = 0 minbinz = NBINZ+1 31 maxbinz = 032 33 x1 = triang(2:3,1,i) 34 edge(:,1) = triang(2:3,2,i)-x1 edge(:,2) = triang(2:3,3,i)-x1 35 edge(:,3) = triang(2:3,3,i)-triang(2:3,2,i) 36 37 a = edge(:,1) b = edge(:,2) 38 det = a(1)*b(2)-a(2)*b(1)39 det = 1.0_MK/det 40 do j=1,3 ! loop over all vertices 41 ! determine index of bin of this vertex 42 idxy = ceiling((triang(2,j,i)-ymin)/dby) if(idxy .EQ. 0) idxy = 1
if(idxy .EQ. NBINY+1) idxy = NBINY ! include boundaries at walls 43 44 45 idxz = ceiling((triang(3,j,i)-zmin)/dbz) if(idxz .EQ. 0) idxz = 146 ! include boundaries at walls if(idxz .EQ. NBINZ+1) idxz = NBINZ 47 ! add this triangle to list of this bin if not already there 48 if(.NOT. ASSOCIATED(bin(idxy,idxz)%list)) then 49 CALL AllocateLL(1,0,idxy,idxz,istat) end if 50 51 NotIn = .FALSE. if(nbin(idxy,idxz) .EQ. 0) then 52 53 NotIn = .TRUE. elseif(bin(idxy,idxz)%list(nbin(idxy,idxz)) .NE. i) then 54 55 NotIn = .TRUE. 56 end if 57 if(NotIn) then 58 p = nbin(idxy,idxz)+1 59 nbin(idxy,idxz) = p ! enlarge this list if nedded 60 if(p .GT. size(bin(idxy,idxz)%list)) then 61 CALL AllocateLL(1,0,idxy,idxz,istat) 62 end if bin(idxy,idxz)%list(p) = i 63 ! update bounding box of cells 64 if(idxy .LT. minbiny) minbiny = idxy 65 if(idxy .GT. maxbiny) maxbiny = idxy if(idxz .LT. minbinz) minbinz = idxz 66 67 if(idxz .GT. maxbinz) maxbinz = idxz end if 68 69 end do ! loop over all bins in the bounding box of the vertices 70 do j=minbiny,maxbiny bb(1) = ymin+(j-1)*dby ! lower y 71 72 bb(2) = ymin+j*dby ! upper y 73 do k=minbinz,maxbinz ! if no list for this bin is yet associated, create one 74 if(.NOT. ASSOCIATED(bin(j,k)%list)) then 75 CALL AllocateLL(1,0,j,k,istat) 76 end if 77 NotIn = .FALSE. ! if this triangle doesnt already belong to this bin, check it 78 if(nbin(j,k) .EQ. 0) then NotIn = .TRUE. 79

80 elseif(bin(j,k)%list(nbin(j,k)) .NE. i) then 81 NotIn = .TRUE. 82 end if 83 if(NotIn) then 84 bb(3) = zmin+(k-1)*dbz ! lower z 85 bb(4) = zmin+k*dbz ! upper z ! edge vectors of this bin (projected onto y/z plane) ! only store two since the other two are identical 86 side(1,1) = bb(2)-bb(1)87 side(2,1) = 0.0_MK 88 side(1,2) = 0.0_MK 89 side(2,2) = bb(4)-bb(3)90 do ix1=1,2 91 y1(1) = bb(ix1)-x1(1)92 do ix2=3,4 93 y1(2) = bb(ix2)-x1(2)94 alpha = det*(b(2)*y1(1)-b(1)*y1(2)) 95 beta = det*(-a(2)*y1(1)+a(1)*y1(2)) if((alpha .GE. 0.0_MK) .AND. (beta .GE. 0.0_MK) & 96 .AND. (alpha+beta .LE. 1.0_MK)) then ! point is inside triangle: add triangle to this bin 97 p = nbin(j,k)+198 nbin(j,k) = p ! enlarge this list if nedded if(p .GT. size(bin(j,k)%list)) then 99 CALL AllocateLL(1,0,j,k,istat) 100 101 end if 102 bin(j,k)%list(p) = i 103 goto 250 ! next bin 104 end if 105 end do 106 end do ! if no vertex of the bin is inside the triangle, check whether any edge of the triangle intersects any edge of ! the bin 107 do ix1=1.3 ! loop over triangles edges v1 = edge(:,ix1) 108 109 do ix2=1.4 ! loop over bins edges ! choose adge vector mod 2 since bin is axis parallel v2 = -side(:,mod(ix2,2)+1) 110 ! system determinant (do not overwrite det !!) $det^{2} = v1(1)*v2(2)-v1(2)*v2(1)$ 111 det2 = 1.0_MK/det2 112 ! determine rhs (starting points of lines depending ! on what edge we are at 113 if(ix2 .EQ. 1) then 114 y1 = (/bb(1), bb(3)/)elseif(ix2 .EQ. 2) then 115 116 y1 = (/bb(1), bb(3)/)117 elseif(ix2 .EQ. 3) then 118 y1 = (/bb(2), bb(3)/)119 elseif(ix2 .EQ. 4) then 120 y1 = (/bb(1), bb(4)/)121 endif 122 if(ix1 .EQ. 1 .OR. ix1 .EQ. 2) then y1 = y1 - x1123 124 else 125 y1 = y1 - triang(2:3,2,i) 126 end if ! solve for intersection point 127 alpha = det2*(v2(2)*y1(1)-v2(1)*y1(2)) 128 beta = det2*(-v1(2)*y1(1)+v1(1)*y1(2)) ! if intersected, add triangle to this bin if((alpha .GE. 0.0_MK) .AND. (beta .GE. 0.0_MK) .AND. & 129 (alpha .LE. 1.0_MK) .AND. (beta .LE. 1.0_MK)) then ! point is inside triangle: add triangle to this bin 130 p = nbin(j,k)+1131 nbin(j,k) = p ! enlarge this list if nedded if(p .GT. size(bin(j,k)%list)) then 132 133 CALL AllocateLL(1,0,j,k,istat) end if 134 135 bin(j,k)%list(p) = i 136 goto 250 ! next bin 137 end if 138 end do 139 end do 140 250 end if 141 end do 142 end do

```
143
     end do
      1-----
                                          _____
      ! prune bin lists to actual size to save memory
                                                    _____
      1 -
              _____
144
     do i=1,NBINY
145
        do j=1,NBINZ
146
           bin(i,j)%list=>reallocate(bin(i,j)%list,nbin(i,j))
147
         end do
      end do
148
      ! Assign triangles to cell lists for nearest neighbor search
      ! cell lists only needed for boundary condition handling
149
     if(BC .GT. 0) then
        do i=1,M
150
                            ! loop over all triangles
           minbinx = NCELLX+1
151
           maxbinx = 0
152
           minbiny = NCELLY+1
153
154
           maxbiny = 0
155
           minbinz = NCELLZ+1
           maxbinz = 0
156
           x13 = triang(1:3,1,i)
157
           a3 = triang(1:3,2,i)-x13
158
           b3 = triang(1:3,3,i)-x13
159
           n3 = normal(1:3,i)
160
           det = -a3(3)*b3(2)*n3(1)+a3(2)*b3(3)*n3(1)+a3(3)*b3(1)*n3(2)- &
161
               a3(1)*b3(3)*n3(2)-a3(2)*b3(1)*n3(3)+a3(1)*b3(2)*n3(3)
           det = 1.0_MK/det
162
           mat(1,1) = -b3(3)*n3(2)+b3(2)*n3(3)
163
           mat(1,2) = b3(3)*n3(1)-b3(1)*n3(3)
164
           mat(1,3) = -b3(2)*n3(1)+b3(1)*n3(2)
165
           mat(2,1) = a3(3)*n3(2)-a3(2)*n3(3)
166
           mat(2,2) = -a3(3)*n3(1)+a3(1)*n3(3)
167
           mat(2,3) = a3(2)*n3(1)-a3(1)*n3(2)
mat(3,1) = -a3(3)*b3(2)+a3(2)*b3(3)
168
169
           mat(3,2) = a3(3)*b3(1)-a3(1)*b3(3)
170
           mat(3,3) = -a3(2)*b3(1)+a3(1)*b3(2)
171
172
           do j=1,3
                            ! loop over all vertices
              ! determine index of cell of this vertex
              idxx = ceiling((triang(1,j,i)-xmin)/dcx)
173
174
              if(idxx .EQ. 0) idxx = 1
                                                     ! include boundaries at walls
              if(idxx .EQ. NCELLX+1) idxx = NCELLX
175
176
              idxy = ceiling((triang(2,j,i)-ymin)/dcy)
177
              if(idxy .EQ. 0) idxy = 1
                                                     ! include boundaries at walls
178
              if(idxy .EQ. NCELLY+1) idxy = NCELLY
179
              idxz = ceiling((triang(3,j,i)-zmin)/dcz)
180
              if(idxz .EQ. 0) idxz = 1
                                                     ! include boundaries at walls
181
              if(idxz .EQ. NCELLZ+1) idxz = NCELLZ
182
              if(.NOT. ASSOCIATED(cell(idxx,idxy,idxz)%list)) then
183
                 CALL AllocateLL(1,idxx,idxy,idxz,istat)
184
              end if
185
              NotIn = .FALSE.
186
              if(ncell(idxx,idxy,idxz) .EQ. 0) then
187
                 NotIn = .TRUE.
188
              elseif(cell(idxx,idxy,idxz)%list(ncell(idxx,idxy,idxz)) .NE. i) then
189
                 NotIn = .TRUE.
190
               end if
191
              if(NotIn) then
192
                 p = ncell(idxx,idxy,idxz)+1
193
                 ncell(idxx,idxy,idxz) = p
                  ! enlarge this list if nedded
                 if(p .GT. size(cell(idxx,idxy,idxz)%list)) then
194
195
                    CALL AllocateLL(1,idxx,idxy,idxz,istat)
196
                 end if
197
                 cell(idxx,idxy,idxz)%list(p) = i
                 ! update bounding box of cells
198
                 if(idxx .LT. minbinx) minbinx = idxx
                 if(idxx .GT. maxbinx) maxbinx = idxx
199
200
                 if(idxy .LT. minbiny) minbiny = idxy
201
                 if(idxy .GT. maxbiny) maxbiny = idxy
202
                 if(idxz .LT. minbinz) minbinz = idxz
203
                 if(idxz .GT. maxbinz) maxbinz = idxz
204
              end if
205
           end do
           ! loop over all cells in the bounding box of the vertices
206
           do j=minbinx,maxbinx
```

207	<pre>bbl(1) = xmin+(j-1)*dcx ! lower x</pre>
208	bbu(1) = xmin+j*dcx ! upper x
209	do k=minbiny,maxbiny
210	bbl(2) = vmin+(k-1)*dcy + lower y
211	bbu(2) = vmin+k*dcv ! upper v
212	do l=minbinz.maxbinz
	if no list for this cell is yet associated, create one
213	if NOT ASSOCIATED (calli k 1)(list)) than
210	CALL Allocatell(1 k l istat)
211	and if
215	
210	NOTIN = .FALSE.
	! II this triangle doesn't aiready belong to this cell, check it
217	if(ncell(j,k,l) .EQ. 0) then
218	Notin = .TRUE.
219	elseif(cell(j,k,l)%list(ncell(j,k,l)) .NE. i) then
220	NotIn = .TRUE.
221	end if
222	if(NotIn) then
223	bbl(3) = zmin+(1-1)*dcz + lower z
224	bbu(3) = zmin+l*dcz ! upper z
225	bbc(1:3) = 0.5*(bbl(1:3)+bbu(1:3)) ! voxels centroid
226	$y_{13}(1:3) = bbc(1:3) - x_{13}(1:3)$
227	alpha = det*(mat(1,1)*y13(1)+mat(1,2)*y13(2) + &
	mat(1,3)*v13(3))
228	heta = det*(mat(2,1)*v13(1)+mat(2,2)*v13(2) + &
	mat(2,3)*v13(3))
229	lambda = det*(mat(3, 1)*v13(1)+mat(3, 2)*v13(2) + &
220	mat(3,3)*v13(3))
	l if intersection point is inside triangle
230	if(alpha GE 0.0 MK) AND (beta GE 0.0 MK) &
200	AND (alpha+ta LE 1 0 MK)) then
	I point where normal through centroid intersects triangle
231	! point where normal through centroid intersects triangle $u_{13}(1,3) = bbc_{11}(1,3)+lambdatn3(1,3)$
231	<pre>! point where normal through centroid intersects triangle y13(1:3) = bbc(1:3)+lambda*n3(1:3) ! if this point is inside the cell add triangle to cell</pre>
231	<pre>! point where normal through centroid intersects triangle y13(1:3) = bbc(1:3)+lambda*n3(1:3) ! if this point is inside the cell, add triangle to cell if((r13(1) LT bbr(1)) AND (r13(2) LT bbr(2)) AND &</pre>
231 232	<pre>! point where normal through centroid intersects triangle y13(1:3) = bbc(1:3)+lambda*n3(1:3) ! if this point is inside the cell, add triangle to cell if((y13(1) .LT. bbu(1)) .AND. (y13(2) .LT. bbu(2)) .AND. &</pre>
231 232	<pre>! point where normal through centroid intersects triangle y13(1:3) = bbc(1:3)+lambda*n3(1:3) ! if this point is inside the cell, add triangle to cell if((y13(1) .LT. bbu(1)) .AND. (y13(2) .LT. bbu(2)) .AND. & (y13(3) .LT. bbu(3)) .AND. (y13(1) .GE. bbl(1)) .AND. & (y13(2) .GE bbl(2)) .AND. (y13(3) .GE bbl(3)) then</pre>
231 232 233	<pre>! point where normal through centroid intersects triangle y13(1:3) = bbc(1:3)+lambda*n3(1:3) ! if this point is inside the cell, add triangle to cell if((y13(1) .LT. bbu(1)) .AND. (y13(2) .LT. bbu(2)) .AND. & (y13(3) .LT. bbu(3)) .AND. (y13(1) .GE. bbl(1)) .AND. & (y13(2) .GE. bbl(2)) .AND. (y13(3) .GE. bbl(3))) then</pre>
231 232 233 234	<pre>! point where normal through centroid intersects triangle y13(1:3) = bbc(1:3)+lambda*n3(1:3) ! if this point is inside the cell, add triangle to cell if((y13(1) .LT. bbu(1)) .AND. (y13(2) .LT. bbu(2)) .AND. & (y13(3) .LT. bbu(3)) .AND. (y13(1) .GE. bbl(1)) .AND. & (y13(2) .GE. bbl(2)) .AND. (y13(3) .GE. bbl(3))) then p = ncell(j,k,l)+1 proclicit k l) = p</pre>
231 232 233 234	<pre>! point where normal through centroid intersects triangle y13(1:3) = bbc(1:3)+lambda*n3(1:3) ! if this point is inside the cell, add triangle to cell if((y13(1) .LT. bbu(1)) .AND. (y13(2) .LT. bbu(2)) .AND. & (y13(3) .LT. bbu(3)) .AND. (y13(1) .GE. bbl(1)) .AND. & (y13(2) .GE. bbl(2)) .AND. (y13(3) .GE. bbl(3))) then p = ncell(j,k,1)+1 ncell(j,k,1) = p </pre>
231 232 233 234	<pre>! point where normal through centroid intersects triangle y13(1:3) = bbc(1:3)+lambda*n3(1:3) ! if this point is inside the cell, add triangle to cell if((y13(1) .LT. bbu(1)) .AND. (y13(2) .LT. bbu(2)) .AND. & (y13(3) .LT. bbu(3)) .AND. (y13(1) .GE. bbl(1)) .AND. & (y13(2) .GE. bbl(2)) .AND. (y13(3) .GE. bbl(3))) then p = ncell(j,k,l)+1 ncell(j,k,l) = p ! enlarge this list if nedded if (contact is a start of the balance is a s</pre>
231 232 233 234 235	<pre>! point where normal through centroid intersects triangle y13(1:3) = bbc(1:3)+lambda*n3(1:3) ! if this point is inside the cell, add triangle to cell if((y13(1) .LT. bbu(1)) .AND. (y13(2) .LT. bbu(2)) .AND. & (y13(3) .LT. bbu(3)) .AND. (y13(1) .GE. bbl(1)) .AND. & (y13(2) .GE. bbl(2)) .AND. (y13(3) .GE. bbl(3))) then p = ncell(j,k,l)+1 ncell(j,k,l) = p ! enlarge this list if nedded if(p .GT. size(cell(j,k,l)%list)) then</pre>
231 232 233 234 235 236	<pre>! point where normal through centroid intersects triangle y13(1:3) = bbc(1:3)+lambda*n3(1:3) ! if this point is inside the cell, add triangle to cell if((y13(1) .LT. bbu(1)) .AND. (y13(2) .LT. bbu(2)) .AND. & (y13(3) .LT. bbu(3)) .AND. (y13(1) .GE. bbl(1)) .AND. & (y13(2) .GE. bbl(2)) .AND. (y13(3) .GE. bbl(3))) then p = ncell(j,k,1)+1 ncell(j,k,1) = p ! enlarge this list if nedded if(p .GT. size(cell(j,k,1)%list)) then CALL AllocateLL(1,j,k,1,istat)</pre>
231 232 233 234 235 236 237	<pre>! point where normal through centroid intersects triangle y13(1:3) = bbc(1:3)+lambda*n3(1:3) ! if this point is inside the cell, add triangle to cell if((y13(1) .LT. bbu(1)) .AND. (y13(2) .LT. bbu(2)) .AND. & (y13(3) .LT. bbu(3)) .AND. (y13(1) .GE. bbl(1)) .AND. & (y13(2) .GE. bbl(2)) .AND. (y13(3) .GE. bbl(3))) then p = ncell(j,k,1)+1 ncell(j,k,1) = p ! enlarge this list if nedded if(p .GT. size(cell(j,k,1)%list)) then CALL AllocateLL(1,j,k,1,istat) end if</pre>
 231 232 233 234 235 236 237 238 	<pre>! point where normal through centroid intersects triangle y13(1:3) = bbc(1:3)+lambda*n3(1:3) ! if this point is inside the cell, add triangle to cell if((y13(1) .LT. bbu(1)) .AND. (y13(2) .LT. bbu(2)) .AND. & (y13(3) .LT. bbu(3)) .AND. (y13(1) .GE. bbl(1)) .AND. & (y13(2) .GE. bbl(2)) .AND. (y13(3) .GE. bbl(3))) then p = ncell(j,k,l) = p ! enlarge this list if nedded if(p .GT. size(cell(j,k,l)%list)) then CALL AllocateLL(1,j,k,l,istat) end if cell(j,k,l)%list(p) = i</pre>
231 232 233 234 235 236 237 238 239	<pre>! point where normal through centroid intersects triangle y13(1:3) = bbc(1:3)+lambda*n3(1:3) ! if this point is inside the cell, add triangle to cell if((y13(1) .LT. bbu(1)) .AND. (y13(2) .LT. bbu(2)) .AND. & (y13(3) .LT. bbu(3)) .AND. (y13(1) .GE. bbl(1)) .AND. & (y13(2) .GE. bbl(2)) .AND. (y13(3) .GE. bbl(3))) then p = ncell(j,k,l)+1 ncell(j,k,l) = p ! enlarge this list if nedded if(p .GT. size(cell(j,k,l)%list)) then CALL AllocateLL(1,j,k,l,istat) end if cell(j,k,l)%list(p) = i end if</pre>
231 232 233 234 235 236 237 238 239 240	<pre>! point where normal through centroid intersects triangle y13(1:3) = bbc(1:3)+lambda*n3(1:3) ! if this point is inside the cell, add triangle to cell if((y13(1) .LT. bbu(1)) .AND. (y13(2) .LT. bbu(2)) .AND. & (y13(3) .LT. bbu(3)) .AND. (y13(1) .GE. bbl(1)) .AND. & (y13(2) .GE. bbl(2)) .AND. (y13(3) .GE. bbl(3))) then p = ncell(j,k,1)+1 ncell(j,k,1) = p ! enlarge this list if nedded if(p .GT. size(cell(j,k,1)%list)) then CALL AllocateLL(1,j,k,1,istat) end if cell(j,k,1)%list(p) = i end if end if</pre>
231 232 233 234 235 236 237 238 239 240 241	<pre>! point where normal through centroid intersects triangle y13(1:3) = bbc(1:3)+lambda*n3(1:3) ! if this point is inside the cell, add triangle to cell if((y13(1) .LT. bbu(1)) .AND. (y13(2) .LT. bbu(2)) .AND. & (y13(3) .LT. bbu(3)) .AND. (y13(1) .GE. bbl(1)) .AND. & (y13(2) .GE. bbl(2)) .AND. (y13(3) .GE. bbl(3))) then p = ncell(j,k,1)+1 ncell(j,k,1) = p ! enlarge this list if nedded if(p .GT. size(cell(j,k,1)%list)) then CALL AllocateLL(1,j,k,1,istat) end if end if end if end if end if end if end if</pre>
231 232 233 234 235 236 237 238 239 240 241 242	<pre>! point where normal through centroid intersects triangle y13(1:3) = bbc(1:3)+lambda*n3(1:3) ! if this point is inside the cell, add triangle to cell if((y13(1) .LT. bbu(1)) .AND. (y13(2) .LT. bbu(2)) .AND. & (y13(3) .LT. bbu(3)) .AND. (y13(1) .GE. bbl(1)) .AND. & (y13(2) .GE. bbl(2)) .AND. (y13(3) .GE. bbl(3))) then p = ncell(j,k,1)+1 ncell(j,k,1) = p ! enlarge this list if nedded if(p .GT. size(cell(j,k,1)%list)) then CALL AllocateLL(1,j,k,1,istat) end if cell(j,k,1)%list(p) = i end if end if end if end if end do</pre>
231 232 233 234 235 236 237 238 239 240 241 242 243	<pre>! point where normal through centroid intersects triangle y13(1:3) = bbc(1:3)+lambda*n3(1:3) ! if this point is inside the cell, add triangle to cell if((y13(1) .LT. bbu(1)) .AND. (y13(2) .LT. bbu(2)) .AND. & (y13(3) .LT. bbu(3)) .AND. (y13(1) .GE. bbl(1)) .AND. & (y13(2) .GE. bbl(2)) .AND. (y13(3) .GE. bbl(3))) then p = ncell(j,k,1)+1 ncell(j,k,1) = p ! enlarge this list if nedded if(p .GT. size(cell(j,k,1)%list)) then CALL AllocateLL(1,j,k,1,istat) end if cell(j,k,1)%list(p) = i end if end if end do end do</pre>
231 232 233 234 235 236 237 238 239 240 241 242 243 244	<pre>! point where normal through centroid intersects triangle y13(1:3) = bbc(1:3)+lambda*n3(1:3) ! if this point is inside the cell, add triangle to cell if((y13(1) .LT. bbu(1)) .AND. (y13(2) .LT. bbu(2)) .AND. & (y13(2) .GE. bbl(2)) .AND. (y13(1) .GE. bbl(1)) .AND. & (y13(2) .GE. bbl(2)) .AND. (y13(3) .GE. bbl(3))) then p = ncell(j,k,1)+1 ncell(j,k,1) = p ! enlarge this list if nedded if(p .GT. size(cell(j,k,1)%list)) then CALL AllocateLL(1,j,k,1,istat) end if cell(j,k,1)%list(p) = i end if end if end if end do end do end do end do</pre>
231 232 233 234 235 236 237 238 239 240 241 242 243 244 245	<pre>! point where normal through centroid intersects triangle y13(1:3) = bbc(1:3)+lambda*n3(1:3) ! if this point is inside the cell, add triangle to cell if((y13(1) .LT. bbu(1)) .AND. (y13(2) .LT. bbu(2)) .AND. & (y13(2) .GE. bbl(2)) .AND. (y13(1) .GE. bbl(1)) .AND. & (y13(2) .GE. bbl(2)) .AND. (y13(3) .GE. bbl(3))) then p = ncell(j,k,1)+1 ncell(j,k,1) = p ! enlarge this list if nedded if(p .GT. size(cell(j,k,1)%list)) then CALL AllocateLL(1,j,k,1,istat) end if cell(j,k,1)%list(p) = i end if end if end if end do end do end do end do end do</pre>
231 232 233 234 235 236 237 238 239 240 241 242 243 244 245	<pre>! point where normal through centroid intersects triangle y13(1:3) = bbc(1:3)+lambda*n3(1:3) ! if this point is inside the cell, add triangle to cell if((y13(1) .LT. bbu(1)) .AND. (y13(2) .LT. bbu(2)) .AND. & (y13(2) .GE. bbl(2)) .AND. (y13(1) .GE. bbl(1)) .AND. & (y13(2) .GE. bbl(2)) .AND. (y13(3) .GE. bbl(3))) then p = ncell(j,k,1) = p ! enlarge this list if nedded if(p .GT. size(cell(j,k,1)%list)) then CALL AllocateLL(1,j,k,1,istat) end if cell(j,k,1)%list(p) = i end if end if end do end do end do end do end do</pre>
231 232 233 234 235 236 237 238 239 240 241 242 243 244 245	<pre>! point where normal through centroid intersects triangle y13(1:3) = bbc(1:3)+lambda*n3(1:3) ! if this point is inside the cell, add triangle to cell if((y13(1) .LT. bbu(1)) .AND. (y13(2) .LT. bbu(2)) .AND. & (y13(2) .GE. bbl(2)) .AND. (y13(1) .GE. bbl(1)) .AND. & (y13(2) .GE. bbl(2)) .AND. (y13(3) .GE. bbl(3))) then p = ncell(j,k,1) = p ! enlarge this list if nedded if(p .GT. size(cell(j,k,1)%list)) then CALL AllocateLL(1,j,k,1,istat) end if cell(j,k,1)%list(p) = i end if end if end if end do end do end do</pre>
231 232 233 234 235 236 237 238 239 240 241 242 243 244 245	<pre>! point where normal through centroid intersects triangle y13(1:3) = bbc(1:3)+lambda*n3(1:3) ! if this point is inside the cell, add triangle to cell if((y13(1) .LT. bbu(1)) .AND. (y13(2) .LT. bbu(2)) .AND. & (y13(2) .GE. bbl(2)) .AND. (y13(1) .GE. bbl(1)) .AND. & (y13(2) .GE. bbl(2)) .AND. (y13(3) .GE. bbl(3))) then p = ncell(j,k,1)+1 ncell(j,k,1) = p ! enlarge this list if nedded if(p .GT. size(cell(j,k,1)%list)) then CALL AllocateLL(1,j,k,1,istat) end if cell(j,k,1)%list(p) = i end if end if end do end do end do end do ! ! prune cell lists to actual size to save memory</pre>
231 232 233 234 235 236 237 238 239 240 241 242 242 242 242	<pre>! point where normal through centroid intersects triangle y13(1:3) = bbc(1:3)+lambda*n3(1:3) ! if this point is inside the cell, add triangle to cell if((y13(1) .LT. bbu(1)) .AND. (y13(2) .LT. bbu(2)) .AND. & (y13(3) .LT. bbu(3)) .AND. (y13(1) .GE. bbl(1)) .AND. & (y13(2) .GE. bbl(2)) .AND. (y13(3) .GE. bbl(3))) then p = ncell(j,k,l)+1 ncell(j,k,l) = p ! enlarge this list if nedded if(p .GT. size(cell(j,k,l)%list)) then CALL AllocateLL(1,j,k,l,istat) end if cell(j,k,l)%list(p) = i end if end if end do end do end do end do end do end do end do</pre>
231 232 233 234 235 236 237 238 239 240 241 242 243 244 245	<pre>! point where normal through centroid intersects triangle y13(1:3) = bbc(1:3)+lambda*n3(1:3) ! if this point is inside the cell, add triangle to cell if((y13(1) .LT. bbu(1)) .AND. (y13(2) .LT. bbu(2)) .AND. & (y13(2) .GE. bbl(2)) .AND. (y13(1) .GE. bbl(1)) .AND. & (y13(2) .GE. bbl(2)) .AND. (y13(3) .GE. bbl(3))) then p = ncell(j,k,1)+1 ncell(j,k,1) = p ! enlarge this list if nedded if(p .GT. size(cell(j,k,1)%list)) then CALL AllocateLL(1,j,k,1,istat) end if cell(j,k,1)%list(p) = i end if end if end do end do end do ! ! prune cell lists to actual size to save memory !</pre>
231 232 233 234 235 236 237 238 239 240 241 242 243 244 245	<pre>! point where normal through centroid intersects triangle y13(1:3) = bbc(1:3)+lambda*n3(1:3) ! if this point is inside the cell, add triangle to cell if((y13(1) .LT. bbu(1)) .AND. (y13(2) .LT. bbu(2)) .AND. & (y13(2) .GE. bbl(2)) .AND. (y13(1) .GE. bbl(1)) .AND. & (y13(2) .GE. bbl(2)) .AND. (y13(3) .GE. bbl(3))) then p = ncell(j,k,1) = p ! enlarge this list if nedded if(p .GT. size(cell(j,k,1)%list)) then CALL AllocateLL(1,j,k,1,istat) end if cell(j,k,1)%list(p) = i end if end if end do end do end do end do i=1,NCELLX do i=1,NCELLX</pre>
231 232 233 234 235 236 237 238 239 240 241 242 243 244 245	<pre>! point where normal through centroid intersects triangle y13(1:3) = bbc(1:3)+lambda*n3(1:3) ! if this point is inside the cell, add triangle to cell if((y13(1) .LT. bbu(1)) .AND. (y13(2) .LT. bbu(2)) .AND. & (y13(2) .GE. bbl(2)) .AND. (y13(1) .GE. bbl(1)) .AND. & (y13(2) .GE. bbl(2)) .AND. (y13(3) .GE. bbl(3))) then p = ncell(j,k,1) = p ! enlarge this list if nedded if(p .GT. size(cell(j,k,1)%list)) then CALL AllocateLL(1,j,k,1,istat) end if cell(j,k,1)%list(p) = i end if end if end do end do end do end do i=1,NCELLX do j=1,NCELLY</pre>
231 232 233 234 235 236 237 238 239 240 241 242 242 242 242 244 245	<pre>! point where normal through centroid intersects triangle y13(1:3) = bbc(1:3)+lambda*n3(1:3) ! if this point is inside the cell, add triangle to cell if((y13(1) .LT. bbu(1)) .AND. (y13(2) .LT. bbu(2)) .AND. & (y13(2) .GE. bbl(2)) .AND. (y13(1) .GE. bbl(1)) .AND. & (y13(2) .GE. bbl(2)) .AND. (y13(3) .GE. bbl(3))) then p = ncell(j,k,1)+1 ncell(j,k,1) = p ! enlarge this list if nedded if(p .GT. size(cell(j,k,1)%list)) then CALL AllocateLL(1,j,k,1,istat) end if cell(j,k,1)%list(p) = i end if end if end do end do end do end do i=1,NCELLX do j=1,NCELLY do k=1,NCELLZ</pre>
231 232 233 234 235 236 237 238 239 240 241 242 243 244 245 246 247 248 249	<pre>! point where normal through centroid intersects triangle y13(1:3) = bbc(1:3)+lambda*n3(1:3) ! if this point is inside the cell, add triangle to cell if((y13(1) .LT. bbu(1)) .AND. (y13(2) .LT. bbu(2)) .AND. & (y13(2) .GE. bbl(2)) .AND. (y13(1) .GE. bbl(1)) .AND. & (y13(2) .GE. bbl(2)) .AND. (y13(3) .GE. bbl(3))) then p = ncell(j,k,1)+1 ncell(j,k,1) = p ! enlarge this list if nedded if(p .GT. size(cell(j,k,1)%list)) then CALL AllocateLL(1,j,k,1,istat) end if cell(j,k,1)%list(p) = i end if end if end do end do end do end do i=1,NCELLX do j=1,NCELLY do k=1,NCELLZ cell(i,j,k)%list=>reallocate(cell(i,j,k)%list,ncell(i,j,k))</pre>
231 232 233 234 235 236 237 238 239 240 241 242 243 244 245 245	<pre>! point where normal through centroid intersects triangle y13(1:3) = bbc(1:3)+lambda*n3(1:3) ! if this point is inside the cell, add triangle to cell if((y13(1) .LT. bbu(1)) .AND. (y13(2) .LT. bbu(2)) .AND. & (y13(2) .GE. bbl(2)) .AND. (y13(1) .GE. bbl(1)) .AND. & (y13(2) .GE. bbl(2)) .AND. (y13(3) .GE. bbl(3))) then p = ncell(j,k,1) = p ! enlarge this list if nedded if(p .GT. size(cell(j,k,1)%list)) then CALL AllocateLL(1,j,k,1,istat) end if cell(j,k,1)%list(p) = i end if end if end do end do end do end do i=1,NCELLX do j=1,NCELLZ cell(i,j,k)%list=>reallocate(cell(i,j,k)%list,ncell(i,j,k)) end do end do</pre>
231 232 233 234 235 236 237 238 239 240 241 242 243 244 245 241 242 243 244 245	<pre>! point where normal through centroid intersects triangle y13(1:3) = bbc(1:3)+lambda*n3(1:3) ! if this point is inside the cell, add triangle to cell if((y13(1) .LT. bbu(1)) .AND. (y13(2) .LT. bbu(2)) .AND. & (y13(2) .GE. bbl(2)) .AND. (y13(1) .GE. bbl(1)) .AND. & (y13(2) .GE. bbl(2)) .AND. (y13(3) .GE. bbl(3))) then p = ncell(j,k,1) = p ! enlarge this list if nedded if(p .GT. size(cell(j,k,1)%list)) then CALL AllocateLL(1,j,k,1,istat) end if cell(j,k,1)%list(p) = i end if end if end do end do end do end do ! ! prune cell lists to actual size to save memory ! do i=1,NCELLX do j=1,NCELLZ cell(i,j,k)%list=>reallocate(cell(i,j,k)%list,ncell(i,j,k)) end do end do end do end do</pre>
231 232 233 234 235 236 237 238 239 240 241 242 242 243 244 245 242 242 242 242 242 242 242 242	<pre>! point where normal through centroid intersects triangle y13(1:3) = bbc(1:3)+lambda*n3(1:3) ! if this point is inside the cell, add triangle to cell if((y13(1) .LT. bbu(1)) .AND. (y13(2) .LT. bbu(2)) .AND. & (y13(2) .GE. bbl(2)) .AND. (y13(1) .GE. bbl(1)) .AND. & (y13(2) .GE. bbl(2)) .AND. (y13(3) .GE. bbl(3))) then p = ncell(j,k,1)+1 ncell(j,k,1) = p ! enlarge this list if nedded if(p .GT. size(cell(j,k,1)%list)) then CALL AllocateLL(1,j,k,1,istat) end if cell(j,k,1)%list(p) = i end if end do end do end do end do !</pre>
231 232 233 234 235 236 237 238 239 240 241 242 243 244 245 244 245 244 245 244 245 251 252 252 255	<pre>! point where normal through centroi intersects triangle y13(1:3) = bbc(1:3)+lambda*n3(1:3) ! if this point is inside the cell, add triangle to cell if((y13(1) .LT. bbu(1)) .AND. (y13(2) .LT. bbu(2)) .AND. & (y13(3) .LT. bbu(3)) .AND. (y13(1) .GE. bbl(1)) .AND. & (y13(2) .GE. bbl(2)) .AND. (y13(3) .GE. bbl(3))) then p = ncell(j,k,1) = p ! enlarge this list if nedded if(p .GT. size(cell(j,k,1)%list)) then CALL AllocateLL(1,j,k,1,istat) end if cell(j,k,1)%list(p) = i end if end if end do end do end do end do i=nd do i=1,NCELLX do i=1,NCELLY do k=1,NCELLZ cell(i,j,k)%list=>reallocate(cell(i,j,k)%list,ncell(i,j,k)) end do end do</pre>

254 END subroutine SortT

C.5.4.10 Allocate dynamic list memory (AllocateLL.f90)

AllocateLL is a helper routine for the handling of the bin and cell lists by SortT and the main program. It implements the whole memory management for the dynamic list data structures as described in section 7.2.6. The input argument flag determines the action to be taken. If the routine is called with flag=0, the initial memory allocation and general set-up of the lists is done. With flag=1, the list associated with cell ix, iy, iz (or bin iy, iz if ix is zero) is enlarged by a

certain number of elements defined by the parameter STEP. If called with flag=2, all lists are deallocated and the memory is freed again. The variable info is a status flag and contains 0 if everything was successful, 1 if an error occurred.

subroutine AllocateLL(flag, ix, iy, iz, info) 1 2 USE globals USE Util 3 IMPLICIT NONE 4 1-----!ALLOCATELL allocates memory for triangle lists. AllocateLL allocates memory for the bin and cell lists. If called with flag=0, initial allocation is done. Calling with flag=1 results in enlargement of existing arrays. Calling with flag=2 deallocates all memory. info returns 0 if allocation was successful, 1 otherwise. See also INIT_PART, POINT_IN_DOMAIN, SORTT, UTIL _____ DIPLOMA THESIS WS01/02 ICOS ETH-ZUERICH 1 -----PROTEIN DIFFUSION INSIDE THE ENDOPLASMIC RETICULUM ! enlarge arrays by STEP entries everytime allocatell(1,ix,iy,iz,info) is called 5 INTEGER, PARAMETER :: STEP = 10 _____ ! Declaration of input/output variables ------! flag: 0=initialize, 1=reallocate, 2=deallocate INTEGER, INTENT(IN) 6 :: flag ! info: 0=everything OK, 1=error occured INTEGER, INTENT(OUT) 7 :: info ! x,y,z indices of cell/bin to (re)allocate list for INTEGER, INTENT(IN) 8 :: ix, iy, iz _____ ! Declaration of local variables _____ ! loop counters 9 INTEGER :: i, j, k ! error status INTEGER 10 :: istat ! length of list 11 INTEGER :: llen 12 info = 0-----! flag=0: Initialize ----if(flag .EQ. 0) then 13 _____ ! Bin lists for point_in_domain (contain all triangles inside bin) 1 ---14 istat = 0 ! allocate matrix to hold lengths of lists ! first index: biny, second index: binz ALLOCATE(nbin(NBINY, NBINZ), STAT=istat) 15 16 if (istat .NE. 0) then 17 WRITE(*,'(2A)') 'Unable to allocate memory for nbin.' 18 info = 119 return 20 end if 21 nbin(1:NBINY,1:NBINZ) = 0 22 istat = 0! allocate array of pointers to lists

! first index: biny, second index: binz 23 ALLOCATE(bin(NBINY, NBINZ), STAT=istat) 24 if (istat .NE. 0) then 25 WRITE(*,'(2A)') 'Unable to allocate memory for bin lists.' 26 info = 127 return 28 end if ! nullify all pointers (associate no lists at the beginning) 29 do i=1,NBINY 30 do j=1,NBINZ 31 nullify(bin(i,j)%list) 32 end do 33 end do ! Cell lists for nearest tri. search (contain only triangles in in-sphere) ! cell list for nearest triangle search is only needed if boundary ! condition handling is enabled 34 if(BC .GT. 0) then ! calculate number of cells in each direction based on cut-off radius 35 NCELLX = int((xmax-xmin)/rc) 36 NCELLY = int((ymax-ymin)/rc) NCELLZ = int((zmax-zmin)/rc) 37 38 if(NCELLX .EQ. 0) NCELLX = 1 ! make at least one cell if(NCELLX .EQ. 0) NCELLX = 1 39 40 if(NCELLX .EQ. 0) NCELLX = 1 ! cell sizes 41 dcx = (xmax-xmin)/real(NCELLX) dcy = (ymax-ymin)/real(NCELLY) 42 dcz = (zmax-zmin)/real(NCELLZ) 43 44 istat = 0! allocate matrix to hold list lengths ! first index: cellx, second: celly, third: cellz ALLOCATE(ncell(NCELLX,NCELLY,NCELLZ), STAT=istat) 45 46 if (istat .NE. 0) then WRITE(*,'(2A)') 'Unable to allocate memory for ncell.' 47 48 info = 149 return 50 end if ncell(1:NCELLX,1:NCELLY,1:NCELLZ) = 0 51 52 istat = 0 ! allocate array of pointers to lists ! first index: cellx, second: celly, third: cellz 53 ALLOCATE(cell(NCELLX,NCELLY,NCELLZ), STAT=istat) 54 if (istat .NE. 0) then 55 WRITE(*,'(2A)') 'Unable to allocate memory for cell lists.' 56 info = 157 return 58 end if ! nullify all pointers in the beginning 59 do i=1,NCELLX 60 do j=1,NCELLY 61 do k=1,NCELLZ 62 nullify(cell(i,j,k)%list) 63 end do 64 end do 65 end do 66 end if 1----! flag=1: Enlarge lists _____ 67 elseif(flag .EQ. 1) then if(ix .EQ. 0) then 68 ! no x index given: operation affects bin list only 69 istat = 070 if(ASSOCIATED(bin(iy,iz)%list)) then ! if list is already assigned, reallocate it 71 llen = size(bin(iy,iz)%list) 72 bin(iy,iz)%list=>reallocate(bin(iy,iz)%list,llen+STEP) 73 else ! if there is no list yet, allocate one 74 ALLOCATE(bin(iy,iz)%list(STEP), STAT=istat) if(istat .NE. 0) then 75 WRITE(*,'(A,I5,A,I5,A)') 'Error reallocating bin list & 76 (',iy,',',iz,')'

```
77
                  info = 1
78
                  return
79
               end if
80
            end if
81
         else
            ! x index given: operation affects cell list only
82
            if(BC .GT. 0) then
83
               istat = 0
               if(ASSOCIATED(cell(ix,iy,iz)%list)) then
84
                  llen = size(cell(ix,iy,iz)%list)
85
                  cell(ix,iy,iz)%list=>reallocate(cell(ix,iy,iz)%list,llen+STEP)
86
87
               else
88
                  ALLOCATE(cell(ix,iy,iz)%list(STEP), STAT=istat)
89
                  if(istat .NE. 0) then
90
                      WRITE(*,'(A,I5,A,I5,A,I5,A)') 'Error reallocating cell &
                        list(',ix,',',iy,',',iz,')'
91
                      info = 1
92
                     return
93
                  end if
94
               end if
            end if
95
96
         end if
      ! flag=2: Deallocate all lists
97
      elseif(flag .EQ. 2) then
         ! bin lists: always
98
         do i=1.NBINY
            do j=1,NBINZ
99
               if(ALLOCATED(bin(i,j)%list)) DEALLOCATE(bin(i,j)%list)
100
101
            end do
         end do
102
         if(ALLOCATED(bin)) DEALLOCATE(bin)
103
         if(ALLOCATED(nbin)) DEALLOCATE(nbin)
104
         ! cell lists: only if boundary condition handling is active
105
         if(BC .GT. 0) then
            do i=1.NCELLX
106
107
               do j=1,NCELLY
108
                  do k=1,NCELLZ
                     if(ALLOCATED(cell(i,j,k)%list)) DEALLOCATE(cell(i,j,k)%list)
109
                  end do
110
111
               end do
112
            end do
113
            if(ALLOCATED(cell)) DEALLOCATE(cell)
114
            if(ALLOCATED(ncell)) DEALLOCATE(ncell)
115
         end if
116
      endif
117
      return
118
      END subroutine AllocateLL
```

C.5.4.11 Determine if point is in domain (point_in_domain.f90)

The next function determines whether the point given by its input argument $p \in \mathbb{R}^3$ is inside or outside the closed surface described by the triangulated set. If the point is inside, the value 1 is returned, 0 if it is outside. If boundary condition handling is switched on (i.e. the parameter **boundarycondition** in the input file is > 0), a value of 2 is returned if the point is inside the domain *and* closer than **cutoff** to its boundary. These are the points that need to be mirrored to the outside. In order to be able to do so, the function in this case also returns the index of the closest triangle in **closest** and the orthogonal distance of the point from this triangle in **lammin**. Bin and cell lists are used to accelerate the regular triangle handling and the search for the closest triangle, respectively.

If the parameter debug is set to .TRUE., a set of additional files is created every time the function is invoked. The file bin.deb contains the edges of the current bin to be visualized with gnuplot. The file tlist.deb contains the vertices of all the triangles that belong to this bin, triang.deb contains the vertices of the

triangles that actually have been intersected by the line $P + \lambda[1,0,0]$, points.deb contains the coordinates of all the intersection points and coeff.deb contains the distances from point P to all the intersection points with the triangles. All files can be visualized using gnuplot and the sp command. Figure 7.3 has for example been created this way.

- 1 function point_in_domain(p, lammin, closest)
- 2 USE globals

```
IMPLICIT NONE
3
```

4

5

6

7

8

9

```
! {\tt POINT\_IN\_DOMAIN} \  \  {\tt checks whether a point is inside the domain or not}.
        if the point given by p(1:3) is inside the closed surface
        defined by triang(...), the function returns 1, otherwise
        0. If the point is inside AND closer to the boundary than rc,
        the value 2 is returned. The return valiable lammin contains the
        orthogonal distance to the closest triangle and closest is its index.
        The latter two are only calculated if BC is not 0.
        See also INIT_PART, SORTT
     DIPLOMA THESIS WS01/02 ICOS
                                                              ETH-ZUERICH
                  PROTEIN DIFFUSION INSIDE THE ENDOPLASMIC RETICULUM
      ----- ivo f. sbalzarini ------
     ! Declaration of external functions
     ! intersects a triangle with a given line
     REAL(MK), EXTERNAL
                                     :: intersect
     ! Declaration of input and output variables
     ! point in question
     REAL(MK), DIMENSION(3), INTENT(IN) :: p
     ! return value
     INTEGER
                                    :: point_in_domain
     ! distance to closest triangle
     REAL(MK), INTENT(OUT)
                                     :: lammin
     ! index of closest triangle
     INTEGER, INTENT(OUT)
                                     :: closest
     ! Declaration of local variables
     1----
                   -----
     ! loop counter
     INTEGER
                                     :: i, j, k, l, tri
     ! angles in spherical coordinates
                                     :: phi, theta
10
     REAL(MK)
     ! edge vectors and direction
11
     REAL(MK), DIMENSION(3)
                                     :: a, b
     ! system matrix
     REAL(MK), DIMENSION(3,3)
12
                                     :: mat
     ! number of intersection points
13
     INTEGER
                                     :: nip
     ! line parameter lambda
14
     REAL(MK)
                                      :: lam
     ! x, y and z indices of current bin/cell
15
     INTEGER
                                      :: idxx, idxy, idxz
     ! produce debuging output ?
16
     LOGICAL
                                     :: debug
     1-----
     ! Follow x direction and intersect with all triangles
                                                       _____
```

```
17
      debug = .false.
      !if(p(1) .gt. 400.0_MK .and. p(1) .lt. 402.0_MK .and. &
      ! p(2) .gt. 386.0_MK .and. p(2) .lt. 388.0_MK .and. &
          p(3) .gt. 5.0_MK .and. p(3) .lt. 5.5_MK) then
          debug = .true.
          print*,'point is: ',p
      !end if
18
      nip = 0
      mat(1,1) = 0.0
19
      mat(2,1) = 0.0
20
21
      idxy = ceiling((p(2)-ymin)/dby)
      if(idxy .EQ. 0) idxy = 1
22
                                            ! account for boundary points
23
      if(idxy .EQ. NBINY+1) idxy = NBINY
24
      idxz = ceiling((p(3)-zmin)/dbz)
25
      if(idxz . EQ. 0) idxz = 1
                                             ! account for boundary points
26
      if(idxz .EQ. NBINZ+1) idxz = NBINZ
27
      if(debug) then
         print*,'bin indices: ',idxy,idxz
28
         print*, 'nbin=', nbin(idxy, idxz)
29
         open(70,file='bin.deb',status='replace')
30
         WRITE(70,*) xmin, ymin+(idxy-1)*dby, zmin+(idxz-1)*dbz
31
         WRITE(70,*) xmin, ymin+(idxy)*dby, zmin+(idxz-1)*dbz
32
33
         WRITE(70,*) xmin, ymin+(idxy)*dby, zmin+(idxz)*dbz
         WRITE(70,*) xmin, ymin+(idxy-1)*dby, zmin+(idxz)*dbz
34
35
         WRITE(70,*) xmin, ymin+(idxy-1)*dby, zmin+(idxz-1)*dbz
36
         WRITE(70, '(A)')
37
         WRITE(70, '(A)')
         WRITE(70,*) xmax, ymin+(idxy-1)*dby, zmin+(idxz-1)*dbz
38
39
         WRITE(70,*) xmax, ymin+(idxy)*dby, zmin+(idxz-1)*dbz
40
         WRITE(70,*) xmax, ymin+(idxy)*dby, zmin+(idxz)*dbz
         WRITE(70,*) xmax, ymin+(idxy-1)*dby, zmin+(idxz)*dbz
41
         WRITE(70,*) xmax, ymin+(idxy-1)*dby, zmin+(idxz-1)*dbz
42
43
         WRITE(70, '(A)')
         WRITE(70, '(A)')
44
         WRITE(70,*) xmin, ymin+(idxy-1)*dby, zmin+(idxz-1)*dbz
45
         WRITE(70,*) xmax, ymin+(idxy-1)*dby, zmin+(idxz-1)*dbz
46
         WRITE(70,'(A)')
47
         WRITE(70, '(A)')
48
         WRITE(70,*) xmin, ymin+(idxy)*dby, zmin+(idxz-1)*dbz
49
50
         WRITE(70,*) xmax, ymin+(idxy)*dby, zmin+(idxz-1)*dbz
         WRITE(70,'(A)')
WRITE(70,'(A)')
51
52
         WRITE(70,*) xmin, ymin+(idxy)*dby, zmin+(idxz)*dbz
53
54
         WRITE(70,*) xmax, ymin+(idxy)*dby, zmin+(idxz)*dbz
55
         WRITE(70,'(A)')
56
         WRITE(70, '(A)')
         WRITE(70,*) xmin, ymin+(idxy-1)*dby, zmin+(idxz)*dbz
57
58
         WRITE(70,*) xmax, ymin+(idxy-1)*dby, zmin+(idxz)*dbz
59
         CLOSE(70)
60
      end if
      ! intersect with all triangles in the same bin
61
      do j=1,nbin(idxy,idxz)
         ! get index of next triangle
62
         i = bin(idxy,idxz)%list(j)
63
         if(debug) then
            open(70,file='tlist.deb',position='append')
64
65
            WRITE(70,*) triang(:,1,i)
            WRITE(70,*) triang(:,2,i)
66
67
            WRITE(70,*) triang(:,3,i)
            WRITE(70,*) triang(:,1,i)
68
69
            WRITE(70, '(A)')
            WRITE(70, '(A)')
70
71
            CLOSE(70)
72
         end if
         ! edge vectors
         a = triang(:,2,i)-triang(:,1,i)
73
         b = triang(:,3,i)-triang(:,1,i)
74
         ! determinant of system matrix
75
         phi = -a(3)*b(2)+a(2)*b(3)
                                               ! for v=(1,0,0)
76
         if (abs(phi) .GE. TOL) then
                                               ! if intersection point exists
            phi = 1.0/phi
77
            ! inverse of system matrix
78
            mat(1,2) = phi*b(3)
            mat(1,2) phi/s(2)
mat(1,3) = -phi/s(2)
mat(2,2) = -phi/s(3)
79
80
            mat(2,3) = phi*a(2)
81
            mat(3,1) = phi*(-a(3)*b(2)+a(2)*b(3))
82
            mat(3,2) = phi*(a(3)*b(1)-a(1)*b(3))
83
```

```
84
           mat(3,3) = phi*(-a(2)*b(1)+a(1)*b(2))
            ! right hand side
85
           a = p-triang(:,1,i)
            ! solution of 3x3 system
86
           b = (mat(:,1)*a(1))+(mat(:,2)*a(2))+(mat(:,3)*a(3))
            ! intersect inside triangle ?
87
           if ((b(1)+b(2) .LE. 1.0_MK) .AND. (b(1) .GE. 0.0_MK) &
                 .AND. (b(2) .GE. TOL) .AND. (b(3) .GE. TOL)) then
88
               if(debug) then
                 print*,'intersected triangle ',i
89
90
                 OPEN(70,file='triang.deb',position='append')
91
                 WRITE(70,*) triang(:,1,i)
                 WRITE(70,*) triang(:,2,i)
92
93
                 WRITE(70,*) triang(:,3,i)
94
                 WRITE(70,*) triang(:,1,i)
95
                 WRITE(70,'(A)')
96
                 WRITE(70, '(A)')
97
                 CLOSE(70)
98
                 open(70,file='points.deb',position='append')
                 write(70,*) triang(:,1,i)+b(1)*(triang(:,2,i)-triang(:,1,i)) &
99
                      +b(2)*(triang(:,3,i)-triang(:,1,i))
100
                 close(70)
101
                 open(70,file='coeff.deb',position='append')
102
                 write(70,*) b
103
                 close(70)
104
              end if
105
              nip=nip+1
106
           end if
        end if
107
108
     end do
      ! Odd or even number of intersections ?
     if(mod(nip,2) .EQ. 0.0) then
109
110
        point_in_domain = 0
111
     else
112
        point_in_domain = 1
113
     end if
     if(debug) print*,'point_in_domain=',point_in_domain
114
      1------
      ! For the points inside: check if they are close to the boundary
      ! and find the nearest triangle as well as the distance to it
                                                                     _____
      1----
              ------
     if(BC .GT. 0) then
115
                                    ! only do this if needed for the B.C.
116
        if(point_in_domain .EQ. 1) then
117
           lammin = HUGE(lam)
118
           closest = 0
            ! determine indices of cell we are in
119
           idxx = ceiling((p(1)-xmin)/dcx)
120
           if(idxx .EQ. 0) idxx = 1
                                                ! account for boundary points
121
           if(idxx .EQ. NCELLX+1) idxx = NCELLX
            idxy = ceiling((p(2)-ymin)/dcy)
122
123
            if(idxy .EQ. 0) idxy = 1
                                                ! account for boundary points
            if(idxy .EQ. NCELLY+1) idxy = NCELLY
124
125
           idxz = ceiling((p(3)-zmin)/dcz)
126
            if(idxz .EQ. 0) idxz = 1
                                                ! account for boundary points
           if(idxz .EQ. NCELLZ+1) idxz = NCELLZ
127
            ! loop over all triangles in this cell and its nearest neighbors
128
           do i=idxx-1,idxx+1
129
              do j=idxy-1,idxy+1
130
                 do k=idxz-1,idxz+1
                    ! if this is a valid cell ...
131
                    if((i .GT. 0 .AND. j .GT. 0 .AND. k .GT. 0) .AND. &
                       (i .LE. NCELLX .AND. j .LE. NCELLY .AND. k .LE. NCELLZ)) then
                        ! ... loop over all triangles in this cell
                       do l=1,ncell(i,j,k)
132
                          ! get nextr triangle from the cell list
133
                          tri = cell(i,j,k)%list(1)
                           ! intersect normal through point with triangle
134
                          lam = intersect(p, normal(1:3,tri), tri)
                          if(lam .GE. 0) then ! if intersection pt. exists
135
                             if(lam .LE. lammin) then
136
                                ! if this is closer than anything before
137
                                lammin = lam
138
                                closest = tri ! => update data
                             end if
139
```

```
140
                           end if
141
                        end do
                        ! if point is closer to surface than rc, set return
                        ! value to 2. NOTE: normal(:,:) are unit vectors
142
                        if(lammin .LE. rc) point_in_domain = 2
                     end if
143
144
                  end do
145
               end do
146
            end do
147
         end if
148
      end if
149
      return
150
      end function point_in_domain
```

C.5.4.12 Create voxel representation (Voxelize.f90)

This routine converts the triangulated surface to a voxel representation according to algorithm 5.1. The geometry is covered by voxels of extension $vx \times vy \times vz$. resx, resy and resz are output variables containing the actual number of voxels that have been generated in each spatial direction. info finally is a status flag which contains 0 on success, 1 otherwise. The binary voxel representation where all the voxels that are intersected by the surface triangulation are set to 1, all others to zero is contained in the global array INTEGER voxel(resx,resy,resz) (see section C.5.4.1) after the call to this routine.

ETH-ZUERICH

subroutine Voxelize(vx,vy,vz,resx,resy,resz,info) 1 USE globals 2 3 IMPLICIT NONE _____ !Voxelize converts the surface to a 3D voxel representation. VOXELIZE finds a voxel representation of the surface using voxels of size vx*vy*vz. The representation is contained in the global binary 3D array voxel (int). resx, resy, resz return the actual number of voxels used in it. info returns 0 on success, else 1 See also INIT_PART, BCDIM _____ DIPLOMA THESIS WS01/02 ICOS PROTEIN DIFFUSION INSIDE THE ENDOPLASMIC RETICULUM

```
! Declaration of input/output variables
                                           _____
     ! desired voxel size (in x, y, z)
4
     REAL(MK), INTENT(INOUT)
                                            :: vx, vy, vz
      ! number of voxels in x,y,z
5
     INTEGER, INTENT(OUT)
                                            :: resx, resy, resz
      ! error varaible
     INTEGER, INTENT(OUT)
6
                                            :: info
      ! Declaration of local variables
      1 ----
     ! loop counters
7
     INTEGER
                                            :: i, j, k, l, ix1, ix2, ix3
     ! minimum/maximum voxel indices in x,y,z of a given triangle
8
     INTEGER
                                            :: ixmin, ixmax, iymin, iymax
     INTEGER
                                            :: izmin, izmax
9
      ! indices of voxel
10
     INTEGER
                                            :: ixx, ixy, ixz
```

! vertex one of triangle, right hand side of eqs 11 REAL(MK), DIMENSION(3) :: x1, y1 ! triangle edge vectors and normal 12 REAL(MK), DIMENSION(3) :: a, b, no ! lower and upper boundaries of voxel and its centroid 13 REAL(MK), DIMENSION(3) :: bbl, bbu, bbc ! error trap 14 INTEGER :: istat ! system matrix REAL(MK), DIMENSION(3,3) 15 :: mat ! determinant REAL(MK) 16 :: det ! system solution REAL(MK) 17 :: alpha, beta, lambda -----! Initialize 18 info = 0resx = ceiling((xmax-xmin)/vx) 19 20 resy = ceiling((ymax-ymin)/vy) 21 resz = ceiling((zmax-zmin)/vz) 22 if(resx .EQ. 0) resx = 1if(resy .EQ. 0) resy = 123 24 if(resz .EQ. 0) resz = 1istat = 0 25 ! initialize array 26 if(ALLOCATED(voxel)) DEALLOCATE(voxel) 27 ALLOCATE(voxel(resx,resy,resz), STAT=istat) if(istat .NE. 0) then 28 29 WRITE(*,'(A)') 'Error in Voxelize: cannot allocate voxel array' 30 info = 131 return 32 end if 33 voxel(1:resx,1:resy,1:resz) = 0 vx = (xmax-xmin)/real(resx) 34 vy = (ymax-ymin)/real(resv) 35 36 vz = (zmax-zmin)/real(resz) 1-----! Discretize geometry onto voxels _____ 1---------do i=1.M 37 ! loop over all triangles 38 ixmin = resx+1 ! reset bounding box indices 39 ixmax = 0iymin = resy+1 40 41 iymax = 0 42 izmin = resz+1 izmax = 0 43 44 x1 = triang(1:3,1,i) ! vertex one of this triangle 45 a = triang(1:3,2,i)-x1 ! edge vectors 46 b = triang(1:3,3,i)-x1 47 no = normal(1:3,i) ! triangles normal det = -a(3)*b(2)*no(1)+a(2)*b(3)*no(1)+a(3)*b(1)*no(2)- & 48 a(1)*b(3)*no(2)-a(2)*b(1)*no(3)+a(1)*b(2)*no(3) 49 det = 1.0_MK/det 50 mat(1,1) = -b(3)*no(2)+b(2)*no(3)! system matrix (see notes) 51 mat(1,2) = b(3)*no(1)-b(1)*no(3) 52 mat(1,3) = -b(2)*no(1)+b(1)*no(2)53 mat(2,1) = a(3)*no(2)-a(2)*no(3)54 mat(2,2) = -a(3)*no(1)+a(1)*no(3)55 mat(2,3) = a(2)*no(1)-a(1)*no(2)56 mat(3,1) = -a(3)*b(2)+a(2)*b(3)57 mat(3,2) = a(3)*b(1)-a(1)*b(3)58 mat(3,3) = -a(2)*b(1)+a(1)*b(2) 59 do j=1,3 ! loop over all vertices ! determine index of voxel of this vertex 60 ixx = ceiling((triang(1,j,i)-xmin)/vx) 61 if(ixx .EQ. 0) ixx = 1! include boundaries at walls 62 if(ixx .EQ. resx+1) ixx = resx 63 ixy = ceiling((triang(2,j,i)-ymin)/vy) 64 if(ixy .EQ. 0) ixy = 1! include boundaries at walls if(ixy .EQ. resy+1) ixy = resy 65 ixz = ceiling((triang(3,j,i)-zmin)/vz) 66 67 if(ixz .EQ. 0) ixz = 1! include boundaries at walls if(ixz .EQ. resz+1) ixz = resz 68 ! set this voxel to one 69 voxel(ixx,ixy,ixz) = 1

```
! update bounding box of voxels
70
             if(ixx .LT. ixmin) ixmin = ixx
71
             if(ixx .GT. ixmax) ixmax = ixx
72
             if(ixy .LT. iymin) iymin = ixy
73
             if(ixy .GT. iymax) iymax = ixy
74
             if(ixz .LT. izmin) izmin = ixz
75
             if(ixz .GT. izmax) izmax = ixz
76
          end do
          ! loop over all voxels in the bounding box of the vertices
         do j=ixmin,ixmax
77
78
             bbl(1) = xmin+(j-1)*vx
                                        ! lower x
79
             bbu(1) = xmin+j*vx
                                         ! upper x
             do k=iymin,iymax
80
                bbl(2) = ymin+(k-1)*vy ! lower y
81
                bbu(2) = ymin+k*vy
82
                                           ! upper y
83
                do l=izmin,izmax
                   bbl(3) = zmin+(1-1)*vz ! lower z
84
                   bbu(3) = zmin+l*vz
                                              ! upper z
85
                   bbc(1:3) = 0.5*(bbl(1:3)+bbu(1:3)) ! voxels centroid
86
                   y1(1:3) = bbc(1:3)-x1(1:3)
87
                   alpha = det*(mat(1,1)*y1(1)+mat(1,2)*y1(2) +
                                                                      &
88
                        mat(1,3)*y1(3))
89
                   beta = det*(mat(2,1)*y1(1)+mat(2,2)*y1(2) +
                                                                      &
                        mat(2,3)*y1(3))
                   lambda = det*(mat(3,1)*y1(1)+mat(3,2)*y1(2) +
90
                                                                       k
                        mat(3,3)*y1(3))
                   y1(1:3) = bbc(1:3) + lambda*no(1:3)
91
                    ! if centroid is over triangle and intersection point
                    ! with triangle is inside voxel, set it to one
                   if ((alpha .GE. 0.0_MK) .AND. (beta .GE. 0.0_MK) &
.AND. (alpha+beta .LE. 1.0)) then
92
                      if(y1(1) .GE. bbl(1) .AND. y1(1) .LE. bbu(1) .AND.
y1(2) .GE. bbl(2) .AND. y1(2) .LE. bbu(2) .AND.
93
                                                                                &
                                                                                &
                          y1(3) .GE. bbl(3) .AND. y1(3) .LE. bbu(3)) voxel(j,k,l) = 1
                   end if
94
95
                end do
96
             end do
97
         end do
98
      end do
99
      END subroutine Voxelize
```

C.5.4.13 Determine box counting dimension (BCdim.f90)

BCdim calculates the box counting dimension of the geometry according to algorithm 5.2. The geometry is first converted to a voxel representation of initial voxel size $vx0 \times vy0 \times vz0$ using the previous routine. The input parameter thresh contains the threshold value for the binarization step as described in section 5.2. info is a status flag that contains 0 if the function returns without errors and 1 otherwise. The return value of the function BCdim is the least squares estimation of the box counting dimension.

1 function BCdim(vx0, vy0, vz0, thresh, info)

```
2 USE globals
```

3 IMPLICIT NONE

```
! Factor by which the box size is increased every reduction step
4
     REAL(MK), PARAMETER
                                          :: fact = 2.0_MK
                  _____
     ! Declaration of input/output variables
     ! box counting dimension
5
     REAL(MK)
                                          :: BCdim
     ! initial voxel size
6
     REAL(MK), INTENT(IN)
                                          :: vx0, vy0, vz0
      ! threshold for voxel values
     REAL(MK), INTENT(IN)
7
                                          :: thresh
     ! error varaible
     INTEGER, INTENT(OUT)
8
                                          :: info
                               _____
                                                        _____
     ! Declaration of local variables
                                   _____
     ! loop counters
9
     INTEGER
                                          :: redstep, i, j, k, isx, isy, isz
     ! index of current voxel
10
     INTEGER
                                          :: idx, idy, idz
     ! indices of voxel in reduced data set
     INTEGER
11
                                          :: ired, ired, kred
     ! number of voxels in x,y,z
     INTEGER
12
                                           :: nvx, nvy, nvz, npix, nones
     ! new number of voxels after reduction step
     INTEGER
                                           :: nvxred, nvyred, nvzred
13
     ! number of reduction steps to perform
     INTEGER
14
                                           :: nred
     ! error trap
15
     INTEGER
                                          :: istat
     ! Array to hold the voxel counts after each reduction step and its log
     INTEGER, DIMENSION(:), ALLOCATABLE
16
                                          :: Nvox
     REAL(MK), DIMENSION(:), ALLOCATABLE
17
                                          :: logN
     ! array to hold the box size of each step
18
     REAL(MK), DIMENSION(:), ALLOCATABLE
                                          :: s
     ! accumulated voxel values for low-pass average and their total number
19
     REAL(MK)
                                          :: val
20
     INTEGER
                                           :: nval
     ! temporary voxel array for resolution reduction
21
     INTEGER, DIMENSION(:,:,:), ALLOCATABLE :: voxtmp
     ! sum of all reduction steps and sum of their squares
22
     REAL(MK)
                                          :: sumi, sumi2
     ! determinant and its inverse
23
     REAL(MK)
                                          :: det, detinv
     ! slope and y-axis intersect of linear fit
24
     REAL(MK)
                                          :: a, b
25
     LOGICAL, DIMENSION(:,:,:), ALLOCATABLE
                                          :: flag
     ! voxel size at current reduction step
26
     REAL(MK)
                                          :: vx, vy, vz
     1-----
     ! Initialize and allocate memory
27
     if(thresh .GT. 1.0_MK .OR. thresh .LT. 0.0_MK) then
28
        \texttt{WRITE}(\texttt{*},\texttt{'}(\texttt{A})\texttt{'}) 'Threshold for BCdim must be between 0 and 1'
29
        info = 1
        return
30
31
     end if
32
     info = 0
33
     BCdim = 0.0_MK
34
     vx = vx0
35
     vy = vy0
     vz = vz0
36
     ! convert surface to voxel representation
     CALL Voxelize(vx0,vy0,vz0,nvx,nvy,nvz,info)
37
38
     if(info .NE. 0) then
39
        WRITE(*,'(A)') 'Error in Voxelize called from BCdim.'
40
       return
41
     end if
     ! reduce until the smallest dimension contains less than 16 voxels
```
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```
42
      nred = ceiling(log(real(minval((/nvx,nvy,nvz/)))/16.0_MK)/log(fact))
43
      istat = 0
44
      ALLOCATE(Nvox(nred+1), s(nred+1), logN(nred+1), STAT=istat)
45
      if(istat .NE. 0) then
46
         WRITE(*,'(A)') 'Error allocating memory for Nvox in BCdim'
47
         info = 1
48
         return
49
      end if
      ! Box counting algorithm by successive reduction of resolution
50
      do redstep=1,nred
51
         ALLOCATE(flag(nvx,nvy,nvz))
         flag = .FALSE.
52
         ! outline voxel set: only keep outer layer of voxels
53
         do i=1,nvx
                                         ! loop over all voxels
54
            do j=1,nvy
               do k=1,nvz
55
                  if(voxel(i,j,k) .NE. 0) then
56
                                                  ! if a voxel is 1, count
                                                  ! the number of its
57
                                                  ! neighbors that are 1 too
                     npix = 0
                     nones = 0
58
                     do isx=-1,1
59
                        idx = i+isx
60
61
                        do isy=-1,1
                           idy = j+isy
62
63
                           do isz=-1,1
                              idz = k+isz
64
65
                              if((idx .GT. 0 .AND. idx .LE. nvx) .AND.
                                                                        &
                              (idy .GT. 0 .AND. idy .LE. nvy) .AND. &
                              (idz .GT. 0 .AND. idz .LE. nvz)) then
                                 ! if inside array: check it
66
                                 npix = npix + 1
                                 if(voxel(idx,idy,idz) .EQ. 1) nones=nones+1
67
68
                              end if
69
                           end do
                        end do
70
                     end do
71
                     ! if all its neighbors are one, mark it for deletion
                     ! since it is not a border voxel
                     if(nones .EQ. npix) flag(i,j,k) = .TRUE.
72
73
                  end if
74
               end do
75
            end do
76
         end do
         ! delete all interior voxels
77
         where(flag) voxel=0
78
         DEALLOCATE(flag)
79
         Nvox(redstep) = sum(voxel)
80
         s(redstep) = vx
81
         vx = fact*vx
82
         vy = fact*vy
83
         vz = fact*vz
84
         WRITE(*,'(A,I2,A,I2,A,I7)') 'BCdim: reduction step ',redstep,' of &
            ',nred,'. voxels: ',Nvox(redstep)
         ! new voxel numbers after reduction step
85
         nvxred = ceiling((1.0_MK/fact)*real(nvx))
86
         nvyred = ceiling((1.0_MK/fact)*real(nvy))
         nvzred = ceiling((1.0_MK/fact)*real(nvz))
87
         ! memory for reduced voxel array
88
         ALLOCATE(voxtmp(nvxred,nvyred,nvzred), STAT=istat)
89
         if(istat .NE. 0) then
            WRITE(*,'(A)') 'Error allocating temporary voxel array in BCdim'
90
            info = 1
91
92
            return
93
         end if
         ! low-pass filtering: repalce each pixel by the average of its neighbors
94
         ired = 0
         jred = 0
95
         kred = 0
96
         ! loop over every second voxel
97
         do i=1,nvx,2
98
            do j=1,nvy,2
99
               do k=1,nvz,2
                  ! reset mean values
```

100	$val = 0.0_{MK}$
101	nval = 0
101	l loop over its neighbors
102	
102	
103	
104	do 159=-1,1
105	idy = j+isy
106	do isz=-1,1
107	idz = k+isz
108	if((idx .GT. 0 .AND. idx .LE. nvx) .AND. &
	(idv .GT. O .AND. idv .LE. nvv) .AND. &
	(idz .GT. 0 .AND. idz .LE. nvz)) then
	l if incide array, undate cliding mean
100	i il inside allay, update situing mean
110	
110	nval = nval + 1
111	end 11
112	end do
113	end do
114	end do
	! mean of surrounding voxel values
115	if(nval .GT. 0) val = val/real(nval)
	! set corresponding voxel in reduced data set to one if mean
	! is larger than threshold, else zero
116	ired = ceiling((1.0 MK/fact)*i)
117	<pre>ired = ceiling((1.0 MK/fact)*i)</pre>
118	$kred = ceiling((1 \cap MK/fact)*b)$
110	$\operatorname{vortmp}(\operatorname{ired} \operatorname{kred}) = 1$
120	vorump(iteu, jieu, reu) - i if(ual IT through) wastma(ined ined ined) - 0
120	ii(vai .Li. thresh) voxtmp(ired, jred, kred) = 0
121	end do
122	end do
123	end do
	! resize voxel array to reduced sizes
124	DEALLOCATE(voxel, STAT=istat)
125	if(istat .NE. 0) then
126	WRITE(*,'(A)') 'Error deallocating voxel array in BCdim'
127	info = 1
128	return
129	end if
130	ALLOCATE(voxel(nvxred.nvvred.nvzred), STAT=istat)
131	if(istat_NE_0) then
132	WRITE(* (A)) From reallocating your array in RCdim?
132	info = 1
100	
104	
135	ena li
	! assign new voxel values
136	voxel = voxtmp
137	nvx = nvxred
138	nvy = nvyred
139	nvz = nvzred
	! deallocate temporary voxel array
140	DEALLOCATE(voxtmp, STAT=istat)
141	if(istat .NE. 0) then
142	WRITE(*, (A)) 'Error deallocating temporary voxel array in BCdim'
143	info = 1
144	return
145	end if
146	end do
140	l also get last data point
1/7	$\frac{1}{2} \frac{1}{2} \frac{1}$
140	
148	s(nred+1) = VX
	! Logarithms of results and output to file
	!
149	do i=1,nred+1
150	$s(i) = log(1.0_MK/s(i))$
151	<pre>logN(i) = log(real(Nvox(i)))</pre>
152	end do
153	OPEN(60, FILE='BC.gnu', STATUS='REPLACE', ACTION='WRITE')
154	do i=1,nred+1
155	WRITE(60.*) s(i), logN(i)
156	end do
157	
101	
	:
	: Least squares fit to estimate the dimension (Slope)

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```
158 sumi = sum(s)
159
     sumi2 = sum(s(:)**2)
160 det = sumi2*(nred+1.0_MK)-sumi*sumi
161 detinv = 1.0_MK/det
162
     a = 0.0 MK
163 b = 0.0_MK
164 do i=0,nred
     a = a + ((nred+1.0_MK)*s(i+1)-sumi)*logN(i+1)
b = b + (sumi2-s(i+1)*sumi)*logN(i+1)
165
166
167
     end do
168
     a = detinv*a
169
     b = detinv*b
170
     WRITE(*, '(A,F8.5,A,F8.5)') 'Box counting results: a = ',a,', b = ',b
      ! the slope is the box counting dimension
171
     BCdim = a
      ! Write gnuplot macro (invoke using: gnuplot plotit.mac)
172 OPEN(60, FILE='plotit.mac', STATUS='REPLACE', ACTION='WRITE')
173
     WRITE(60,*) 'set key top left'
174
     WRITE(60, '(A,F9.6,A,F9.6,A)') 'p ''BC.gnu'' w p, ',a,'*x+',b,' w l'
175 WRITE(60,*) 'pause -1'
176
    CLOSE(60)
      ! free memory and terminate
      1---
     DEALLOCATE(Nvox)
177
     DEALLOCATE(s)
178
     DEALLOCATE(logN)
179
     if(ALLOCATED(voxel)) DEALLOCATE(voxel)
180
181
    return
182 END function BCdim
```

C.5.4.14 Convert text to upper case (UpperCase.f90)

UpperCase.f90 takes the two input arguments, string and ilen, and converts the character string string of length ilen to all upper case letters. The result is again contained in the variable string.

```
1
    subroutine UpperCase(string,ilen)
    IMPLICIT NONE
2
    !UpperCase Converts a string to all upper case characters.
      UPPERCASE takes the string "string" of length ilen and converts
      every character of it to upper case.
      See also ReadParams
      todo:
    1=====
        DIPLOMA THESIS WS01/02 ICOS
                                              ETH-ZUERICH
             PROTEIN DIFFUSION INSIDE THE ENDOPLASMIC RETICULUM
         1-----
         _____
    ! Input/Output arguments
                   _____
    ! the string to be converted
3
    CHARACTER(LEN=*), INTENT(INOUT)
                                :: string
    ! length of the string
```

```
4
    INTEGER, INTENT(IN)
                                     :: ilen
    1-----
    ! Declaration of local variables
                             _____
    ! loop counter
5
    INTEGER
                                     :: i, j
    ! alphabet boundaries and shift lower->upper
6
    INTEGER
                                    :: i1,i2,i3,iadd
                    _____
    ! uppercase
                        _____
    ! determine alphabet boundaries
7
    i1 = IACHAR(ia') - 1
       = IACHAR('z') + 1
8
    i2
    i3 = IACHAR('A')
9
    ! shift to upper case
    iadd = i3 - i1 - 1
10
    ! shift all lower case characters to upper case
11
    do i=1,ilen
12
      j = IACHAR(string(i:i))
      if (j.GT.i1.AND.j.LT.i2) then
    string(i:i) = CHAR(j+iadd)
13
14
      end if
15
    end do
16
    1----
              _____
    ! return
    1---
    9999 CONTINUE
17
18
    return
```

```
19 END subroutine UpperCase
```

C.5.4.15 Dynamic list structure handling (Util.f90)

Util.f90 contains an overloaded set of dynamic list reallocation functions as described in section 7.2.6 and [Numerical Recipes in Fortran 90 (1996)]. Currently, the function is overloaded for real vectors (reallocate_rv), real matrices (reallocate_rm), integer vectors (reallocate_iv), integer matrices (reallocate_im) and character vectors (reallocate_hv). However, only the one for integer vectors is used in this work.

```
module Util
1
2
   USE globals
   IMPLICIT NONE
3
                         !Util contains common utility functions.
     UTIL currently only implements an overloaded dynamic array
     reallocation function
     See also INIT_PART, SORTT
     todo:
   ! DIPLOMA THESIS WS01/02 ICOS
                                      ETH-ZUEBICH
                   _____
           PROTEIN DIFFUSION INSIDE THE ENDOPLASMIC RETICULUM
   1------
   ! Interface declaration of contained functions
```

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```
1------
Δ
     INTERFACE reallocate
5
        MODULE PROCEDURE reallocate_rv,reallocate_rm,reallocate_iv, &
             reallocate_im,reallocate_hv
6
     END INTERFACE
7
     CONTAINS
      ! Reallocate a pointer to a new size preserving its previous contents
      !--
      ! for real vectors
8
     function reallocate_rv(p,n)
9
       REAL(MK), DIMENSION(:), POINTER
                                                  :: p, reallocate_rv
10
       INTEGER(I4B), INTENT(IN)
                                                  :: n
       INTEGER(I4B)
11
                                                   :: nold, ierr
       ALLOCATE(reallocate_rv(n), STAT=ierr)
12
       if(ierr .NE. 0) then
13
          WRITE(*,'(A)') 'Error allocating memory in reallocate_rv'
14
15
          return
16
       end if
17
       if(.NOT. ASSOCIATED(p)) return
       nold = size(p)
18
       reallocate_rv(1:min(nold,n))=p(1:min(nold,n))
19
       DEALLOCATE(p)
20
21
     END function reallocate_rv
      ! for integer vectors
22
     function reallocate_iv(p,n)
23
       INTEGER(I4B), DIMENSION(:), POINTER
                                                  :: p, reallocate_iv
       INTEGER(I4B), INTENT(IN)
24
                                                   :: n
25
       INTEGER(I4B)
                                                   :: nold. ierr
26
       ALLOCATE(reallocate_iv(n), STAT=ierr)
27
       if(ierr .NE. 0) then
    WRITE(*,'(A)') 'Error allocating memory in reallocate_iv'
28
29
          return
30
       end if
       if(.NOT. ASSOCIATED(p)) return
31
       nold = size(p)
32
33
       reallocate_iv(1:min(nold,n))=p(1:min(nold,n))
34
       DEALLOCATE(p)
35
     END function reallocate_iv
     ! for character vectors
36
     function reallocate_hv(p,n)
37
       CHARACTER(1), DIMENSION(:), POINTER
                                                  :: p, reallocate_hv
38
       INTEGER(I4B), INTENT(IN)
                                                   :: n
39
       INTEGER(I4B)
                                                   :: nold, ierr
40
       ALLOCATE(reallocate_hv(n), STAT=ierr)
41
       if(ierr .NE. 0) then
42
          WRITE(*,'(A)') 'Error allocating memory in reallocate_hv'
43
          return
44
        end if
45
       if(.NOT. ASSOCIATED(p)) return
46
       nold = size(p)
47
       reallocate_hv(1:min(nold,n))=p(1:min(nold,n))
48
       DEALLOCATE(p)
49
     END function reallocate_hv
      ! for real matrices
50
     function reallocate_rm(p,n,m)
51
       REAL(MK), DIMENSION(:,:), POINTER
                                                  :: p, reallocate_rm
       INTEGER(I4B), INTENT(IN)
52
                                                   :: n, m
       INTEGER(I4B)
53
                                                   :: nold, mold, ierr
54
       ALLOCATE(reallocate_rm(n,m), STAT=ierr)
       if(ierr .NE. 0) then
55
56
          WRITE(*,'(A)') 'Error allocating memory in reallocate_rm'
57
          return
58
       end if
       if(.NOT. ASSOCIATED(p)) return
59
       nold = size(p,1)
60
       mold = size(p, 2)
61
62
       reallocate_rm(1:min(nold,n),1:min(mold,m))=p(1:min(nold,n),1:min(mold,m))
       DEALLOCATE(p)
63
```

END function reallocate_rm 64 ! for integer matrices 65 function reallocate_im(p,n,m) INTEGER(I4B), DIMENSION(:,:), POINTER INTEGER(I4B), INTENT(IN) 66 :: p, reallocate_im 67 :: n, m 68 INTEGER(I4B) :: nold, mold, ierr ALLOCATE(reallocate_im(n,m), STAT=ierr) 69 if(ier .NE. 0) then
 WRITE(*,'(A)') 'Error allocating memory in reallocate_im' 70 71 72 return 73 74 end if if(.NOT. ASSOCIATED(p)) return 75 76 nold = size(p,1)
mold = size(p,2)
reallocate_im(1:min(nold,n),1:min(mold,m))=p(1:min(nold,n),1:min(mold,m)) 77 78 DEALLOCATE(p) 79 END function reallocate_im

80 END module Util

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